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JANUS: A RECONFIGURABLE SYSTEM FOR SCIENTIFIC COMPUTING

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Contents

In	Introduction 1							
1	Intr	oducti	ion to reconfigurable computing	5				
	1.1	Gener	al purpose architectures	6				
	1.2	Doma	in-specific architectures	$\overline{7}$				
	1.3	Applie	cation-specific architectures	8				
	1.4	Progra	ammable logic, FPGA	9				
	1.5	Recon	figurable Computing	11				
		1.5.1	Pervasiveness of RC	13				
		1.5.2	The Hartenstein's point of view	13				
		1.5.3	Man does not live by hardware only	14				
	1.6	Non e	xhaustive history of RC	15				
		1.6.1	Common features	16				
		1.6.2	Fix-plus machine (Estrin)	17				
		1.6.3	Rammig Machine	17				
		1.6.4	Xputer (Hartenstein)	18				
		1.6.5	PAM, VCC and Splash	19				
		1.6.6	Cray XD1	21				
		1.6.7	RAMP (Bee2)	21				
		1.6.8	FAST (DRC)	22				
		1.6.9	High-Performance Reconfigurable Computing: Maxwell and Janus	23				
2	Mo	nte Ca	rlo methods for statistical physics	29				
	2.1			29				
		2.1.1	•	33				
		2.1.2	-	35				
		2.1.3		36				
	2.2			38				
		2.2.1		38				
		2.2.2	-	39				

		2.2.3	Metropolis algorithm	ļ						
		2.2.4	How to use the Metropolis algorithm for spin systems 42)						
		2.2.5	Another MC algorithm: the heat bath	j						
		2.2.6	Parallel tempering techniques	2						
	2.3	Nume	rical requirements)						
		2.3.1	Implementation and available parallelism	j						
		2.3.2	Techniques on a general purpose processor	,						
		2.3.3	Random numbers	I						
3	Janus architecture at large 5'									
	3.1	Janus	project	ļ						
	3.2	Questi	ions leading Janus's development	J						
		3.2.1	Why many nodes on a board?	ļ						
		3.2.2	Why an Input/Output processor?	J						
		3.2.3	How are organized communications between Janus boards and							
			Janus host?	-						
		3.2.4	How are organized communications within a Janus board? 64	2						
		3.2.5	Why a nearest neighbours network?	j						
		3.2.6	Why do boards have no direct link among them self? 67	,						
		3.2.7	Why only 17 nodes per board? 68	ļ						
		3.2.8	Why do the nodes have no off chip memory?	į						
		3.2.9	Which clock frequency and why? 69	I						
	3.3	SP firm	mware: spin glass	I						
		3.3.1	Parallelism	I						
		3.3.2	Algorithm Implementation	-						
	3.4	SP firm	mware: parallel tempering	2						
	3.5	SP firm	mware: graph coloring	,						
		3.5.1	Memory organization	,						
		3.5.2	Janus limitations in graph coloring	I						
4	Arc	hitectı	ral details of Janus 85	,						
	4.1	FPGA	.: different flavours)						
	4.2	Struct	ure of a Janus board	,						
		4.2.1	SP	,						
		4.2.2	IOP	J						
		4.2.3	PB	-						
		4.2.4	Janus box	2						
	4.3	The I	$OP in depth \ldots $	ļ						
		4.3.1	Clock handling: topClock	2						
		4.3.2	Double data rate: iddrBus and oddrBus	,						

		4.3.3 Message routing: StreamRouter					
		4.3.4 Memory controller: memExt					
		4.3.5 SP reconfiguration interface: mainProgInt					
		4.3.6 SP communication: spInt					
		4.3.7 Synchronization device: syncInt $\ldots \ldots \ldots$					
	4.4	SP firmwares for Janus test					
	4.5	Engineering problems					
5	Per	formance and results 111					
	5.1	Useful concepts					
		5.1.1 About the equilibrium					
		5.1.2 Correlation $\ldots \ldots \ldots$					
		5.1.3 Order parameters: magnetization and overlap					
	5.2	First run details					
	5.3	Janus performance					
	5.4	Physics results overview					
		5.4.1 Non-equilibrium dynamics of a large EA spin glass 119					
		5.4.2 The 4-state Potts model and its phase structure					
Co	onclu	sions 129					
\mathbf{A}	Not	es on the IOP communication strategy 133					
	A.1	Overview of the IOP structure					
	A.2	First idea: the stuff byte					
		A.2.1 Which stuff-value?					
		A.2.2 Performance problem					
	A.3	Second idea: the tagged stream					
		A.3.1 Remarks					
	A.4	Third idea: the encapsulated stream					
		A.4.1 Pros and cons					
Ri	Ringraziamenti 139						

Introduction

The widespread diffusion of field programmable gate arrays (FPGA) and their remarkable technological developments have allowed reconfigurable computing to play an increasingly important role in computer architectures. This trend should be very evident in the field of high performance, as the possibility to have a huge number of gates that can be configured as needed opens the way to new approaches for computationally very intensive tasks.

Even if the reconfigurable approach has several advantages, its impact so far has been limited. There are several reasons that may explain why reconfigurable computing is still in a corner:

- i) the costs in terms of time and the specific technical skills needed to develop an application using programmable logic is by far higher than those associated to programming an application in an appropriate programming language, even considering a reasonable amount of (possibly platform-specific) optimization for performance;
- ii) strictly correlated with the point above, software tools performing a (more or less) automatic and transparent transitions from a conventional computer approach to a reconfigurable computing structure are still in an embryonic phase and in a too fragmentary stage of development;
- iii) the interface between reconfigurable devices and standard processors is often not standard and therefore almost any project developing a system housing an FPGA side by side to a conventional architecture defines a new data-exchange protocol, so even the simplest communication primitives cannot be standardized. While this lack of standardization does provide opportunities for those willing (and able to) consider innovative designs, it is only perceived as a further obstacle for any attempt to provide a standard development environment for reconfigurable systems.

In spite of these drawbacks, there are several areas of computational sciences where the reconfigurable approach is able to provide such a large performance boost as to provide adequate compensation to the disadvantages described above. Hardware resources of recent FPGA generations allow us to map complex algorithms directly within just one device, configuring the available gates in order to perform a computationally heavy tasks with high efficiency. In some cases, moreover, these applications have computing requirements that are large enough to justify even huge development efforts.

A paradigmatic example of this situation is the study of a particular class of theoretical physics models called *spin glasses* performed using Monte Carlo methods. The algorithms relevant in this field are characterized by *i*. large intrinsic parallelism that allows one to implement a trivial SIMD approach (i.e. many Monte Carlo update engines within a single FPGA); *ii*. relatively small size of the computational data base (~2 MByte), that can be fully stored into on-chip memories; *iii*. large use of good quality random numbers (up to 1024 per clock cycle); *iv*. integer arithmetic and simple logic operations. Careful tailoring of the architecture to the specific features of the algorithms listed above makes it possible to reach impressive performance levels: just one FPGA has the same performance as ~ 1000 standard PC with a recent state-of-the-art processor (this will be explained in details in chapter 5).

The Janus project was started approximately 3 years ago in order to harvest all the potential advantages offered by reconfigurable computing for spin-glass simulations. Janus is a collaboration among the Spanish Universities of Zaragoza, Madrid and Extremadura, the BIFI Institute of Zaragoza and the Italian Universities of Ferrara and Roma I with the industrial partnership of the Eurotech Group. The main aim of the project is to build an FPGA based supercomputer strongly oriented to study and solve the computational problems associated to the simulation of the spin systems introduced above.

Janus is a system composed of three logical layers. The hardware layer includes several (16 in the first system, deployed in December 2007) boards each housing 17 FPGA-based subsystem: 16 so-called scientific processors (SPs) and one input/output processor (IOP). A standard PC (called the Janus host) connects to up to two Janus boards and controls their functionalities via the IOP module.

At the software layer we find the communication libraries, developed in order to allow the user to interface his applications with Janus, and the physics libraries, a set of routines written in C that simplifies the operations of setting up of a lattice spin simulation on Janus. By we also developing these libraries, in some sense we define an interface between our FPGA-based system and a general purpose processorbased computer. We obviously need such an interface to operate our machine, but we concede that in this way we give our fair contribution to increasing the entropy of reconfigurable computing interfaces (on the other hand, since we deal with statistical physics, we know very well that decreasing entropy is a more formidable tasks than writing a PhD thesis).

The third layer is composed of the firmware for the FPGAs running the simulation codes, that is the set of parametric SP firmware that implements different spin models and the IOP-based firmware that includes several IO interfaces, a memory controller, a configuration controller driving the configuration of the SPs and several debug interfaces. All these firmware modules are handcrafted in VHDL. No automatic translation tools from high level languages to hardware description languages has been used, since we consider optimization of the usage of logic resources and computational performance as our primary goal.

The Janus project started in 2004 with preliminary meetings between Italian researchers and the Spanish group that built in the '80s another FPGA based machine, called *Spin Update Engine* (SUE). During my laurea degree I studied spin models and implemented a Monte Carlo simulation engine on an FPGA: this initial experiment, that we named *SuperSUE*, assessed the viability and the expected performances of a smassively parallel system based on latest generation FPGAs based.

In summer 2006 the Eurotech group assembled the first three prototype Janus boards, using Xilinx Virtex-4 LX160 FPGAs. One year later, the first Janus rack, powered by 256 Xilinx Virtex-4 LX200 based computational nodes was tested. Acceptance tests on this large system ended before Christmas 2007 and in march 2008 we performed the first large scale run (a simulation stretching uninterruptedly for over 25 days with just one system crash of a couple of hour due to severe weather conditions that caused a power failure). The results of this run were reported in our application for the 2008 Gordon Bell Prize. Unfortunately, at least one Gordon Bell referee made the argument that only floating point performance is relevant for that award.

The Janus installation in Zaragoza was officially unveiled in June 2008, and, since then, has been always in operation, running several physics codes.

I was involved in all the phases of the project. In details during my first two PhD years I was involved at the hardware layer of the project, developing the overall hardware design of the system, working on the detailed structure of all its subsystems and developing procedures for hardware tests. This period was characterized by a strong interaction with engineers of the Eurotech group, the company that actually built our hardware. Another relevant work of this period was the development of the firmware for the IOP and some firmwares modules for the SPs, that we needed to test all implemented hardware functionalities. The final period of my PhD studies was dedicated to realize a complete test bench in order to validate the system as it was assembled.

I also worked on some preliminary studies for the Janus implementation of an efficient firmware for random graph coloring.

This thesis has 5 chapters, structured as follows:

In Chapter 1 I review basic concepts of reconfigurable computing and I include a non exhaustive history of the projects that marked developments in this area.

Chapter 2 is an introduction to the physics concepts and simulation algorithms for which Janus has been designed and built. I introduce the Edward-Anderson spin model, the graph coloring problem and the Monte Carlo algorithms used to investigate them (Metropolis, Heat Bath and Parallel Tempering).

Chapter 3 is dedicated to a general discussion of the Janus architecture. I briefly describe hardware components and then I discuss a set of important architectural questions that I handled during the development of Janus prototypes.

In Chapter 4 I highlight selected significant details about the Janus basic hardware elements, about the IOP and SP firmware and a brief description of the test environment developed to check the Janus boards.

Chapter 5 is a short review of the most important physics results obtained so far with Janus, including a detailed analysis of performances.

My work is wrapped up in the concluding chapter.

When you got nothing, you got nothing to lose.

Bob Dylan

1

Introduction to reconfigurable computing

Research in the architecture of computer systems has always played a central role in the computer science and high performance computing communities. The investigation goals vary according to the target applications, the price of the final equipment, the programmability and the scalability of the system and many others.

Until a few years ago for processors to be used in parallel machines for high-performance computing the focus was placed on high clock rates, parallelism of different chips and high communication bandwidth at the expense of power. Recently, however, attention has focused on multi core and many core architectures with the aim of taking advantage of the presence on-chip of more than one complex calculation unit, adding to the historical problem of the needs of bandwidth between chips the new challenge of a careful handling of parallelism among cores within a single chip.

On the other hand in the embedded systems environment the governing factor during development is in many cases the price of the final equipment and the main aim is to use optimized components in order to contain costs and optimize power consumption. A small microcontroller is used, for instance, to control data acquisition from sensors and provide data to a collector system at a very low frequency. In those systems the architectures are focused on containing power consumption and costs and are in some cases carefully tailored for specific areas.

A third example in which the architecture plays a key role is the environment in which some specific duties should be executed extremely efficiently depending on a small set of constraints. An example could be a rover used to explore a given area: its architecture will be carefully tailored to the specific application of image collection, elaboration and transmission and, at the same time, to obstacle detection.

We consider the three examples given above as illustrating of the big scenario of the architectures that can be analyzed by splitting it in three main groups:

- general purpose architectures based on the Von Neumann computing paradigm;
- *domain-specific architectures* designed for class of applications having common features;
- application-specific architectures developed for only one specific application;

1.1 General purpose architectures

In 1945, the mathematician John Von Neumann demonstrated in a study of computation [1] that is possible to have a simple fixed architecture able to execute any kind of computation, given a properly programmed control, without the need for hardware modification. The Von Neumann contribution was universally adopted and quickly became the groundwork of future generations of high-speed digital computers. One of the reasons for the acceptance of the Von Neumann approach is its simplicity of programming that follows the sequential way of human thinking. The general structure of a Von Neumann machine as shown in Figure 1.1 consists of:

- A memory for storing program and data. Harvard architectures contain two parallel accessible memories for storing program and data separately.
- A control unit (also called control path) featuring a program counter that holds the address of the next instruction to be executed.
- An arithmetic and logic unit (also called data path) in which instructions are executed.

A program is coded as a set of instructions to be executed sequentially, instruction after instruction. At each step of the program execution, the next instruction is fetched from the memory at the address specified in the program counter and decoded. The required operands are then collected from the memory before the instruction is executed. After execution, the result is written back into the memory. In this process, the control path is in charge of setting all signals necessary to read from and write to the memory, and to allow the data path to perform the right computation. The data path is controlled by the control path, which interprets the instructions and sets the control signals accordingly to execute the desired operation.

In general, the execution of an instruction on a Von Neumann computer can be done in five cycles: Instruction Read (IR) in which an instruction is fetched from the

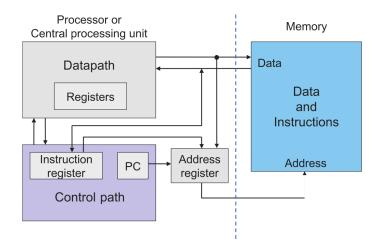


Figure 1.1 – A scheme of the Von Neumann computer architecture (source [2]).

memory; Decoding (D) in which the meaning of the instruction is determined and the operands are localized; Read Operands (R) in which the operands are read from the memory; Execute (EX) in which the instruction is executed with the read operands; Write Result (W) in which the result of the execution is stored back to the memory. In each of those five cycles, only the part of the hardware involved in the computation is activated. The rest remains idle. For example if the IR cycle is to be performed, the program counter will be activated to get the address of the instruction, the memory will be addressed and the instruction register to store the instruction before decoding will be also activated. Apart from those three units (program counter, memory and instruction register), all the other units remain idle.

Decades of research in computer architectures developed techniques to optimize the organization of processors as the ones described above. The extraction of the instruction level parallelism, the so called *pipelining*, the multiple issue architectures, the branch prediction or the optimized instruction scheduling are only a few examples of the breakthroughs in this subject. For an exhaustive overview of architectural optimizations see [3].

1.2 Domain-specific architectures

A domain-specific processor is a processor tailored for a class of algorithms. As mentioned in the previous section, the data path is tailored for an optimal execution of a common set of operations that mostly characterizes the algorithms in the given class. Also, memory access is reduced as much as possible. Digital Signal Processor (DSP) are among the most used domain-specific processors.

A DSP is a specialized processor used to speed-up computation of repetitive, numerically intensive tasks in signal processing areas such as telecommunication, multimedia, automobile, radar, sonar, seismic, image processing, etc. The most often cited feature of DSPs is their ability to perform one or more multiply accumulate (MAC) operations in single cycle. Usually, MAC operations have to be performed on a huge set of data. In a MAC operation, data are first multiplied and then added to an accumulated value. The normal Von Neumann computer would perform a MAC in 10 steps. The first instruction (multiply) would be fetched, then decoded, then the operand would be read and multiply, the result would be stored back and the next instruction (accumulate) would be read, the result stored in the previous step would be read again and added to the accumulated value and the result would be stored back. DSPs avoid those steps by using specialized hardware that directly performs the addition after multiplication without having to access the memory.

Because many DSP algorithms involve performing repetitive computations, most DSP processors provide special support for efficient looping. Often a special loop or repeat instruction is provided, which allows a loop implementation without expending any instruction cycles for updating and testing the loop counter or branching back to the top of the loop. DSPs are also customized for data with a given width according to the application domain. For example if a DSP is to be used for image processing, then pixels have to be processed. If the pixels are represented in Red Green Blue (RGB) system where each colour is represented by a byte, then an image processing DSP will not need more than 8 bit data path. Obviously, the image processing DSP cannot be used again for applications requiring 32 bits computation.

This specialization of DSP's functions increases the performance of the processor and improves device utilization, but reduces the execution efficiency of an arbitrary application.

1.3 Application-specific architectures

Although DSPs incorporate a degree of application-specific features such as MAC and data width optimization, they still hide the Von Neumann approach and, therefore, remain sequential machines. Their performance is limited. If a processor has to be used for only one application, which is known and fixed in advance, then the processing unit could be designed and optimized for that particular application. In this case, we say that the hardware "fits" itself to the application. This kind of approach is useful, for instance, when a processor has to perform tasks defined by a standard, such as encoding and decoding of an audio/video stream.

A processor designed for only one application is called an Application-Specific Processor (ASIP). In an ASIP, the instruction cycles (IR, D, R, EX, W) are eliminated: there is no fetch of instructions because the instruction set of the application is directly implemented in hardware, or, in other words the algorithm to perform is hardwired in a custom processor. Therefore a data stream works as input, the processor performs the required computation and the results can be collected at the outputs of the processor.

ASIPs use a spatial approach to implement only one application. The gates building the final processor are configured so that they constitute all the functional units needed for the computation of all parts of the application. This kind of computation is called *spatial computing* [4]. Once again, an ASIP that is built to perform a given computation cannot be used for other tasks other than those for which it has been originally designed.

ASIPs are usually implemented as single chips called *Application-Specific Integrated Circuit*, ASIC, or using devices housing programmable logic. This approach arose in the late '80's with the widespread commercial availability of reconfigurable chips called *Field Programmable Gate Arrays*, FPGAs.

1.4 Programmable logic, FPGA

The FPGA is a regularly tiled two-dimensional array of logic blocks. Each logic block includes a Look-Up Table (LUT), a simple memory that can store an arbitrary *n*-input boolean function. The logic blocks communicate through a programmable interconnection network that includes both nearest neighbor as well as hierarchical and long path wires. The periphery of the FPGA contains I/O blocks to interface between the internal logic blocks and the I/O pins. This simple, homogeneous architecture has evolved to become much more heterogeneous, including on-chip memory blocks as well as DSP blocks such as multiply/multiply-accumulate units.

There are several sorts of FPGAs, including those that can be programmed only once, but the application-specific architectures may require that the device can be reconfigured on-the-fly during a run or between separate runs to obtains different behaviours. Depending on the needs of the applications it is possible to use devices basing their reconfiguration on SRAM (faster) or FLASH (slower) but this is only a technological detail. In both cases this means that the configuration of the FPGA, the *object code* defining the algorithm loaded onto the device, is stored in an on-chip storage device. By loading different configurations into this configuration device, different algorithms can be executed. The configuration determines the boolean function computed by each logic block and the interconnection pattern between logic and I/O blocks.

FPGA designers have developed a large variety of programmable logic structures for FPGAs since their invention in the mid-1980's. For more than a decade, much of the programmable logic used in FPGAs can be generalized as shown in Figure 1.2. The basic logic element generally contains some form of programmable combinational logic, a flip-flop or latch, and some fast carry logic to reduce the area and delay costs

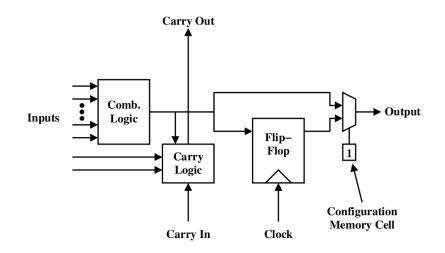


Figure 1.2 – A generic programmable logic block (source [5]).

for implementing carry logic. In our generic logic block, the output of the block is selectable between the output of the combinational logic or the output of the flip-flop. The figure also illustrates that some form of programming, or configuration, memory is used to control the output multiplexer; of course, configuration memory is used throughout the logic block to control the specific function of each element within the block.

In addition to the relatively *fine-grained* configurability provided by FPGAs and similar devices, the drive to reduce the power, area, and/or delay costs of fine-grained reconfigurability has led to a number of what may be called *coarse-grained* reconfigurable logic devices. Instead of providing configurability at the level of individual gates, flip-flops or look-up tables (LUTs), these coarse-grained architectures often provide arithmetic logic units (ALUs) and other larger functions that can be combined to perform computations. In the extreme, the functions might be as large as microprocessor cores such as in the Raw chip from MIT [6].

With their introduction in 1985, FPGAs have been an alternative for implementing digital logic in systems. The earlier use of the FPGAs were to provide a denser solution for glue logic within systems, but now they have expanded their applications to the point that it is common to find FPGAs as the central processing devices within systems. The reason of their increased diffusion and use lies mainly in the available resources embedded within a single chip: most of the FPGA family of the main brand offers in fact today not only logic resources, but also embedded memories, DSP block, high speed IO pins, hardwired IP core for the interface with PCI or other standard communication protocol. Compared with application-specific integrated circuits (ASICs), FPGAs have several advantages for their users, including: quick time to market, being a standard product; no non-recurring engineering costs for fabrication; pre-tested silicon for use by the designer; and reprogrammability, allowing designers to upgrade or change logic through in-system programming. By reconfiguring the device with a new circuit, design errors can be fixed, new features can be added, or the function of the hardware can be entirely retargeted to other applications. Of course, compared with ASICs, FPGAs cost more per chip to perform a particular function so they are not good for extremely high volumes. Also, an FPGA implementation of a function is slower than the fixed-silicon options.

1.5 Reconfigurable Computing

From the discussion in sections 1.1 1.2 1.3, where we introduced three different kinds of processing units, we can identify two main means to characterize processors: flexibility and performance.

The computers based on Von Neumann paradigm are very flexible because they are in principle able to compute any kind of task: therefore we refer to them with the terminology general purpose processors. Although there are many kind of optimizing procedures and tricks their general purpose orientation has a cost in terms of performance: for instance, the five steps (IR, D, R, EX, W) needed to perform one instruction becomes a major drawback, in particular if the same instruction has to be executed on huge sets of data; moreover their intrinsic sequential structure is useful for the programmer because it is similar to the human thought process but is a natural hindrance for a possible parallel computing approach for some applications. With this architecture we have thus a high level of *flexibility* because the hardware structure is fixed and, in many cases, is hidden to the programmer by the compiler that play the role to "fit" the application in the hardware in order to be executed. We could use the catchphrase: with general purpose processors the application must always fits in the hardware.

On the other side the application-specific architectures bring high performance because they are optimized for a particular application. The instruction set required for that application can then be built in a chip, but we pay a high cost in terms of flexibility. In this case the important goals are the performance of the processor and the hardware is shaped by the application. From this is possible we can invent the opposite catchphrase: in presence of application-specific architectures the hardware always fits in the application.

Between these two extreme positions, general purpose processors and applicationspecific processors, there is, architecturally speaking, an interesting *space* in which we find different types of processors. We can classify them depending on their performance and their flexibility.

If we consider, after this analysis, the features of the FPGAs introduced briefly in section 1.4 we can easily see that they allow us to implement hardware architectures that merge the flexibility of a general purpose processor and the performance of an application-specific processor with the comfort of the reconfigurability. In other words, the boost that FPGA technology gives to researchers studying the architectures a powerful tool to try to efficiently fill the space between general purpose and application-specific processors. We consider therefore FPGAs as the way to build a reconfigurable hardware or reconfigurable device or *Reconfigurable Processing Unit*, RPU, in analogy with the Central Processing Unit, CPU. Following this, the study of computation using reconfigurable devices is commonly called *Reconfigurable Computing*.

For a given application, at a given time, the spatial structure of the device will be modified such as to use the best computing approach to speed up that application. If a new application has to be computed, the device structure will be modified again to match the new application. Contrary to the Von Neumann computers, which are programmed by a set of instructions to be executed sequentially, the structure of reconfigurable devices are changed by modifying all or part of the hardware at compiletime or at run-time, usually by downloading a so called bitstream into the device. In this sense we call *configuration* or *reconfiguration* the process of changing the structure of a reconfigurable device respectively at star-up-time or at run-time.

Other than this difference of approach the major operative differences between reconfigurable and processor-based computing are:

- The FPGA is configured into a customized hardware implementation of the application. The hardware is usually data path driven, with minimal control flow; processor-based computing depends on a linear instruction stream including loops and branches.
- The reconfigurable computer data path is usually pipelined so that all function units are in use every clock cycle. The microprocessor has the potential for multiple instructions per clock cycle, but the delivered parallelism depends on the instruction mix of the specific program, and function units are often underutilized.
- The reconfigurable computer can access many memory words in each clock cycle, and the memory addresses and access patterns can be optimized for the application. The processor reads data through the data cache, and efficiency of the processor is determined by the degree to which data is available in the cache when needed by an instruction. The programmer only indirectly controls the cache-friendliness of the algorithm, as access to the data cache is hidden from the instruction set architecture.
- The FPGA has in principle no constraints about the size of data words: the words of the data path can have arbitrary length. Using the general purpose architectures the length of the data words is fixed and the programmer has no fine control on it.

To summarize, reconfigurable computing is concerned with decomposing applications into spatially parallel, tiled, application-specific pipelines, whereas the traditional general purpose processor interprets a linear sequence of instruction, with pipelining and other forms of spatial parallelism hidden within the microarchitecture of the processor.

Progress in reconfiguration has been amazing in the last two decades. This is mostly due to the wide acceptance of the Field Programmable Gate Array (FPGAs) that are now established as the most widely used reconfigurable devices.

1.5.1 Pervasiveness of RC

There are two main fields in which the reconfigurable computing has been mostly accepted, developed and used: embedded computing and scientific computing.

There are *social reasons* for that analyzed in [7]: people with expertise developing embedded systems have hardware background and see in FPGAs a cheap solution to develop custom systems with the possibility to fix with new firmware releases their projects so that the development and the update of an FPGA-based system is easier and faster than with other components. Also this category of people with hardware background is in many cases familiar with hardware description languages and with hardware implementation techniques and can obtain impressive boost of performance with a relative low prices and power consumption.

From scientific computing come problems often with very special requirements and in many cases it is possible to implement their algorithms directly within an FPGA. Depending on the project the FPGAs can be configured as the co-processor of a general purpose processor or a main-custom processor. A limiting factor for to this kind of project comes often from the absence of a specific hardware background of the people involved: in some cases in fact the need to develop a design using hardware description languages can require a long development period and a drastic change of the paradigm of programming. Because of this some groups and companies develop and sell tools for the *translation* of standard codes, like C, for instance, to various hardware description language. These automatic tools of translation can speed up and make easier the development process, but sometimes have big limitations on the code structures that can be translated and the extraction of parallelism may not always be efficient.

1.5.2 The Hartenstein's point of view

The most common architectural approach in computer science is the Von Neumann paradigm and the most used processors are based on the general purpose architectures. The reconfigurable computing approach requires a deep change of programming paradigm in comparison with the Von Neumann: R. Hartenstein starting from 1990 gave a formalisation of it in [8, 9, 10]. Theoretically speaking, there is in the RC environment a different view of the software and the hardware. In the Von Neumann approach the hardware resources are fixed, only one source code is needed and the compiler processing it generates an instruction-stream to be stored in the program memory waiting the scheduled time to be executed.

The RC approach requires two different type of source code, *configware* and *flowware*. Configware is commonly written using some abstraction of an hardware description language and is synthesised using a tool that translates the code describing the custom architecture into logic gates, maps it on the FPGA resources and produces as output a so called bitstream, a configuration file to properly set up the FPGA resources. Flowware is code written with a high level programming language generating a stream of data used as input for the custom architecture implemented within the FPGA. As sometimes the platform housing the reconfigurable device is not a "standard" mother board with "standard" communication protocols, the flowware implements also communication interfaces and other features useful for the system. Reconfigurable systems require moreover that configuration/reconfiguration of the logic is performed external to the device: in some cases a PROM is used to set up the reconfigurable device on boot, but it is useful to have the possibility to reconfigure the devices "on the fly". This require that flowware is able to perform this task too. Figure 1.3 summarize the theoretical schemes of two approach: Von Neumann and reconfigurable computing.

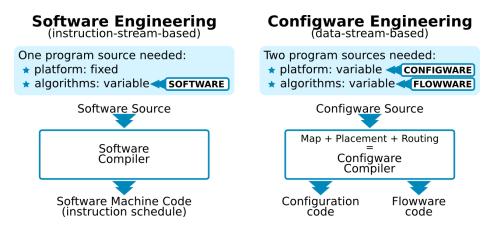


Figure 1.3 – On the left the organization of Software Engineering in the classic Von Neumann point of view; on the right a schema of the Configure Engineering theorized by Hartenstein (adapted from [7]).

1.5.3 Man does not live by hardware only

On the theoretical side Hartenstein tries to formalize the principles of reconfigurable computing. Following this idea and thanks to the increasing resources offered by the FPGAs many people built therefore systems based on programmable logic and many of these projects developed interfaces to connect general purpose architecture and FPGA based systems [11, 12, 13, 14] trying to define a standard interface in order to:

- access reconfigurable hardware resources without introducing undesirable dependencies on hardware;
- avoid client code changes whenever the hardware is revised;
- leave the programmer free to know or ignore the hardware details of interfaces or low level protocols;
- develop a set of libraries optimized for scientific computing and reprogrammable logic based coprocessors [12].

All these efforts are focused on solving a basic problem coming from reconfigurable computing and formalized by Hartenstein: which is the way to efficiently use the huge degrees of freedom coming from the FPGA while maintaining programmability accessible to a large part of the computer science community?

A first approach to solve it is to completely ignore the programmability and the generalization of the design in order to obtain the best performances from the logics: this approach is commonly adopted by groups or projects developing efficient applicationspecific architectures that are not interested in developing a general purpose machine, but only a performance oriented custom machine.

The second opposite approach comes from groups and projects studying systems oriented to a large enough set of applications and in which the availability of as friendly a programming environment as possible can justify a considerable loss of performance.

1.6 Non exhaustive history of RC

Like most technologies, reconfigurable computing systems are built on a variety of existing technologies and techniques. It is always difficult to pinpoint the exact moment a new area of technology comes into existence or even to pinpoint which is the first system in a new class of machines. Popular scientific history often gives simple accounts of individuals and projects that represent a turning point for a particular technology, but in reality the story is usually more complicated.

The large number of exotic high-performance systems designed and built over a very short time makes this area particularly difficult to document, but there is also a problem specific to them. Much of the work was done inside various government agencies, particularly in the United States, and was never published. In these cases, all that can be relied on is currently available records and publications.

1.6.1 Common features

Reconfigurable systems are distinguished from other cellular multiprocessor systems. Array processors, in particular Single Instruction Multiple Data Processors (SIMD), are considered architecturally distinct from reconfigurable machines, in spite of many similarities. This distinction arises primarily from the programming techniques. Array processors tend to have either a shared or dedicated instruction sequencer and take standard instruction-style programming code. Reconfigurable machines tend to be programmed spatially, with different physical regions of the device being configured at different times. This necessarily means that they will be slower to reprogram than cellular multiprocessor systems but should be more flexible and achieve higher performance for a given silicon area.

Although FPGA-based systems are very different each others, it is possible to extract some shared features concerning both architectures of systems housing reconfigurable devices and architecture of designs within the FPGAs.

The systems using FPGA are in many cases configured as master-slave systems: a general purpose architecture works as master, runs the flowware and handles communications with the slave system with one or more reconfigurable devices. In general slave systems are custom boards housing in addition to one or more FPGAs other components like for instance external memory, communications devices, special I/O devices etc. Although the structure of the system housing reconfigurable logics could change among different projects, the global scheme of a system involved FPGAs can be summarize in general as show in Figure 1.4.



Figure 1.4 – A generic reconfigurable system is composed by a **Slave System** housing a programmable device (FPGA) and a **Master System** allowing the (re-)configuration of the programmable device and handling the communications with it.

As said before, reconfigurable devices are programmed spatially, thus different regions have different tasks and each region can be programmed (configured) in different times. Despite this large degree of freedom, some constraints coming from the chip vendors are fixed for the developer: position of the I/O blocks, distribution of internal memories, clock drivers and clock trees are common problems for a developer using FPGA. Moreover, even if a programmer finds the infrastructure already built, a fixed structure is forced by the design of the system housing the FPGA so that some areas of the chip are reserved for logic blocks performing specific task, as for instance I/O interfaces or memory controllers. A common feature of all reconfigurable computing projects is therefore the presence of some kind of spatial constraints.

1.6.2 Fix-plus machine (Estrin)

In 1959, Gerald Estrin, a computer scientist of the university of California at Los Angeles, introduced the concept of reconfigurable computing. The following fragment of an Estrin publication in 1960 [15] on the fix-plus machine, defines the concept of reconfigurable computing paradigm.

"Pragmatic problem studies predicts gains in computation speeds in a variety of computational tasks when executed on appropriate problemoriented configurations of the variable structure computer. The economic feasibility of the system is based on utilization of essentially the same hardware in a variety of special purpose structures. This capability is achieved by programmed or physical restructuring of a part of the hardware."

To implement this vision, Estrin designed a computing system, the fix-plus machine, that like many reconfigurable computing systems available today, was composed of a fixed architecture (a proto-general purpose processor) and a variable part consisting of logic operators that could be manually changed in order to execute different operations.

1.6.3 Rammig Machine

In the year 1977, Franz J. Rammig, a researcher at the university of Dortmund proposed a concept for editing hardware [16]. The goal was:

"investigation of a system, which, with no manual or mechanical interference, permits the building, changing, processing and destruction of real (not simulated) digital hardware."

Rammig realised his concept by developing a hardware editor similar to today's FPGA architecture. The editor was build upon a set of modules, a set of pins and a one-to-one mapping function on the set of pins. The circuitry of a given function was then defined as a *string* on an alphabet of two letters (w = wired and u = unwired). To build the hardware editor, *selectors* were provided with the modules' outputs connected to the input of the selectors and the output of the selectors connected to the input of the modules. The overall system architecture is shown in Figure 1.5.

The implementation of the *{wired, unwired}* property was done through a programmable crossbar switch, made upon an array of selectors. The bit strings were provided by storing the selector control in registers, and by making these registers accessible from a host computer, the PDP11 in those days. The modules were provided

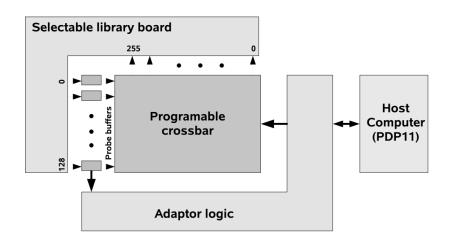


Figure 1.5 – Structure of the Rammig Machine (source [2]).

on a library board similar to that of Estrin's Fix-Plus. Each board could be selected under software control. The mapping from module I/Os to pins was realized manually, by a wiring of the provided library boards, i.e. fixed per library board.

1.6.4 Xputer (Hartenstein)

The Xputer's concept was presented in early 1980s by Reiner Hartenstein, a researcher at the University of Kaiserslautern in Germany [8].

The goal was to have a very high degree of programmable parallelism in the hardware, at the lowest possible level, to obtain performance not possible with the Von Neumann computers. Instead of sequencing the instructions, the Xputer would sequence data, thus exploiting the regularity in the data dependencies of some class of applications like image processing, where repetitive processing is performed on a large amount of data. An Xputer consists of three main parts: the data sequencer, the data memory and the *reconfigurable ALU*, rALU, that permits the run-time configuration of communication at levels below instruction set level. Within a loop, data to be processed were accessed via a data structure called the *scan window*. Data manipulation was done by the rALU that had access to many scan windows. The most essential part of the data sequencer was the *generic address generator*, GAG, that was able to produce address sequences corresponding to the data of up to three nested loops. An rALU subnet that could be configured to perform all computations on the data of a scan window was required for each level of a nested loop.

The general XPuter architecture is presented in Figure 1.6. This shows the realization of the XPuter as a *map oriented machine*, MoM [17]. The overall system was made upon a host processor, whose memory was accessible by the MoM. The rALU subnets received their data directly from local memory or from the host main memory via the MoM bus. Communication was also possible among the rALUs via direct serial

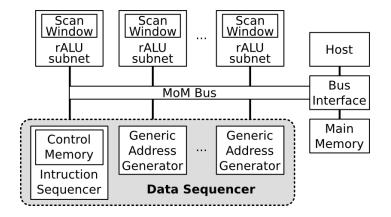


Figure 1.6 – General architecture of the XPuter as implemented in the Map oriented Machine (MOM-3) prototype.

connections. Several XPuters could also be connected to provide more performance.

For executing a program, the hardware had to be configured first. If no reconfiguration took place at run-time, then only the data memory would be necessary. Otherwise, a configuration memory would be required to hold all the configurations to be used at run-time.

The basic building block of the reconfigurable ALU was the so-called *reconfigurable datapath unit*, rDPU. Several rDPUs were used within an rALU for data manipulation. Each rDPU had two registered inputs and two registered outputs with a data width of 32 bit. Input data were provided either from the north or from the west, while the south and east were used for the output. Besides the interconnection lines for the rALUs, a global I/O-Bus is available for the connection of designs to the external world. The I/O bus was principally used for accessing the scan windows.

The control implemented a program that is loaded on reconfiguration to control different units of the rALU. Its instruction set consisted of instructions for loading the data as well as instructions for collecting results from the field. Application of the XPuters was in image processing, systolic array and signal processing.

1.6.5 PAM, VCC and Splash

In the late 1980s, PAM, VCC, and Splash, three significant general-purpose systems using multiple FPGAs, were designed and built. They were similar in that they used multiple FPGAs, communicated to a host computer across a standard system bus, and were aimed squarely at reconfigurable computing.

The *Programmable Active Memories*, PAM, project [18] at Digital Equipment Corporation (DEC) initially used four Xilinx XC3000-series FPGAs. The original Perle-0 board contained 25 Xilinx XC3090 devices in a 5×5 array, attached to which were four independent banks of fast static RAM (SRAM), arranged as $64K \times 64$ bits, which were

controlled by an additional two XC3090 FPGA devices. This wide and fast memory provided the FPGA array with high bandwidth. The Perle-0 was quickly upgraded to the more recent XC4000 series. As the size of the available XC4000-series devices grew, the PAM family used a smaller array of FPGA devices, eventually settling on 2×2 . Based at the DEC research lab, the PAM project ran for over a decade and continued in spite of the acquisition of DEC by Compaq and then the later acquisition of Compaq by Hewlett-Packard. PAM, in its various versions, plugged into the standard PCI bus in a PC or workstation and was marked by a relatively large number of interesting applications as well as some groundbreaking work in software tools. It was made available commercially and became a popular research platform.

The Virtual Computer from the Virtual Computer Corporation, VCC, [19] was perhaps the first commercially available reconfigurable computing platform. Its original version was an array of Xilinx XC4010 devices and I-Cube programmable interconnect devices in a checkerboard pattern, with the I-Cube devices essentially serving as a crossbar switch. The topology of the interconnection for these large FPGA arrays was an important issue at this time: With a logic density of approximately 10K gates and input/output (I/O) pins on the order of 200, a major concern was communication across FPGAs. The I-Cube devices were perceived as providing more flexibility, although each switch had to be programmed, which increased the design complexity. The first Virtual Computer used an 8×8 array of alternating FPGA and I-Cube devices. The exception was on the left and right sides of the array, which exclusively used FPGAs, which consumed 40 Xilinx XC4010 FPGAs and 24 I-Cubes. Along the left and right sides were 16 banks of independent $16 \times 8K$ dual-ported SRAM, and attached to the top row were 4 more banks of standard single-ported $256K \times 32$ bits SRAM controlled by an additional 12 Xilinx XC4010 FPGAs. While this system was large and relatively expensive, and had limited software support, VCC went on to offer several families of reconfigurable systems over the next decade and a half.

The Splash system [20, 21], from the Supercomputer Research Center (SRC) at the Institute for Defense Analysis, was perhaps the largest and most heavily used of these early systems. Splash was a linear array consisting of XC3000-series Xilinx devices interfacing to a host system via a PCI bus. Multiple boards could be hosted in a single system, and multiple systems could be connected together. Although the Splash system was primarily built and used by the Department of Defense, a large amount of information on it was made available. A Splash 2 [22] system quickly followed and was made commercially available from Annapolis Microsystems. The Splash 2 board consisted of two rows of eight Xilinx XC4010 devices, each with a small local memory. These 16 FPGA/memory pairs were connected to a crossbar switch, with another dedicated FPGA/memory pair used as a controller for the rest of the system. Much of the work using Splash concentrated on defense applications such as cryptography and pattern matching, but the associated tools effort was also notable, particularly some of the earliest high-level language (HLL) to hardware description language (HDL) translation software targeting reconfigurable machines. Specifically, the data parallel C compiler and its debug tools and libraries provided reconfigurable systems with a new level of software support.

PAM, VCC, and Splash represent the early large-scale reconfigurable computing systems that emerged in the late 1980s. They each had a relatively long lifetime and were upgraded with new FPGAs as denser versions became available. Also of interest is the origin of each system. One was primarily a military effort (Splash), another emerged from a corporate research lab (PAM), and the third was from a small commercial company (Virtual Computer). It was this sort of widespread appeal that was to characterize the rapid expansion of reconfigurable computing systems during the 1990s.

1.6.6 Cray XD1

While the number of small reconfigurable coprocessing boards would continue to proliferate as commercial FPGA devices became denser and cheaper, other new hardware architectures were produced to address the needs of large-scale supercomputer users. Unlike the earlier generation of boards and systems that sought to put as much reconfigurable logic as possible into a single unified system, these machines took a different approach. In general, they were traditional multiprocessor systems, but each processing node in them consisted of a very powerful commercial desktop microprocessor combined with a large commercial FPGA device. Another factor that made these systems unique is that they were all offered by mainstream commercial vendors.

The first reconfigurable supercomputing machine from Cray, the XD1 [23], is based on a chassis of 12 processing nodes, with each node consisting of an AMD Opteron processor. Up to 6 reconfigurable computing processing nodes, based on the Xilinx Virtex-2 Pro devices, can also be configured in each chassis, and up to 12 chassis can be combined in a single cabinet, with multiple cabinets making larger systems. Hundreds of processing nodes can be easily configured with this approach.

1.6.7 RAMP (Bee2)

Around 2005 the choice of the computer hardware industry to focus production on single-chip multiprocessors gave a boost to the idea of developing a system able to simulate highly parallel architectures at hardware speeds. The *Research Accelerator* for Multiple Processors, RAMP, is the open-source FPGA-based project that arose from this idea [24, 25]: its main aim is to develop and share the hardware and software necessary to create parallel architectures.

The computational support of RAMP project is the system BEE2 [26] using Xilinx Virtex-2 Pro FPGAs as primary and only processing elements. A peculiarity of this system is the PowerPC 405 embedded on the FPGA that makes it possible to minimize latency between microprocessor and reconfigurable logic while maximizing the data throughput. Moreover the BEE2 system is an example of an FPGA-based system that does not require explicitly a master system checking over the tasks of the reconfigurable logic: each FPGA embeds in fact general purpose processors able to control itself.

Each BEE2 compute module consists of five Xilinx Virtex-2 Pro-70 FPGA chips, each directly connected and logically organized into four compute FPGAs and one control FPGA. The control FPGA has additional global interconnect interfaces and control signals to the secondary system components, while the compute modules are connected as a 2×2 mesh.

The architecture of the BEE2 leaves some degrees of freedom and using the 4X Infiniband physical connections, the compute modules can be wired into many network topologies, such as a 3D mesh. For applications requiring high-bisection-bandwidth random communication among many compute modules, the BEE2 system is designed to take advantage of commercial network switch technology, such as Infiniband or 10G Ethernet. The compute module runs the Linux OS on the control FPGA with a full IP network stack. Moreover each BEE2 system is equipped with high bandwidth memories (DDR, DDR2) and other I/O interfaces.

As well as being a hardware architecture project, RAMP aims to support the software community as it struggles to take advantage of the potential capabilities of parallel microprocessors, by providing a malleable platform through which the software community can collaborate with the hardware community.

1.6.8 FAST (DRC)

FPGA-Accelerated Simulation Technologies (FAST) [27], is a today's project developed by the University of Texas at Austin that attempts to speed up the simulation of complex computer architectures. It gives a methodology to build extremely fast, cycle-accurate full system simulators that run real applications on top of real operating systems. Current state of the project allows one to boot unmodified Windows XP, Linux 2.4 and Linux 2.6 and run unmodified applications on top of those operating systems at simulation speeds in the 1.2 MIPS range, between 100 and 1000 times faster than Intel's and AMD's cycle-accurate simulators (e.g. which is fast enough to type into Microsoft Word). I knew people of this project during my visit at the University of Texas in summer 2008.

The hardware platform used to develop this project is a DRC development system (DS2002). This machine contains a dual-socket motherboard, where one socket contains an AMD Opteron 275 (2.2GHz) and the other socket contains a Xilinx Virtex-4

LX200 (4VLX200) FPGA. The Opteron communicates to the FPGA via HyperTransport. The functional model runs on the Opteron and the timing model runs on the FPGA. DRC provides libraries to read and write from the FPGA. Interesting feature regarding the hardware platform is the fact that an FPGA uses a standard socket of a general purpose processor and not custom interfaces.

1.6.9 High-Performance Reconfigurable Computing: Maxwell and Janus

The high-performance computing field is traditionally dominated by clusters of general purpose processors and a common approach of the scientists is to find a machine as fast as possible to run a given code, if possible with no changes of it. This approach do not require in principle an understanding of the architecture or the hardware features of the machine running the code. On the other hand it is known that code optimization with respect to architectural details improves the performance of applications.

Despite that there are studies about the viability of reconfigurable supercomputing [28] and some projects with relevant results in this field.

The FPGA High Performance Computing Alliance (FHPCA) [29] was established in 2004 and is dedicated to the use of Xilinx FPGAs to deliver new levels of computational performance for real-world industrial applications. Led by EPCC, the supercomputing centre at The University of Edinburgh, the FHPCA is funded by Scottish Enterprise and builds on the skills of Nallatech Ltd, Alpha Data Ltd, Xilinx Development Corporation, Algotronix and ISLI.

Maxwell [30, 31, 32] is a high-performance computer developed by the FHPCA to demonstrate the feasibility of running computationally demanding applications on an array of FPGAs. Not only can Maxwell demonstrate the numerical performance achievable from reconfigurable computing, but it also serves as a testbed for tools and techniques to port applications to such systems.

The unique architecture of Maxwell comprises 32 blades housed in an IBM Blade Center. Each blade comprises one 2.8 GHz Xeon with 1 Gbyte memory and 2 Xilinx Virtex-4 FPGAs each on a PCI-X subassembly developed by Alpha Data and Nallatech. Each FPGA has either 512 Mbytes or 1 Gbyte of private memory. Whilst the Xeon and FPGAs on a particular blade can communicate with each other over the PCI bus (typical transfer bandwidths in excess of 600 Mbytes/s), the principal communication infrastructure comprises a fast Ethernet network with a high-performance switch linking the Xeons together and RocketIO linking the FPGAs. Each FPGA has 4 RocketIO links enabling the 64 FPGAs to be connected together in an 8×8 toroidal mesh. The RocketIO has a bandwidth in excess of 2.5 Gbits/s per link.

Together these two principal interconnect subsystems enable the efficient imple-

mentation of parallel codes where there is a need both for intensive numerical processing and for fast data communication between the cooperating processing elements.

The Parallel Toolkit developed by EPCC supports the decomposition of a numerically intensive application into a set of cooperating modules running on the array of Xeons in much the same way that many applications can be decomposed to run on a cluster of PCs. Each module can then be further analysed to identify the numerical "hot spots" which are then implemented on the FPGAs taking advantage of the fast RocketIO linking the FPGAs for fast communications. The implementation of the numerically intensive parts of the applications is accomplished using a combination of tools such as DIME-C from Nallatech, Handel-C from Celoxica and VHDL available from several vendors including Xilinx.

Janus is a project among universities of Italy and Spain with the main goal to realize a FPGA-based parallel system optimized a specific class of statistical physics simulations. Janus is composed of 256 Xilinx Virtex-4 FPGA organized in sets of 16 each connected via raw-ethernet Gigabit channel to a standard PC.

Both these projects give to the computer science community an efficient proof that reconfigurable computing can be used in order to obtain high performance machines as with a general purpose environment, in the Maxwell case, as in a very special and application specific case, in the Janus case. I discuss in depth details of Janus in the next chapters.

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They tried to take me to a doctor, but its too late for me. Then they took me to a preacher that they saw on their tv who said that for a small donation my lost soul would be saved I said I don't think so preacher, I'll come back another day.

Bon Jovi

2

Monte Carlo methods for statistical physics

This chapter is about the use of computers to solve problems in statistical physics. In particular, it is about Monte Carlo methods, which form the largest and most important class of numerical methods used for solving statistical physics problems.

In the opening section, I look first at what we mean by statistical physics, giving a brief overview of the discipline called statistical mechanics with special care to two examples: the Edward-Anderson model and the Potts model applied to the problem of a random graph coloring. These two models are in fact implemented with impressive boost of performances on the Janus supercomputer. The material presented here is largely inspired by references [1, 2] and [3].

In section 2.2 I introduce Monte Carlo methods in general and I explain how they can be used to explore statistical mechanics problems.

The last section will be dedicated to the computational features of Monte Carlo implementation on a standard architecture. A special care will be given to the problem of random numbers generation.

2.1 Statistical Physics

Statistical mechanics is primarily concerned with the calculation of properties of condensed matter systems. The crucial difficulty associated with these systems is that they are composed of very many parts, typically atoms or molecules. These parts are usually all the same or of a small number of different types and they often obey quite simple equations of motion so that the behaviour of the entire system can be expressed mathematically in a straightforward manner. However the complexity of the problem makes it impossible to solve the mathematics exactly. A standard example is that of a volume of gas in a container. One litre of, for instance, oxygen at standard temperature and pressure consists of about 3×10^{22} oxygen molecules, all moving around and colliding with one another and the walls of the container. One litre of air under the same conditions contains the same number of molecules, but they are now a mixture of oxygen, nitrogen, carbon dioxide. The atmosphere of the Earth contains 4×10^{24} litres of air, or about 1×10^{44} molecules, all moving around and colliding with each other and with the environment. It is not feasible to solve Hamilton's equations for these systems because there are too many equations, and yet when we look at the macroscopic properties of the gas, they are very well-behaved and predictable. Clearly, there is something special about the behaviour of the solutions of these many equations that "averages out" to give us a predictable behaviour for the entire system. For example, the pressure and temperature of the gas obey quite simple laws although both are measures of rather gross average properties of the gas. Statistical mechanics attempts to avoid the problem of solving the equations of motion and to compute these gross properties of large systems by treating them in a probabilistic fashion. Instead of looking for exact solutions, we deal with the probabilities of the system being in one state or another hence the name statistical mechanics. Such probabilistic statements are extremely useful, because we usually find that for large systems the range of behaviours of the system that have a non negligible probability to occur is very small; all the reasonably probable behaviours fall into a narrow range, allowing us to state with extremely high confidence that the real system will display behaviour within that range.

The typical paradigm for the systems we will be studying in this section is one of a system governed by a Hamiltonian function \mathcal{H} which gives us the total energy of the system in any particular state. We only consider systems that have discrete sets of states each with its own energy, ranging from the lowest, or ground state energy E_0 upwards, E_1 , E_2 , E_3 , ..., possibly without limit.

If the system were in insulation energy would be conserved, which means that the system would stay in the same energy state all the time (or if there were a number of degenerate states with the same energy, maybe it would make transitions between those) However there is another component to our paradigm, and that is the thermal reservoir. This is an external system which acts as a source and sink of heat, constantly exchanging energy with our Hamiltonian system in such a way as always to push the temperature of the system, defined as in classical thermodynamics, towards the temperature of the reservoir. In effect the reservoir is a weak perturbation on the Hamiltonian, which we ignore in our calculation of the energy levels of our system, but which pushes the system frequently from one energy level to another. We can incorporate the effects of the reservoir in our calculations by giving the system a dynamics, a rule whereby the system changes periodically from one state to another. The exact nature of the dynamics is dictated by the form of the perturbation that the reservoir produces in the Hamiltonian. However, there are a number of general conclusions that we can reach without specifying the exact form of the dynamics, and we will examine these first.

Suppose our system is in a state μ . Let us define $R(\mu \to \nu)dt$ to be the probability that it is in state ν a time dt later. $R(\mu \to \nu)$ is the transition rate for the transition from μ to ν . The transition rate is normally assumed to be time independent and we will make that assumption here. We can define a transition rate like this for every possible state ν that the system can reach. These transition rates are usually all we know about the dynamics, which means that even if we know the state μ that the system starts off in, we need only wait a short interval of time and it could be in any one of a very large number of other possible states. This is where our probabilistic treatment of the problem comes in. We define a set of weights $w_{\mu}(t)$ which represent the probability that the system will be in state μ at time t. Statistical mechanics deals with these weights, and they represent our entire knowledge about the state of the system. We can write a master equation for the evolution of $w_{\mu}(t)$ in terms of the rates $R(\mu \to \nu)$ thus:

$$\frac{dw_{\mu}}{dt} = \sum_{\nu} [w_{\nu}(t)R(\nu \to \mu) - w_{\mu}(t)R(\mu \to \nu)]$$
(2.1.1)

The first term on the right-hand side of this equation represents the rate at which the system is undergoing transitions into state μ the second term is the rate at which it is undergoing transitions out of μ into other states. The probabilities $w_{\mu}(t)$ must also obey the sum rule

$$\sum_{\mu} w_{\mu}(t) = 1 \tag{2.1.2}$$

for all t, since the system must always be in some state. The solution of Equation 2.1.1, subject to the constraint 2.1.2, tells us how the weights w_{μ} vary over time.

We must now consider how the weights w_{μ} relate to the macroscopic properties of the system which we study. If we are interested in some quantity Q, which takes the value Q_{μ} in state μ , then we can define the expectation of Q at time t for our system as

$$\langle Q \rangle = \sum_{\mu} Q_{\mu} w_{\mu}(t). \qquad (2.1.3)$$

Clearly this quantity contains important information about the real value of Q that we might expect to measure in an experiment. For example, if our system is definitely in one state τ then $\langle Q \rangle$ will take the corresponding value Q_{τ} . And if the system is equally likely to be in any of three states, and has zero probability of being in any other state, then $\langle Q \rangle$ is equal to the mean of the values of Q in those three states. However, the precise relation of $\langle Q \rangle$ to the observed value of Q must be considered more closely. There are two ways to look at it. The first is to imagine having a large number of copies of our system all interacting with their own thermal reservoirs and making transitions between one state and another all the time. $\langle Q \rangle$ is then a good estimate of the number we would get if we were to measure the instantaneous value of the quantity Q in each of these systems and then take the mean of all of them. This is a conceptually sound approach to defining the expectation of a quantity. However this definition is not closely related to what happens in a real experiment. In a real experiment we normally only have one system and we make all our measurements of Q on that system, though we probably don't just make a single instantaneous measurement, but rather integrate our results over some period of time. There is another way of looking at the expectation value which is similar to this experimental picture. This is to envisage the expectation as a time average of the quantity Q. Imagine recording the value of Q every second for a thousand seconds and taking the average of those one thousand values. This will correspond roughly to the quantity calculated in Equation 2.1.3 as long as the system passes through a representative selection of the states in the probability distribution w_{μ} , in those thousand seconds. Obviously we will get an increasingly accurate fit between our experimental average and the expectation $\langle Q \rangle$ as we average over longer and longer time intervals. Conceptually there is a weakness in this approach as we do not have a rigorous definition of what we mean by a "representative selection of the states". There is no guarantee that the system will pass through anything like a representative sample of the states of the system in the time during which we observe it. It could easily be that the system only moves from one state to another on longer time scales and so it remains in the same state for all of our measurements. Or maybe it changes state very rapidly, but because of the nature of the dynamics spends long periods of time in small portions of the state space. This can happen for example if the transition rates $R(\mu \rightarrow \nu)$ are only large for states of the system that differ in very small ways, so that the only way to make a large change in the state of the system is to go through very many small steps. This is a very common problem in a lot of the systems. Another potential problem with the time average interpretation of 2.1.3 is that the weights $w_{\mu}(t)$, which are functions of time, may change considerably over the course of our measurements, making the expression invalid. For equilibrium systems the weights are by definition not time-varying, so this problem does not arise. Despite these problems however, this interpretation of the expectation value of a quantity is

the most widely used and most experimentally relevant interpretation. The calculation of expectation values is one of the fundamental goals of statistical mechanics, and of Monte Carlo simulation in statistical physics.

2.1.1 Spin Glass

Spin glass materials are diluted magnetic materials with long range interactions with oscillating signs, that can be accurately described by statistical mechanical models where the quenched couplings have zero average and are randomly distributed: they have played a crucial role in the development of a new paradigm [4]. Recently the replica symmetry breaking (RSB) solution of the mean field theory has been proven to be correct [5, 6], but many of its implications, already at the mean field level, have yet to be unveiled, and its relations with realistic models of three dimensional materials are not yet clear.

This paradigm is of remarkable interest for at least four different reasons. First the theoretical structure that emerges from the RSB is very complex and very new, different from the most part of typical features of physical systems. Second it sheds at least some lights on materials in the universality class of spin glasses (that is far larger of the one including magnetic samples). Third it will hopefully allow us to advance in our understanding of structural glasses [7], where in the absence of quenched disorder it is the intrinsic complexity of the Hamiltonian which dynamically induces frustration. Fourth it allows us to analyze (and sometimes to brilliantly solve) a large number of non physics problems: optimization, financial markets and complex networks can be studied thanks to the tools and algorithms created for studying spin glasses.

The typical Hamiltonian is

$$\mathcal{H} = -\sum J_{ij} s_i s_j , \qquad (2.1.4)$$

where the s variables are the basic degrees of freedom of the model and can take the values ± 1 (models where the variables are defined on a μ dimensional spheres are for example also of interest), *i* and *j* are sites of a *D* dimensional lattice (typically simple cubic) and the couplings *J* are quenched random variables defined on the links of the lattice (among first neighboring sites in *D* dimensions or among all couples of sites for the mean field theory) that can take both signs and have a zero average (they can be for example ± 1 with equal probability or normal variables).

We consider a spin glass system as a collection of spins (i.e. magnetic moments) that presents a frozen disordered low temperature state, instead of the uniform or periodic pattern usually found in magnetic systems. In order to produce such a state, two key ingredients turn out to be necessary: frustration (i.e. competition among the different interactions between magnetic moments, so that they cannot be all satisfied simultaneously) and disorder in the form of quenched randomness (i.e. each interaction

term has a roughly equal a priori probability to be ferromagnetic or antiferromagnetic, and its value is fixed on experimental times).

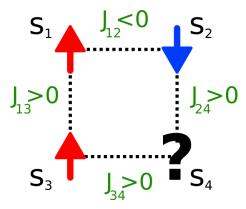


Figure 2.1 – Example of frustration in a 2D Ising spin system.

Figure 2.1 shows a simple case of frustration. In this example we consider four interacting Ising spins, represented by arrows pointing up or down (the only two available states for an Ising spin). The interaction values J are depicted as positive or negative on the bonds, indicating respectively a ferromagnetic or antiferromagnetic interaction between the spins connected by the bond. Two spins coupled by a positive interaction J > 0 tend to align themselves parallel to each other, while an antiferromagnetic interaction tends to favor an anti-parallel alignment. Figure 2.1 shows a case where the competition between the interactions makes it impossible to minimize the energy of all the bonds at the same time. The system is thus said to be frustrated.

The other feature of a spin glass system is the disorder usually given by the couplings J distributed with a given probability distribution or by an external random field applied to the system.

The appearance of many competing configurations in spin glasses is caused by the random distribution of interactions within the system. To explain the presence of different types of interactions within a single sample of material it is instructive to consider the first kind of glassy systems studied. These were mostly diluted solutions of magnetic impurities of a transition metal (such as Fe or Mn) in non-magnetic (noble) metal hosts (as Au or Cu).

The impurities cause a magnetic polarization of the surrounding electrons of the host metal. This interaction oscillates rapidly with the distance, producing a polarization field which is positive at some distances and negative at others. The surrounding impurity moments interacts with this magnetic field and try to align themselves in accordance to it. Because of the oscillatory behavior of these magnetic fields, and the random placement of the impurities in the host, some of the interactions end up being positive, and thus favor parallel alignment of the spins, while some others are negative, preferring antiparallel alignment. This, together with the variable distance between impurities, results in a set of random, competing interactions.

2.1.2 Edward-Anderson model

This model was introduced by Edwards and Anderson in 1975. One of their most important contributions consisted in being able to capture the very essence of the spin glass phenomenon within a model apparently simple and beautifully minimal, and yet offering a rich and complex phenomenology at the same time. This boosted the theoretical work on the field, and paved the way for other simplified models.

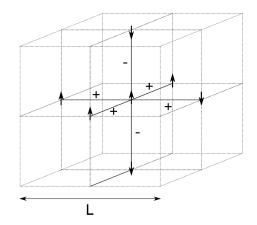


Figure 2.2 – Three dimensional cubic lattice. The magnetic moments, placed in all lattice nodes, are Ising-like variables ± 1 (i.e. up/down). In this case the interactions between nearest neighbors (represented by continuous lines in the figure) are taken from a bimodal distribution, and can be either ferromagnetic (+) or antiferromagnetic (-) (source [3]).

The Edwards-Anderson (EA) model is defined on a regular three dimensional cubic lattice of size L, usually with periodic boundary conditions and with the spins S_i lying on the $V = L^3$ sites i of the lattice (see Figure 2.2). Its Hamiltonian can be written as

$$\mathcal{H} = -\sum_{\langle i,j \rangle} J_{ij} s_i s_j , \qquad (2.1.5)$$

where $\langle i, j \rangle$ means that the sum is performed only over nearest neighbors, and the interactions J_{ij} are independent random variables. In general the precise form of the distribution of the J_{ij} is not very important, as far as it produces a disordered and frustrated system, but typically one uses either a Gaussian distribution with mean J_0 and standard deviation unity

$$P(J_{ij}) = \frac{1}{\sqrt{2\pi}} e^{-(J_{ij} - J_0)/2},$$
(2.1.6)

or a double delta distribution with zero mean

$$P(J_{ij}) = \frac{1}{2} [\delta(J_{ij} - 1) + \delta(J_{ij} + 1)].$$
(2.1.7)

The spins \mathbf{s}_i are classical *m*-component vectors (here indicated with bold font). The simplest version of the model considers the case with m = 1, the so-called Ising model, but a lot of theoretical work has been done also with m = 2 (the XY model), m = 3 (the Heisenberg model) and even higher number of components. In the literature the name *Edwards-Anderson models* is often used to indicate exclusively the version with Ising-like (i.e. m = 1 and $s_i = \pm 1$) spins and nearest neighbors interactions. We will usually refer to this model in the following sections, and will therefore use the scalar notation s_i for the spin, instead of the vectorial \mathbf{s}_i .

The EA model can be easily generalized to accommodate the contribution of an external magnetic field h whose value can vary randomly from one point of the lattice to another. It is indicated as h_i (as a function of the site i) and the resulting generalized Hamiltonian function is

$$\mathcal{H} = -\sum_{\langle i,j \rangle} J_{ij} s_i s_j - \sum_i h_i s_i . \qquad (2.1.8)$$

The apparent simplicity of the EA model, yet showing such a rich behavior, motivated a lot of researchers to develop alternative models to help advancements in spin glass theory. Many variations and different models have been introduced and studied during the last decades, such as, for instance the Sherrington-Kirkpatrick (SK) model [8, 4] in which each spin interacts with all others spins and not only with the nearest neighbours.

2.1.3 The Potts model and random graph coloring

The standard p-state Potts glass model is defined by the Hamiltonian

$$\mathcal{H} = -\sum_{\langle i,j \rangle} J_{ij} s_i s_j , \qquad (2.1.9)$$

where $s_i \in \{1, 2, ..., p\}$, and the couplings $\{J\}$ are random values. As it appears, an energy term $-J_{ij}$ is gained (or lost, depending on the sign of J_{ij}) if the two spins are in the same state, while the contribution is zero for spin in different states. The spins lie on a three-dimensional lattice of linear size L.

It is common to redefine the Hamiltonian by making use of the simplex representation of the Potts spins [9] which maps the p states of the (scalar) spins s_i to vectors $\mathbf{s_i}$ in a (p-1)-dimensional space, pointing to the corners of a p – simplex². The pvectors defined in this way satisfy the relation

$$s_a s_b = \frac{p\delta_{ab} - 1}{p - 1} \tag{2.1.10}$$

where a and b are scalar Potts states (thus $a, b \in \{1, 2, ..., p\}$).

This choice leads to the modified Hamiltonian

$$\mathcal{H} = -\sum_{\langle i,j\rangle} J'_{ij} \mathbf{s}_i \mathbf{s}_j , \qquad (2.1.11)$$

which is analogous to a vector spin model in (p-1) dimensions, with the substantial difference that the spin vectors can take only a discrete set of values. Furthermore, except for the case with p = 2 (which is just the EA-Ising model), given a spin s_i its inverse $-s_i$ is not allowed.

The simplex Hamiltonian is commonly used to study this class of models, since it helps to define the relevant observables and order parameters. The couplings J_{ij} are usually drawn from a Gaussian or a bimodal distribution, as for the EA model. It is not clear whether the choice on the distribution leading the disorder does affect or not the system behavior. It has been suggested [10] that the Gaussian and the bimodal Potts glass might belong to different universality classes, though numerical simulations like those in [11] and [12] show incompatible results and a quite different behavior.

The Potts model is relevant not only because it represents a generalization of EA model but also because it is an alternative formulation of the coloring problem on a random graphs. A random graph is obtained by starting with a set of n vertices and adding edges between them with a given different probability distributions. The Erdös-Rényi model for the random graphs is a common model [13, 14] in which the graph is denoted as G(V, E), a set of V vertices and E edges occurring with a probability p. For any graph G = (V, E), the set E of edges of G may be understood as a binary relation on V. This is the adjacency relation of G, in which vertices a and b are related precisely if $\{a, b\} \in E$, so ab is an edge of G. Conversely, every symmetric relation R on V gives rise to (and is the edge set of) a graph on V (see [15] for an exhaustive review). If we consider a set of p colors and we assume that each vertex has one of these colors, we can define a *conflict* in a colors configuration when two node connected by an edge (adjacent vertices) have the same color. The coloring problem consists of finding a color configuration of V, the set of the vertices, so that each vertex has a color and each edge has no conflicts. The reader can surely guess that also in the coloring problem there may exist the sort of frustration close to the definition given in Figure 2.1.

Analytic work [16] and numerical studies in [17] [18] prove that there exist interesting relations between the Potts model studying glassy systems in physics and the Erdös-Rényi model solving the coloring problem on random graphs. The basic idea of this equivalence is given in Figure 2.3: the p states of a spin in a Potts model correspond to p colors of a vertex in a random graph model. Therefore it is possible to define a Hamiltonian function as in 2.1.9 for a random graph too and to study a random graph with techniques similar to the ones used for spin systems. The only

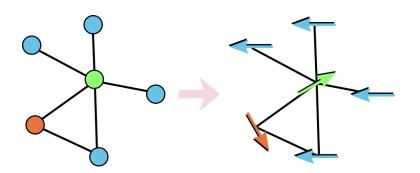


Figure 2.3 – Analogy between random graph coloring problem and Potts model.

difference lies in the geometry of the lattice: in a spin system the lattice is regular and each spin has a fixed number of neighbours with which it interacts (in a 3D lattice the neighbours are 6); by contrast in a random graph the number of adjacent nodes of a given node is given by a distribution of probability and is not fixed.

From the perspective of the scientist studying the random graph problem, this feature related to the topology of the system is not constrictive and approaching a random graph with spin model techniques is useful and interesting because it allows one to discover and prove important properties of random graphs. On the other hand, if we look at the topology of the system with the eyes of the computer scientist, we note immediately that while in spin system the regular lattice can be very efficiently mapped into a two-dimensional structure (i.e. a memory in a computer), an irregular structure of information as given by the random graph problem represents a non trivial challenge if we think about the manner to efficiently store and access data in a memory.

I studied this kind of problems during the prototyping phase of the Janus system [19] as described in details in Section 3.5.

2.2 Monte Carlo in general

2.2.1 Markov processes

Markov processes are stochastic processes whose futures are conditionally independent of their pasts given their present values. More formally, a stochastic process $\{X_t, t \in \mathscr{T}\}$, with $\mathscr{T} \subseteq \mathbb{R}$, is called a Markov process if, for every s > 0 and t,

$$(X_{t+s}|X_u, u \leqslant t) \sim (X_{t+s}|X_t) \tag{2.2.1}$$

In other words, the conditional distribution of the future variable X_{t+s} , given the entire past of the process $\{X_u, u \leq t\}$, is the same as the conditional distribution of X_{t+s} given only the present X_t . That is, in order to predict future states, we only need to know the present one. Property (2.2.1) is called the Markov property.

Depending on the index set \mathscr{T} and state space \mathscr{E} (the set of all values the $\{X_t\}$ can take), Markov processes come in many different forms. A Markov process with a discrete index set is called a *Markov chain*. A Markov process with a discrete state space and a continuous index set (such as \mathbb{R} or \mathbb{R}_+) is called a Markov jump process.

2.2.2 Markov chains

Consider a Markov chain $X = \{X_t, t \in \mathbb{N}\}$ with a discrete (that is, countable) state space \mathscr{E} . In this case the Markov property (2.2.1) is:

$$P(X_{t+1} = x_{t+1} | X_0 = x_0, ..., X_t = x_t) = P(X_{t+1} = x_{t+1} | X_t = x_t)$$
(2.2.2)

for all $x_0, ..., x_{t+1} \in \mathscr{E}$ and $t \in \mathbb{N}$. We restrict ourselves to Markov chains for which the conditional probability

$$P(X_{t+1} = j \mid X_t = i), i, j \in \mathscr{E}$$
(2.2.3)

is independent of the time t. Such chains are called *time-homogeneous*. The probabilities in (2.2.3) are called the *(one-step) transition probabilities* of X. The distribution of X_0 is called the initial distribution of the Markov chain. The one-step transition probabilities and the initial distribution completely specify the distribution of X. If we consider the product rule of probability for any sequence of events A_1, A_2, \ldots, A_n ,

$$P(A_1 \cdots A_n) = P(A_1)P(A_2|A_1)P(A_3|A_1A_2) \cdots P(A_n|A_1 \cdots A_{n-1}), \qquad (2.2.4)$$

using the abbreviation $A_1A_2 \cdots A_k \equiv A_1 \cap A_2 \cap \cdots \cap A_k$ and the Markov property 2.2.1 we obtain

$$P(X_0 = x_0, \dots, X_t = x_t) =$$

$$= P(X_0 = x_0)P(X_1 = x_1 | X_0 = x_0) \cdots P(X_t = x_t | X_0 = x_0 \dots X_{t-1} = x_{t-1})$$

$$= P(X_0 = x_0)P(X_1 = x_1 | X_0 = x_0) \cdots P(X_t = x_t | X_{t-1} = x_{t-1})$$
(2.2.5)

Since \mathscr{E} is countable, we can arrange the one-step transition probabilities in an array. This array is called the *(one-step)* transition matrix of X denoted by T. For example, when $\mathscr{E} = \{0, 1, 2, ...\}$ the transition matrix T has the form

$$T = \begin{pmatrix} p_{00} & p_{01} & p_{02} & \cdots \\ p_{10} & p_{11} & p_{12} & \cdots \\ p_{20} & p_{21} & p_{22} & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$
(2.2.6)

Note that the elements in every row are positive and sum up to unity. Another convenient way to describe a Markov chain X is through its transition graph. Figure 2.4

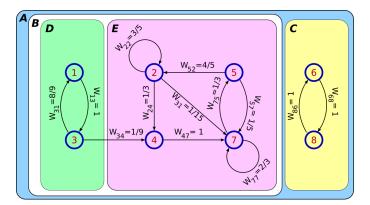


Figure 2.4 – A transition graph: each numbered circle represents a state and each directional link between states has a probability w that is the probability that the system passes from a given state to another (following the arrow direction). The transition graph A can be split in 2 independent Markov chains, B and C. Moreover the chain D can be reduced to the chain E because no transitions are possible from a state of E to a state of D. Markov chain E is instead irreducible (adapted from [20]).

shows a graph in which each node is a *state* and the edges link vertices with non-zero transition probability. For instance, if we are in state 2 at a given time t, the configuration at time t + 1 could be state 4, with probability $w_{24} = \frac{1}{3}$, or state 7, with probability $w_{27} = \frac{1}{15}$, or to stay in state 2, with probability $w_{22} = \frac{3}{5}$. From properties of the probability we assume that

$$\sum_{j} w_{ij} = 1 \tag{2.2.7}$$

and the probability to be in a given state i at a given time t + 1 is

$$p_i(t+1) = \sum_j p_j(t) w_{ji},$$
(2.2.8)

where $p_j(t)$ is the probability to be in the state, j, at the previous time, t; w_{ji} represents the probability to evolve from state j to state i in the last time step. From (2.2.7) and (2.2.8) we obtain again (2.2.5) in the form

$$p_i(t+1) - p_i(t) = \sum_j p_j(t)w_{ji} - \sum_j p_i(t)w_{ij}$$
(2.2.9)

and from this form the compact matrix notation used in (2.2.6) follows.

2.2.3 Metropolis algorithm

In this section I review the powerful generic method, called Markov chain Monte Carlo (MCMC), for approximately generating samples from an arbitrary distribution. This is typically not an easy task, in particular when X is a random vector with dependent components. The MCMC method is due to Metropolis et al. [21]. They were motivated by computational problems in statistical physics, and their approach uses the idea of generating a Markov chain whose limiting distribution is equal to the desired target distribution. There are many modifications and enhancement of the original Metropolis algorithm, most notably the one by Hastings [22]. Nowadays, any approach that produces an ergodic Markov chain whose stationary distribution is the target distribution is referred to as MCMC or Markov chain sampling [23]. The most prominent MCMC algorithms are the Metropolis-Hastings and the Gibbs samplers, the latter being particularly useful in Bayesian analysis. Finally, MCMC sampling is the main ingredient in the popular simulated annealing technique [24] for discrete and continuous optimization.

The main idea behind the Metropolis-Hastings algorithm is to simulate a Markov chain such that the stationary distribution of this chain coincides with the target distribution. To motivate the MCMC method, assume that we want to generate a random variable X taking values in $\mathscr{X} = 1, \ldots, m$, according to a target distribution $\{\pi_i\}$, with

$$\pi_i = \frac{b_i}{C}, \quad i \in \mathscr{X}, \tag{2.2.10}$$

where it is assumed that all b_i are strictly positive, m is large, and the normalization constant $C = \sum_{i=1}^{m} b_i$ is difficult to calculate. Following Metropolis et al. [21], we construct a Markov chain $\{X_t, t = 0, 1, ...\}$ on X whose evolution relies on an arbitrary transition matrix $\mathbf{Q} = (q_{ij})$ in the following way:

- When $X_t = i$, generate a random variable Y satisfying $P(Y = j) = q_{ij}, j \in \mathscr{X}$. Thus, Y is generated from the *m*-point distribution given by the *i*-th row of Q.
- If Y = j, let

$$X_{t+1} = \begin{cases} j & \text{with probability } \alpha_{ij} = \min\{\frac{\pi_j q_{ji}}{\pi_i q_{ij}}, 1\} = \min\{\frac{b_j q_{ji}}{b_i q_{ij}}, 1\} \\ i & \text{with probability } 1 - \alpha_{ij} \end{cases}$$
(2.2.11)

It follows that $\{X_t, t = 0, 1, ...\}$ has a one-step transition matrix $\mathbf{P} = (p_{ij})$ given by

$$p_{ij} = \begin{cases} q_{ij}\alpha_{ij} & \text{if } i \neq j \\ 1 - \sum_{k \neq i} q_{ik}\alpha_{ik} & \text{if } i = j \end{cases}$$
(2.2.12)

Now it is easy to check that, with α_{ij} as above,

$$\pi_i p_{ij} = \pi_j p_{ji}, \ i, j \in \mathscr{X}$$

In other words, (2.2.12) satisfies the detailed balance, and hence the Markov chain is time reversible and has stationary probabilities $\{\pi_i\}$. Moreover, this stationary distribution is also the limiting distribution if the Markov chain is irreducible and aperiodic. Note that there is no need for the normalization constant C in (2.2.10) to define the Markov chain. The extension of the above MCMC approach for generating samples from an arbitrary multidimensional probability density function f(x) (instead of n,) is straightforward. In this case, the nonnegative probability transition function q(x, y) (taking the place of q_{ij} above) is often called the proposal or instrumental function. Viewing this function as a conditional probability density function one also writes q(y|x) instead of q(x, y). The probability $\alpha(x, y)$ is called the acceptance probability. The original Metropolis algorithm [21] was suggested for symmetric proposal functions, that is, for q(x, y) = q(y, x). Hastings modified the original MCMC algorithm to allow non-symmetric proposal functions. Such an algorithm is called a Metropolis-Hastings algorithm. We call the corresponding Markov chain the Metropolis-Hastings Markov chain. In summary, the Metropolis-Hastings algorithm, which, like the acceptancerejection method, is based on a trial-and-error strategy, is comprised of the following iterative steps.

Metropolis-Hastings Algorithm

Given the current state \mathbf{X}_t :

- 1. Generate $\mathbf{Y} \sim q(\mathbf{X}_t, y)$.
- 2. Generate $U \sim U(0,1)$ and deliver

$$\mathbf{X}_{t+1} = \begin{cases} \mathbf{Y} & \text{if } U \leqslant \alpha(\mathbf{X}_t, \mathbf{Y}) \\ \mathbf{X}_t & \text{otherwise,} \end{cases}$$
(2.2.13)

where

$$\alpha(x,y) = \min\{\rho(x,y),1\} \quad with \quad \rho(x,y) = \frac{f(y)q(y,x)}{f(x)q(x,y)}.$$
(2.2.14)

By repeating Steps 1 and 2, we obtain a sequence $\mathbf{X}_1, \mathbf{X}_2 \dots$ of dependent random variables, with \mathbf{X}_t approximately distributed according to f(x), for large t.

Since this algorithm is of the acceptance-rejection type, its efficiency depends on the acceptance probability $\alpha(x, y)$. Ideally, one would like q(x, y) to reproduce the desired probability distribution function f(y) as faithfully as possible. This clearly implies maximization of $\alpha(x, y)$. A common approach [23] is to first parameterize q(x, y) as $q(x, y; \theta)$ and then use stochastic optimization methods to maximize this with respect to θ .

2.2.4 How to use the Metropolis algorithm for spin systems

The Metropolis algorithm consists of a set of rules derived from the energy of a given configuration on a 3D EA lattice of side L.

Let $\{s_i\}$ with $i = 1, ..., L^3$ a configuration of a three-dimensional lattice of spins (indicated with s as usual). The energy of the configuration $\{s_i\}$ is given by

$$E(\{s_i\}) = -\sum_{\langle i,j \rangle} s_i J_{ij} s_j \tag{2.2.15}$$

We consider now a spin s_i , we propose a change for it s'_i (remember that in EA model variables have only two values $s \in \{+1, -1\}$) and we calculate the energy of the configuration $\{s'_i\}$

$$E(\{s'_i\}) = -\sum_{\langle i,j \rangle} s'_i J_{ij} s_j \qquad (2.2.16)$$

If the energy change $\Delta E = E(\{s'_i\}) - E(\{s_i\})$ is negative the new state is automatically accepted, while if the energy value grows the new state is accepted only with a certain probability. In practice, the probability of accepting a configuration change is:

$$P(s_i \to s'_i) = \begin{cases} 1 & \text{if } E(\{s'_i\}) < E(\{s_i\}) \\ e^{-\beta \Delta E} & \text{if } E(\{s'_i\}) > E(\{s_i\}) \end{cases}$$
(2.2.17)

The probability $P(s_i \to s'_i)$ from Eq. (2.2.17) is then compared with a random value r in order to decide whether to accept or reject the proposed state change s'. This ends a spin update; once all L^3 have been updated we say that a Monte Carlo sweep (MCS) has completed. Is important to note that the site to be updated can be chosen randomly or following a given order and the new value obtained after the update is used for the following updates (i.e. it is not necessary to retrieve a copy of whole old system until the end of a MCS).

2.2.5 Another MC algorithm: the heat bath

In the case of the heat bath algorithm we directly select the new value of the spin with a probability proportional to the Boltzmann factor and regardless of the value of the spin lying in the site we are updating.

The probabilities to assign to a given spin s_i the new value +1 or -1 is defined as

$$P_{\rm HB}(s_i = +1) = \frac{e^{-\beta E_+}}{e^{-\beta E_+} + e^{-\beta E_-}}, \qquad (2.2.18)$$

$$P_{\rm HB}(s_i = -1) = 1 - P_{HB}(s_i = +1)$$
(2.2.19)

where E_+ and E_- are the local energies of the two spin configurations for spin s_i pointing up $(s_i = +1)$ or down $(s_i = -1)$, respectively.

In detail, we consider a site i of the lattice and we calculate its local energy as the sum of the values of its nearest neighbors (weighted by the related couplings J)

$$E_i = \sum_{\langle j \rangle} s_j J_{ij}$$

Using E_i we calculate the probability $P_{\text{HB}}(s_i)$ as in 2.2.18 and compare it with a random number R. If $R \leq P_{\text{HB}}(s_i)$ then we assign +1 as new value of the spin, else the new value will be -1.

As the calculation of the new value of a spin s_i does not depend on the current value of s_i , a few terms of the energy function change and computational load is lower.

This method is most useful in circumstances where the Metropolis-like approach described in Section 2.2.4 has a very low acceptance rate, i.e. when a new spin proposed in the Metropolis algorithm is accepted with very low probability. The effect is that using Heat Bath some systems reach a particular state of equilibrium, called thermalization, with fewer Monte Carlo steps than with Metropolis.

2.2.6 Parallel tempering techniques

The free energy landscapes of complex systems are characterized by many local minima which are separated from each other by energy barriers. In studying such systems, we have to take into account of each configuration for these local minima and the fluctuation around it. The characteristic time in which the system escapes from a local minimum, however, increases rapidly as temperature decreases. This situation causes "hardly-relaxing" problem in using conventional Monte Carlo simulations based on a local updating (see Figure 2.5 for a graphical representation of this situation).

Various new algorithms have been proposed in the last years to overcome this problem (see [25] for a review). In 1996 Hukushima and Nemoto [26] proposed a method called *parallel tempering*, previously called also *replica exchange* [27, 28], in which M non-interacting replicas of the system are simulated simultaneously at a range of temperatures $\{T_1, T_2, \ldots, T_M\}$. After a fixed number of Monte Carlo sweeps a sequence of swap moves, the exchange of two replicas at neighbouring temperatures, T_i and T_{i+1} , is suggested and accepted with a probability

$$p(E_i, T_i \to E_{i+1}, T_{i+1}) = \min\{1, \exp(\Delta\beta\Delta E)\}$$

$$(2.2.20)$$

where $\Delta\beta = 1/T_{i+1} - 1/T_i$ is the difference between the inverse temperatures and $\Delta E = E_{i+1} - E_i$ is the difference in energy of the two replicas. At a given temperature, an accepted swap move effects a global update as the current configuration of the system is exchanged with a replica from a nearby temperature. For a given replica, the swap moves induce a random walk in temperature space. This random walk allows the replica to overcome free energy barriers by wandering to high temperatures where equilibration is rapid and returning to low temperatures where relaxation times can be long. The simulated system can thereby efficiently explore complex energy landscapes that can be found in frustrated spin systems and spin glasses.

Recently Katzgraber et al. in [29] proposed a way to maximize the efficiency of parallel tempering Monte Carlo by optimizing the distribution of temperature points

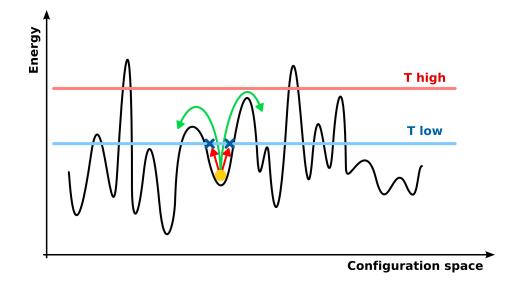


Figure 2.5 – Energy landscape of a complex system such as spin glass. The configuration A (circle) is trapped in a minimum. A low temperature makes transitions between configurations with large energy differences unlikely; by moving the configuration to a higher temperature we increase the probability to escape from the local minimum.

in the simulated temperature set such that round-trip rates of replicas between the two extremal temperatures in the simulated temperature set (i.e. T_1 and T_M) are maximized. The optimized temperature sets are determined by an iterative feedback algorithm that is closely related to an adaptive algorithm, introduced in [30], that explores entropic barriers by sampling a broad histogram in a reaction coordinate and iteratively optimizes the simulated statistical ensemble defined in the reaction coordinate to speed up equilibration.

Katzgraber's work allows one de facto to automatize the process of temperature swap needed by parallel tempering method so that temperature points are sampled near the bottleneck of a simulation.

2.3 Numerical requirements

From previous sections the theoretical infrastructure that holds the simulations implemented on the Janus supercomputer should be clear. In this section I introduce the computing requirements of them and we will see in detail the reasons leading us to implement some algorithms of statistical physics such as the EA model and Potts model, directly in hardware, using the reconfigurable computing approach explained in Chapter 1.

Special attention will be paid to the generation of random numbers and its implementation in hardware because the Monte Carlo simulations are strongly dependent on it: a typical run requires in fact $\sim 10^{10} \cdot L^3$ random numbers.

```
k < SIDE
for
        ( k=0
                                       <u>k</u>++
                                                   ł
                           j<SIDE; j++ )
; i<SIDE; i++
   <u>for</u> ( j=Ó
       <u>for</u> ( i=0
                                                     )
                                                           {
                                        the index
                calculating
                                                           0 f
                                                                 the
                                                                         LUT
                                 [(i+1)%SIDE][
[(i-1)%SIDE][
                                                                          Jx[(i+1)%SIDE]
Jx[(i-1)%SIDE]
                                                        Гj
                                                            ][k]
][k]
                       spin
           idx =
                       spin
                                                                                                                   +
                                                                      * Jy[i][(j+1)%SIDE][k]
* JY[i][(j-1)%SIDE][k]
* Jz[i][j][(k+1)%SIDE]
* Jz[i][j][(k-1)%SIDE]
                                         (j+1)%SIDE][k]
                       spin
                                 [i]
                                 [i][(j-1)%SIDE][k]
[i][j][(k+1)%SIDE]
[i][j][(k-1)%SIDE]
                                                                                                                   +
                       spin
                       spin
                                                                                                                   +
                       spin
                 comparing a random number with the value
of the HBT look up table of index 'idx'
(rand() < HBT(idx)) {
           if
                         [i][j][k]
                spin
                                                +1 ;
               <u>else</u>
                         [i][j][k] = -1;
               spin
           }
       }
   }
}
```

Code 2.1 – Heat Bath C code

2.3.1 Implementation and available parallelism

In previous sections I presented the theoretical models in which part of the physics community is interested and I gave an overview of the "tools" used to study them, but we still have the question How?: how do we write a Monte Carlo simulation for spin lattice? In frame Code 2.1 I present a non-optimized C code of a simulation to run the Heat Bath algorithm.

Inspection of Code 2.1 shows some features that are shared with other spin models implementation (such as Edward-Anderson and Potts):

- The code kernel has a regular structure, associated with regular loops. At each iteration the same set of operations is performed on data stored at locations whose addresses can be predicted in advance.
- Types of variables are "unusual": 3D array called *spin* and J are array of integers, but the information stored in each variable could be coded in only one (or just a few) bits. The physical values $\{+1, -1\}$ can be easily mapped with no loss of generality on to the set $\{1, 0\}$. Moreover data processing is associated with logical operations performed on bit-variables denoting the site variables.
- If we consider a lattice of L³ spins, the data base associated with the computation is: L³ bits to store spin informations; 3 ⋅ L³ bits to store J informations. Note that couplings are symmetric, so we do not need to store them for all directions x₊, x₋, y₊, y₋, z₊, z₋: positive directions x₊, y₊, z₊ are enough. This means that a typical data base associated with the main computational core of a lattice with side L = 64 is ~ 1 Mbit. If we consider the size of the data base, it is important to remember that a huge quantity of random numbers are needed (see Section 2.3.3 for more details).

If we think of the physical spin variables in set {+1, −1} we can easily calculate the possible value of the local energy for a given spin s_i, ∑_{⟨i,j⟩} J_{ij}s_j. Assuming that J ∈ {+1, −1} the possible values for the local energy in the EA model are only 7: {−6, −4, −2, 0, +2, +4, +6}. This allows a first trivial optimization in order to avoid the calculation of the exponential, e^{-βΔE}; we just read its value from a look up table (LUT) with 7 entries that can be initialized before the start of the simulation. It is clear that the use of a LUT is useful when the number of its entries is not too big: in the case of a Gaussian Potts model, in which J ∈ ℝ it is impossible to implement it and the calculation of the exponential is required.

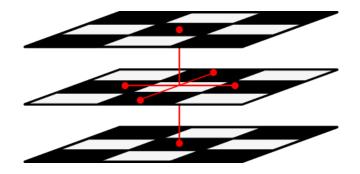


Figure 2.6 – Checkerboard organization of a piece of spin lattice. When spins sitting on white sites are updating, the black ones (nearest neighbours belonging to the 3 directions) should have a fixed value.

What is not explicitly written in the sample Code 2.1 (but easily checked by inspection) is that a very large amount of parallelism is available: each site interacts with its nearest neighbors only, while the correctness of the procedure requires that when one site is being updated its neighbors are kept fixed. As a consequence, up to one-half of all the sites, organized in a checkerboard structure as in Figure 2.6, can in principle be operated upon in parallel. The same sequence of operation is performed while updating any site, so almost perfect load-balancing is also possible.

2.3.2 Techniques on a general purpose processor

As Monte Carlo algorithms works on discrete lattice, there are some tricks commonly used to speed up performances and to parallelize them on general purpose architectures.

A basic idea useful on clusters is to split the lattice into sub-lattices that can be independently updated. This method is well known and has as its main problem the fact that nodes lying on boundaries have neighbors processed by a different CPU. Shared memory or exchange of the "halo" nodes are usually solutions for these kinds of problems. In the case of the spin glass simulation we saw e.g. in Figure 2.6 that up to $L^3/2$ spins can be updated concurrently, thus a high level of parallelism is available also while the update process is running. Today multi-core and SIMD processors are widely available and, as they can support a high level of parallelism, it seems interesting to try to port to them using the Metropolis algorithm described above. The technique used to reach this goal is the *multispin coding* [31, 32, 33, 34] which stores the states of several variables (spins in the case of the spin models) in the bits of a single word of data. So the Monte Carlo algorithm is written in terms of operations which act on an entire word at a time. The result is that it is possible to perform operations on several variables at once, rather than performing them sequentially, which improves the speed of our calculation, thereby giving us better results for a given investment in CPU time.

We will consider the recently introduced IBM CBE and Intel Nehalem architecture as representative examples of a class of multi-core processors with the following features:

- a small set of cores are integrated on a single processor.
- Each core support SIMD-like vector instructions operating on W-bit vector datawords. Each vector data-word can be partitioned in two or more scalar data-words of different sizes.
- each core has access to a local private memory.
- each core can access data stored on the private memories of the other cores.
- all cores share (arbitrarily) large main memory external to the chip.
- memory accesses and data-transfers can be performed concurrently with computation.
- memory access performance has a strong dependence on "distance": latency to the local memory is constant and small. Access to the local memories of the other cores has a significantly longer latency, while access to the shared memory is still much slower.

It is possible to evaluate the performance of a program in terms of *spin update* time: the time required to update a single spin of a given system. It is equal to the wall clock time for a Monte Carlo step divided by L^3 .

The goal is to exploit the highest possible level of parallelism, so using the highest granularity of the architecture allows one to update the highest number of spins in parallel. We need in principle only one bit of the smallest available scalar variable, the spin, (or few bits in the case of the Potts model), so a large amount of memory is wasted. For example, if the architecture allows the use of V bytes in parallel, then $7 \times V$ bits are not used, and more generally, if scalar variables are w bit long, then $(w-1) \times V$ bits are wasted.

Taking a different approach, it is possible to use a scalar register of w bits to represent corresponding spins of w different systems. This technique, called *asynchronous multi-spin coding* allows several systems to evolve concurrently. It has no effect in shortening the wall-clock time for the simulation of a single system (it usually even worsen it!) but is the most efficient approach when large statistics is needed. Multispin coding is possible because the algorithm can be coded in such a way that all the operations (that in principle are sums, multiplications and comparisons between integer numbers) can be mapped to logical bit-wise instructions with a very small overhead.

On the other hand *synchronous multi-spin coding* is a similar approach in which the bits in a word represent several spins (or other variables) on the lattice of a single system so that one multi-spin coded Monte Carlo step updates many sites at once. This approach has the bottleneck of random numbers that can be generated in parallel: in average each update process requires, in fact, an independent random number.

Combinations of SMSC and AMSC allow to perform parallel updates of spins of the same system (with the upper limit of the random numbers generation) and, at the same time, of spins belonging to different systems, sharing the random numbers and increasing the statistics. Figure 2.7 shows an example of a multi-spin coding implementation in which a 128 bit vector is used to perform updates in parallel of 32 systems (AMSC) and for each of them 4 spins are updated in parallel (because we assume that the multi-core processor considered for this example allows us to generate 4 random numbers in parallel).

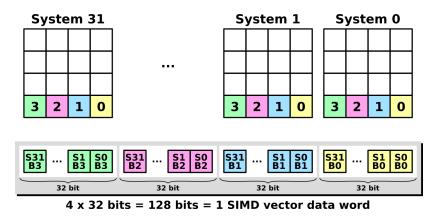


Figure 2.7 – Coding of the spins of 32 systems in a vector data word of 128 bit allowing the parallel updates of 4 set of 32 spins.

2.3.3 Random numbers

The efficiency of Monte Carlo methods depends largely on the random numbers used to drive the updates: this determines the imperative need to implement a very reliable pseudo-random number generator (RNG), that produces a sequence of numbers under

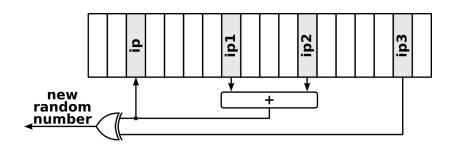


Figure 2.8 – Generic scheme of the Parisi-Rapuano RNG. In out implementation we set ip1 = 24, ip2 = 55 and ip3 = 61.

the selected distribution, with no known or evident pathologies. We use the Parisi-Rapuano shift register method [35] defined by the rules:

$$I(k) = I(k - 24) + I(k - 55)$$

$$R(k) = I(k)I(k - 61),$$
(2.3.1)

where I(k-24), I(k-55) and I(k-61) are 32-bit words of a set, called the wheel that we initialize with externally generated random values. The index k is incremented at every step, though its value is defined modulo the size of the wheel. I(k) is the new element of the updated wheel, and R(k) is the generated pseudo-random value. An implementation of this algorithm is shown in Figure 2.8.

The exact period and other features of this random generator are not known; we only check that no periodicity appears generating 10^{13} numbers. If we consider that Janus produce one random number per clock cycle (note that we produce up to 1024 random numbers but with independents generators) we can estimate a lower limit time of the random number goodness:

$$10^{13} \text{ numbers} \times 16 \text{ ns} \approx 44 \text{ hours}$$
 (2.3.2)

to avoid periodicity problems we periodically (every $\sim 10^6$ Monte Carlo updates of whole lattice equivalent to less than 3 hours) refresh the wheel. Moreover we monitor physical quantities that are extremely sensitive to random number quality: they have unpredictable behavior in presence of problems connected with period or correlation so they are perfect candidates to check the goodness of random numbers. Figure 5.1(b) for instance shows the graph of the correlation at different temperature after 10^{19} Monte Carlo steps: in presence of correlation phenomena the correlation function does not reach a plateaux or presents oscillations that we do not find in our studies.

The easy rule described in (2.3.1) implemented via standard programming language in a function can be called every time we need a random number. However access to the wheel I has a performance cost but the presence in the new architecture of long registers (up to 128 bits) allows us to generate $4 \cdot 32$ -bit random numbers in parallel. From a hardware implementation point of view a wheel uses many hardware resources (in our case where we use the three pointer values 24, 55 and 61 we need to store 62 numbers), and the random number generator is a system bottleneck, since the number of updates per clock cycle is determined by how many random values we are able to produce. A large performance improvement comes from the implementation of the wheel through logic elements (as opposed to memory blocks), as the former can be written in cascade structured combinatorial logic that may be tuned to produce several numbers per clock cycle. We can exploit this feature and use a limited number of wheels to produce more random numbers (and therefore more updates) per clock. Remember that to produce one random number we must save the result of the sum of two values and then perform the XOR with a third value. The wheel is then shifted and the computed sum fills the empty position. All this is done with combinatorial logic, so one can produce various pseudo-random numbers simply replicating these operations and, of course, increasing logic complexity. A schematic representation of a simplified case is given in Figure 2.9.

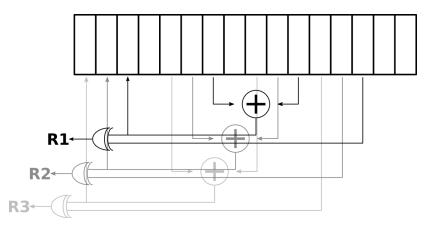


Figure 2.9 – Parallel implementation of the random number generator. For graphical reason the example shows a wheel generating 3 random numbers in parallel. Duplication and shift of the pointers allow parallel generation of numbers, but logic complexity (and use of hardware resources) grows when producing more numbers.

The logical complexity of the implementation depends on the parameters of (2.3.1)and on the quantity of random numbers that we need. For example to perform Nupdates per clock cycle with the EA model we need N random numbers, while $2 \cdot N$ random numbers are needed for the same number of updates of a Potts model. Such a large number of RNG wheels would saturate easily the FPGA resources, and we would have to reduce the number N of parallel updates. To avoid this problem we have implemented the optimized version discussed above, so that each wheel can produce many random numbers in each cycle. At present we use one wheel to generate up to 96 numbers per clock (so even this way we need more wheels to be active at the same time, in order to compute all the random values required by the algorithm). With respect to the choice of 32-bit random numbers, we have verified that this word size is sufficient for the models that we want to simulate. Indeed we compare random numbers with transition probabilities calculated as $U = e^{-\beta \Delta E_i}$ where β is the inverse of the temperature and $\Delta E_i = E_{s'_i} - E_{s_i}$ energy fluctuation after an update step calculated as in (2.2.15). Typical values assumed by β are ~ 1 (see values in Section 5.2). We can therefore calculate the smallest value U corresponding to the biggest value of energy variation ΔE :

$$\Delta E = 12 \ \to \ e^{-12} \approx 6 \cdot 10^{-6} \tag{2.3.3}$$

The number of bits to write the integer representation of ΔE is

$$N_b(\Delta E) = \log_2 6 \cdot 10^6 \approx 22$$
 bits

If we consider an error of 10% the random number should be able to represent a number roughly hundred times smaller than the smallest energy difference, so it needs

$$N_b(R) = \log_2 6 \cdot 10^8 \approx 29$$
 bits.

Mainly we use on Janus 32 bits to be compliant with standard computer systems.

Other models may require larger (i.e. with more bits) random numbers. Generating random numbers of larger size (e.g., 40 or even 64-bit) is rather straightforward, at the price, of course, of a larger resource usage, which implies a reduced number of parallel updates. Obviously the same problem would affect PC simulations as well. Most of the models studied so far have been simulated with 32-bit random numbers.

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We gotta get out while we're young 'cause tramps like us, baby we were born to run.

Bruce Springsteen

3 Janus architecture at large

In this chapter I introduce the Janus project and the FPGA-based system built at the heart of it. A short introductive section is dedicated to explain the structure of the system, the main components and their functions.

After highlighting the Janus system the idea is to retrace the architectural questions that led to the Janus systems and were solved by me and my group in Ferrara during prototype and development phase of Janus (the main activity of my PhD). Therefore a relatively long part of this chapter will be dedicated to a detailed overview of the problems connected to the development of a massively parallel dedicated machine such as Janus.

The last three sections are dedicated to the hardware implementation of two statistical physics models: the Edward Anderson spin glass model, the spin model using the parallel tempering technique and the Potts model on an irregular lattice, the last one with the main aim to solve the graph coloring problem. The theoretical details of these problems were introduced in previous chapter, therefore they are here presented in terms of their implementation in a hardware description language on a programmable device.

Details about input output, control system, firmware, VHDL code etc. are postponed to the next chapter in order to keep here general ideas strongly connected with the needs that encourage an heterogeneous group of physicists, mathematicians and computer scientist to build a dedicated machine not only in order to solve open problems of statistical physics but also to study new architectural solutions and new algorithms.

3.1 Janus project

The Janus project is a collaboration between BIFI (Institute for Biocomputation and Physics of Complex Systems), universities of Spain (Badajoz, Madrid and Zaragoza) and Italy (Ferrara and Roma 1) with the industrial partnership of Eurotech Group. The main aim of the project is to build an FPGA based supercomputer strongly oriented to study and solve the problems related with the spin systems introduced in Chapter 2.



Figure 3.1 – The Janus rack during the installation for the unveiling in Zaragoza in Spring 2008.

The Janus supercomputer is a modular system composed of several Janus boards. Each board houses 17 FPGA-based subsystems: 16 so-called scientific processors (SPs) and one input/output processor (IOP). Janus boards are driven by a PC (Janus host). The first system, deployed in December 2007, is composed by 16 boards and 8 Janus hosts and is depicted in Figure 3.1 during the first installation and the unveiling in Zaragoza.

The 17 FPGA chips (generically referred to as *nodes*) are housed in small daughter cards plugged into a mother board so that the 16 SPs form a 2D nearest neighbour grid with toroidal enclosure as shown in Figure 3.2a. Figure 3.2b exhibits a board housed in a Janus case with appropriate cooling systems and cables.

The advantage of using removable modules as opposed to a direct connection of the FPGA to the main boards, is that it is easy to substitute single nodes in case of damage. This choice also facilitates a basic idea leading the Janus project: chips are strongly dependent to the technology progress, so could be possible to upgrade Janus

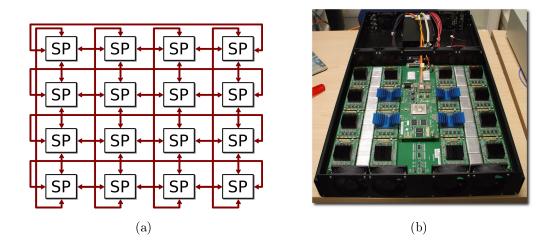


Figure 3.2 – (a) Topology of a Janus board: each SP communicates with its nearest neighbours in the plane of the board. (b) Janus board housed in a Janus box.

system with relative small effort.

Each SP is connected with the nearest neighbour SP and with IOP via hardwired dedicated lines creating two networks: a point to point network between IOP and SPs used for initializations and controls and a 2D toroidal mesh in order to allow trivial parallelization among SPs.

The host PC play a key role of master device: a set of purpose-made C libraries are written using low levels of linux operating system in order to access the raw Gigabit ethernet level (excluding protocols and other unhelpful layers adding latencies to communications). Moreover two software environments are available: an interactive shell written in Perl mostly used for testing and debugging or short preliminary run and a set of C libraries strongly oriented to the physicists making relative easy to access the hardware resources of Janus to peoples with no hardware experience.

3.2 Questions leading Janus's development

The main aim of this section is to give a comprehensive overview of the problems and challenges that were widely discussed and solved during the 2 years of Janus's development.

The section is deliberately structured as a set of questions with their respective answers, mimicking the situations created during internal discussions and collaboration in Janus's meetings and during poster sessions or conferences talks.

3.2.1 Why many nodes on a board?

There are three primary classes of techniques for creating a parallel algorithm to perform a simulation on a *d*-dimensional lattice: trivial parallelization, functional decomposition and domain decomposition.

Trivially parallel problems are ones which can be split up into separate tasks which are completely unrelated and can be performed by separate programs running on isolated processors without the need for inter-processor communication.

For more complicated problems we have two ways to go. Functional decomposition refers to the breaking up of a task into a variety of different jobs: each processor in a program which carries out such a task is performing a different type of work, and the contributions of all the processors put together perform the complete task. We could imagine for example delegating separate processors in a Monte Carlo calculation to evaluation of a Hamiltonian or other observable, construction of clusters, updating lattice variables, or any number of other elementary chores. This however is not a very common approach to the kinds of calculations needed in parallel systems such as Janus.

Monte Carlo simulations almost always use domain decomposition, which means breaking up the calculation into separate simulations of different regions of the system under study. Since Monte Carlo simulations are usually concerned with systems defined on a lattice this approach is often quite straightforward to implement. The processors in such a calculation all perform essentially the same task on different parts of the lattice, and indeed are often running exactly the same program.

Following this general idea, Janus is structured as a collection of FPGA-based processors, running the same firmware, able to process each one a portion of a given lattice.

3.2.2 Why an Input/Output processor?

In computer science is widely use the concept of *loose coupling* referring to a relationship in which one module interacts with another module through a stable interface and does not need to be concerned with internal implementation of other modules. In this definition the term *module* is sufficiently general to encompass functions of a high level language, or hardware components such as a CPU or an FPGA. Conversely when interaction of two modules requires informations about their implementation, we are in presence of *tight coupling*.

Systems that do not exhibit tight coupling might experience the following developmental difficulties:

- Change in one module forces a ripple of changes in other modules.
- Modules are difficult to understand in isolation.

- Modules are difficult to reuse or test because dependent modules must be included.

The basic idea for Janus was to develop a loosely coupled machine in which the Janus host (intended as a module) should interact with SPs while ignoring in principle the algorithm implemented within the FPGA. This statement was a guide in the development of the Janus hardware and for this reason we chose to develop a custom Input/Output Processor, *IOP*, to play the key role of a stable interface between the Janus host and SPs, thus assuring the loose coupling between them.

It is clear that the development of an interface module is not the only way to efficiently realize a loosely coupled machine: e.g. Maxwell uses the PCI buses to allow communications between the general purpose processors and the FPGAs constituting the system.

The choice of a custom module such as IOP comes from features of the class of algorithms that we plan to run on Janus: Monte Carlo simulations for spin lattices requires in fact a small data exchange with relative low latency and very long run time. If we consider, for instance, a spin lattice of side L = 64, the data to transfer from the Janus host to the SP in order to start a simulations is: L^3 bits for the spins, $3 \cdot L^3$ bits for the Js, 7 words of 32 bits in order to initialize the LUT of the energies and 61 words of 32 bits for initializing of the random number generator.

$$64^3 + (3 \cdot 64^3) + (7 \cdot 32) + (61 \cdot 32) \approx 1$$
 Mbits

If we consider the worst case in which each SP run a completely different history to the others, we have to transfer ~ 16 Mbit (2 MB) from the Janus host to each of the FPGA structures storing data. Another Janus requirement could be to upload new firmware for the SPs: this requires a transfer of ~ 5 MB.

So we have that typical bandwidth requirements are of the order of some MB, therefore relatively small, but what are the latency requirements? Some Monte Carlo techniques requires that every P Monte Carlo steps spin configurations (or alternatively LUT) are changed using a simple rule depending on the temperature (parallel tempering). In such cases is very important to have a large bandwidth, but much more important is a low latency in order to exchange efficiently the configuration and so not waste computing time.

3.2.3 How are organized communications between Janus boards and Janus host?

The programming framework developed for Janus is intended to meet the requirements of the prevailing operating modes of Janus, i.e. supporting (re)configuration of SPs,

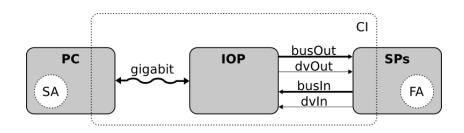


Figure 3.3 – Framework of an application running on the Janus system.

initialization of memories and data structures within the FPGA, a monitor system during run, memory interface and some other debug functions.

Applications running on the Janus system can be thought of as split into two subapplications, one, called *software application* SA, written for example in C, and running on the Janus host. The other, called *firmware application* FA, written for example in VHDL, and running on the nodes of a Janus board. As shown in Figure 3.3, the two entities, SA and FA are connected together by a *communication infrastructure* CI, that is a logical block including physically the IOP and which allows the exchange of data and performs synchronization operations, directly and in a transparent way.

The CI abstracts the low-level communication between SA and FA applications, implemented in hardware by the IOP and its interfaces. It includes three main components:

- a C communication library linked by the SA,
- a communication firmware running on the IOP processor, interfacing both the host PC and the SP processor,
- a VHDL library linked by the FA.

Firmware running on the IOP communicates with the host PC via a dual Gigabit channel, using the standard RAW-ethernet communication protocol. The choice of RAW-ethernet is based on the following considerations:

- the standard TCP/IP protocol provides a lot of functionalities which are not needed by our purposes;
- an hardware implementation of the TCP/IP stack is non trivial, uses far more hardware resources and is more time critical than a RAW-ethernet implementation;
- the TCP/IP protocol introduces overheads in the communication.

To guarantee reliability of the communication in the direction IOP to Janus host, we adopt the *Go-back-N* protocol [1]. In this protocol, N frames are sent without waiting for an acknowledge from the receiver. After N frames have been transmitted, the sender waits for an acknowledge. If the acknowledge is positive it starts to send the next N frames, otherwise it send again the last N frames. The receiver waits for N frames, and after receiving all the frames sends an acknowledge to the sender. If some frames are lost, typically dropped by the network card because of CRC error, or by the Linux operating system because of network buffer overflow, the receiver gets a timeout and send back a not-acknowledge, asking re-transmission of the last N frames. Using the Go-back-N protocol we reach approximately the 90% of the full Gigabit bandwidth, transferring messages whose length is of the order of 1 MB, using the maximum data payload per frame, 1500 bytes, and N = 64, see Figure 3.4.

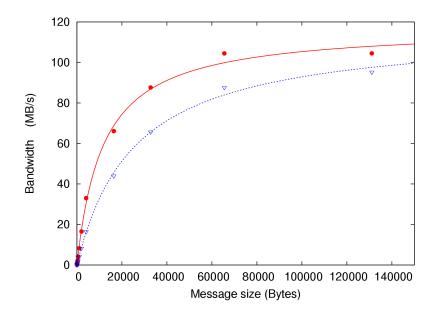


Figure 3.4 – Measured transfer bandwidth of the IOP processor. Red bullets are for write operations (Janus host to nodes), blue triangles are for read operations (nodes to Janus host). The lines are fits to a simple functional form $B(s) = \frac{s}{\alpha + \beta s}$ where s is the message size.

In the other direction we did not adopt any communication protocol since the IOP interface, barring hardware errors, ensures that no packets are lost. Incoming frames are protected by standard ethernet CRC code, and errors are flagged by the IOP processor.

Data coming from the SA application are interpreted by the IOP as commands for itself, to set, for instance, internal control registers, or are routed to one or more SPs according to specific control information packed together with the data as described in 3.2.4. Data coming from the FA application are packed as burst of N frames, according to the Go-back-N protocol.

Developers of applications running on the Janus system have to write the SA and FA relaying on the CI communication infrastructure, to make the two applications collaborative. A typical SA application configures, using the functions provided by the communication library, the SP processors with the firmware corresponding to the FA program, loads input data for the FA application, and it waits for the incoming results. On the other side, a typical FA application wait for incoming requests, performs some calculation and sends back the results.

The approach adopted to develop application for the Janus system explained in this section is to keep deliberately the structure of the IOP as easy as possible in order to have a low latency stable (and hopefully standard) interface between SA running on Janus host and FA running on SPs. This approach is coherent to the concept of loose coupling introduced in 3.2.2 and could allow the use of development toolkits to automatically decompose an application into cooperating modules.

3.2.4 How are organized communications within a Janus board?

The current firmware configuration of the IOP focuses on the implementation of an interface between SPs and Janus host. Although a programmable device is used to implement it, IOP is not intended as a reconfigurable processor: it basically allows streaming of data from the host to the appropriate destination (and back), under the complete control of the Janus host and its configuration should remain fixed.

The IOP configured structure is naturally split into 2 worlds:

- IOlink block handles the I/O interfaces between IOP and host-PC (Gigabit channels, serial and USB ports). Its supports the lower layers of the Gigabit ethernet protocol, ensures data integrity and provides a bandwidth close to the theoretical limit;
- *multidev block* contains a logical device associated with each hardware subsystem that may be reached for control and/or data transfer: there is a memory interface for the staging memory, a programming interface to configure the SPs, an SP interface to handle communication with the SPs (after they are configured) and several service/debug interfaces.

A block diagram of the IOP is shown in Figure 3.5.

Gigabit ethernet feeds the IOlink with a stream composed of 32 bit words with a clock period of 8 ns (125 MHz). Internal entities of the IOlink synchronize the stream and return a 16 bit data stream (dataIn) and a data valid (dvIn) clocked by the main system clock with a period of 16 ns (62.5 MHz). See section 3.2.9 for more detail about clock frequencies.

In order to route the input stream coming from IOlink we considered three different possibilities explained in detail in Appendix A. All three ideas that we considered use a module *Stream Router*, SR that scans the data stream, recognizes a "specific" semi-word associated with a device and sets high a valid signal for the device selection (devSelVal).

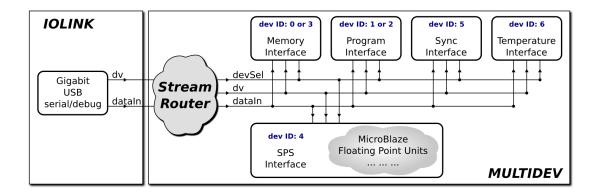


Figure 3.5 – Schematic logical representation of the IOP processor: on the left the |O| lolink that handles the I/O interfaces between IOP and the Janus host PC; on the right the devices of the multidev including the streamRouter. Each devices is identified using an ID. The gray bubble including MicroBlaze and floating point units is a possible expansion of the IOP.

The common assumption leading the development of a communication system is therefore that each device on the right side in Figure 3.5 is identified with a bit of a 16-bit word (bitwise mode) used as a mask. The "specific" semi-word of the previous paragraph is also a 16 bit word of the data stream coming from the IOlink that the SP recognizes and flags as a bitwise device selector.

Another task performed by the SR module is to recognize the other valid data (with informations for the devices) and forward them with a relevant data valid (dvOut) to the devices on the right.

The problem of recognizing destination information encoded within the stream is solved using the so called *encapsulate stream protocol*: a simple method in which each message has a header (violet frame in Figure 3.6) giving to the SR the following informations:

- the first word is always the device mask and is associated with a high value of the devSelVal signal;
- the second and the third word give to the SR the message length that will be registered and decreased each time that a valid word is processed by SR. When the length assumes the value zero then delivering of the current message to the multidev ends and next input data labelled with a valid flag is always considered as a new device mask, starting a new message.

An example of stream of data is shown in Figure 3.6.

Figure 3.7 shows the time diagram of the encapsulate stream protocol.

As the masks and the lengths are known at the moment of the message packing and is not required that they are inserted in the stream at the run time, the encapsulate

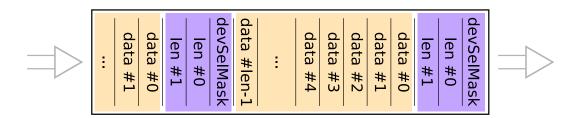


Figure 3.6 – Sample of an encapsulate stream.

clk .						
dvIn .						
dataln ⁻	devSelMask	Len #0	Len #1	data #0	data #i	data #len-1
devSelVal .						
dvOut .						
dataBus ⁻				(data #0)	data #i	data #len-1

Figure 3.7 – Time diagram of encapsulate stream.

stream protocol allows us to exploit buffers during the message arrangement on Janus host and assures therefore full bandwidth during the transfer.

A length composed by 32 bit (2 word of 16 bit) requires a 32 bit hardware counter in the FPGA but allows a stream of $\sim 10^9$ words of 16 bit or, in other words, a maximal message length of ~ 8 GB. If we consider uploading to each SP information about spins and couplings of a 3D lattice of side L = 64 using a single message we obtain a message size of:

$$(16 \cdot 4 \cdot 64^3) + \varepsilon \simeq 2 \text{ MB}$$

that is small enough. Note that a length of 16 bit is however too small for a such kind of data transfer.

As the stream router distinguishes only between words encoding information of the message destination and words building the message, could arise the question: "What happens if the machine goes out of control during a transfer?" In fact a stop/abort command for the machine is not allowed during communication using this protocol, but the simplicity of the implementation as in the hardware side as in the software side is a more relevant factor and this is the reason persuading us to use it on Janus.

3.2.5 Why a nearest neighbours network?

As presented in Chapter 2 the spin model simulations considered for the Janus system require only nearest neighbours interactions, therefore a nearest neighbours network of processors was the natural choice to implement a trivial partitioning of a given lattice and consequently an easy way to parallelize. Following this first trivial answer, it is possible to think about the historical reasons that guided the Janus collaboration to realize a such network. In view of the resources available on Virtex-4 LX160 devices used with the first prototype of Janus, we estimated to be able to update ~ 100 spins in parallel into a single FPGA. This led us to think that to simulate a big lattice it would be necessary to split it among different processors in order to increase ~ 10 times the parallel spin update rate. The nearest neighbours networks was therefore needed.

However the use of new FPGAs for the production system (we substitute Virtex-4 LX160 with Virtex-4 LX200) and the introduction of some optimizations in the VHDL code allowed us to increase by about one order of magnitude the number of parallel updates performed on a single FPGA so that it is possible to update a stand alone "big" lattice (i.e. up to 80^3) in a single node.

3.2.6 Why do boards have no direct link among them self?

There are examples of other parallel dedicated machines composed of standard processor linked with custom processors housed on boards with different features and goals (see for instance [2, 3]). In many of these cases each custom processor is considered as a node of a structured network (e.g. ring, mesh, crossbar, torus), therefore a link among nodes is required and thus, if nodes are housed on main boards, a link among boards is necessary. In the case of Janus however the development choice was to avoid a direct board connection.

Estimated performance for a board equipped with 16 FPGA was enough to justify the costs of the systems and the adding of a back-plane for the board to board connections would have increased costs, complexity of design and development time and could not be justified. As described in [4] in fact, development of a custom and dedicated machine makes sense if performance obtained and development time are balanced: increasing the complexity of Janus's design in order to add a direct board to board connection would have entailed a delay of some months to the delivery of the system while adding a negligible increase to the performance of the system.

Another reason to leave out board to board connection is the costs that would have a such architectural choice. A goal of the Janus project was in fact to take part to the Gordon Bell Prize award as an entry in the cost-efficiency machine section [5]. Cost-control was therefore a weak development guideline.

In the end the building of a backplane would also conflict to another idea of the Janus project, that was to build a modular supercomputer able to run in a rack but that could be in some sense also *portable*: a Janus board out of the rack is in fact ~ 30 cm $\times 30$ cm and can run connected to a common power plug and linked via a Gigabit cable to a standard laptop.

3.2.7 Why only 17 nodes per board?

In the development phase of a large system such as Janus it is important to consider the technology factor, by which we mean studying data sheets and application notes in order to know the limits and thresholds of the technology that one decides to use.

In the case of an FPGA based system the number of nodes housed in a board is not arbitrary, but arises from a technology factor. Each SP in fact has a set of 10 hardwired lines connecting it with the IOP and a set of 16 differential pairs (32 lines in all) to exchange data with the nearest neighbours. The number of signals hardwired in a board is very high and many physical layers are required in order to distribute all the signals with equalized delays and obtain a completely synchronous system. In the case of Janus the SP and IOP modules are composed of 12 layers while the mother board needs 16 layers to allow the wiring of signals.

Another limiting factor is the power supply problem. We decided for Janus to power each board independently via 48V input voltage in order to be compliant with most standard industrial power supplies, and distribute this voltage on the main board and convert it to 2.5V for each node using a DCDC converter. On board each node we convert than 2.5V to 1.2V so that we have both voltages to power I/O and core of the FPGA. Distribution and conversion of voltages are critical design factors and core voltage is a very critical limit for the VHDL design implemented on the FPGA.

The development strategy was a good choice for Janus and the use of consolidated technology makes the project not overly challenging on the engineering side. Outstanding help came moreover from the ability and experience of employees of Eurotech, the industrial partner of the Janus project.

3.2.8 Why do the nodes have no off chip memory?

The question of the on chip vs off chip memory is a never ending story and an open question in many large systems. New multi-core architectures, for instance Intel's Nehalem, adopt a cache level shared among cores and a faster (and smaller) cache level local for each core. Moreover they integrate on chip a memory controller in order to speed up access and increase bandwidth to main memory.

In the case of a reconfigurable computer each FPGA offers at present up to $\sim 5-8$ MB of on chip memory. The choice of the Janus collaboration was to consider this memory as sufficient and take advantage of the very high bandwidth given by the fact that the memories were within the FPGA. The reason of this drastic direction is justified by the fact that the data base of spin glass systems is in any case very small and embedded memories allow Janus to be a reference machine for these models for next 5 years. Data base of a spin system of lattice L needs in fact a quantity of memory of $\sim 4 \cdot L^3$ bits. If we search values of L solving the constraint $4 \cdot L^3 < 6$ MB

we found that the biggest lattice that can be stored within an FPGA has $L \simeq 116$. We suggest therefore that on chip memory is not a limiting factor for simulation of spin systems on Janus.

The presence of staging memories off chip introduce the problem of bandwidth between memory and FPGA that added to the low work frequency of Janus raises a problem of IO bandwidth related to the packaging of the FPGA. To update N spins per clock cycle we need in fact to access $4 \cdot N$ bits and if we suppose N = 256 (a fourth of the current update rate per clock cycle) we should have a bandwidth of 1024 bits per clock cycle. This is impossible to realize because the number of the IO pins on an FPGA used for Janus is ~ 1000. Using a double or quadruple data rate 2 or 4 data words per clock cycle are exchanged with the memory and could be possible to implement an interface with external memory with an acceptable bandwidth and a limited usage of IO pins, but part of the logic would need to be used to implement it and the timing constraints would increase the complexity of the synthesis process. It is remarkable that all this effort is required to achieve a fourth of the actual bandwidth.

If the on chip memory seems enough for the high-efficiency implementation of the spin models, a problem arises when we try to implement on Janus a model having a bigger data base or with an irregular structure. Section 3.5 describes a preliminary implementation of the random graph coloring problem on Janus and are described some difficulties encountered to store graph informations within embedded memories.

The choice to build computation nodes with no off chip memory can be viewed as an easy way to have a high performance machine for spin systems that on the downside introduces a limiting factor to the reconfigurability and generality of the Janus supercomputer.

3.2.9 Which clock frequency and why?

The main board of Janus is equipped with a variable frequency oscillator distributing a clock signal of 62.5 MHz to each daughter board via equalized lines.

The idea was to use for the Janus nodes a clock frequency not too high in order to keep the synthesis process of the VHDL code relative easy and non challenging in terms of time constrains. A low frequency moreover allows one to program a Janus board as a synchronous system and to consider each FPGA synchronous with each of the others avoiding tricky synchronization process during data transfer among nodes.

A good compromise frequency was 62.5 MHz directly coming from half of the Gigabit ethernet work frequency (125 MHz). Therefore we distribute to each FPGA a 62.5 MHz clock and then, only on the IOP with a clock multiplier obtained by built-in Digital Clock Manager (DCM), we obtain the 125 MHz to use for the I/O.

In this way the bandwidth are completely balanced because Gigabit ethernet produce 32 bit word each clock cycle (8 ns); Janus distribute words of 16 bits for each clock cycle, but works at half of the frequency (16 ns per clock cycle) so that the information distribution is balanced.

A posteriori we discovered that the choice of a low frequency played a key role in the Janus development. High clock frequency combined with very high density designs produce in fact a nightmare in terms of power consumption of the FPGAs: In some cases we obtain in fact the switching off of the chip after few seconds of run. Because of this we had to improve the power supply when we decide to adopt Virtex-4 LX200 instead LX160 (more details in 4.2.1 and in 4.2.3).

3.3 SP firmware: spin glass

3.3.1 Parallelism

The guiding line of our implementation strategy is to try to express all the parallelization opportunities allowed by the FPGA architecture, matching as much as possible the potential for parallelism offered by spin systems. Let us start by remembering that, because of the locality of the spatial interaction [6], the lattice can be split in two halves in a checkerboard scheme (we are dealing with a so-called *bipartite lattice* depicted in figure Figure 2.6), allowing in principle the parallel update of all white (or black) sites at once. Additionally, one can further boost performance by updating in parallel more copies of the system. We do so by updating at the same time two spin lattices (see later for further comments on this point).

The hardware structure of FPGAs allows exploitation of the full parallelism available in the algorithm, with the only limit of logic resources. As we explain below, the FPGAs that we use (Virtex-4 LX160 and Virtex-4 LX200, manufactured by Xilinx) have enough resources for the simultaneous update of half the sites for lattices of up to 80³ sites. For larger systems there are not enough logic resources to generate all the random numbers needed by the algorithm (one number per update, see below for details), so we need more than one clock cycle to update the whole lattice. In other words, we are in the very rewarding situation in which: i) the algorithm offers a large degree of allowed parallelism, ii) the processor architecture does not introduce any bottleneck to the actual exploitation of the available parallelism, iii) performance of the actual implementation is only limited by the hardware resources contained in the FPGAs.

We have developed a parallel update scheme, supporting 3D lattices with $L \geq 16$, associated with the Hamiltonian of (2.1.4). One only has to tune a few parameters to adjust the lattice size and the physical parameters defined in \mathcal{H} . We regard this as an important first step in the direction of creating flexible enough libraries of application codes for an FPGA-based computers. The number of allowed parallel updates depends on the number of logic cells available in the FPGAs. For the Ising-like codes developed so far, we update up to 1024 sites per clock cycle on a Xilinx Virtex-4 LX200, and up to 512 sites/cycle for the Xilinx Virtex-4 LX160. The algorithm for the Potts model requires more logic resources and larger memories, so performances lowers to 256 updates/cycle on both the LX200 and LX160 FPGAs.

3.3.2 Algorithm Implementation

We now come to the description of the actual algorithmic architecture, shown in Figure 3.8.

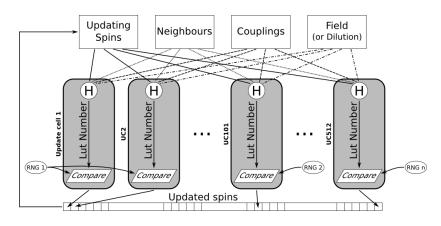


Figure 3.8 – Parallel update scheme. The spins that must be updated, their neighbors, the couplings (J) and all other relevant values are passed to the update cells where the energy is computed. The result is used as a pointer to a Look-up Table (LUT). The associated value is compared with a random number (RNG), and following the comparison, the updated spin value is computed.

In short, we have a set of update cells (512 in Figure 3.8): they receive as input all the variables and the parameters needed to perform all required arithmetic and logic operations, and compute the updated value of the spin variable. Data (variables and parameters) are kept in memory and are fed to the appropriate update cell. Updated values are written back to memory, to be used for subsequent updates.

The choice of an appropriate storage structure for data and the provision of enough data channels to feed all update cells with the data they need is a complex challenge; designing the update cells is a comparatively minor task. Hence we describe first the memory structures of our codes, followed by some details on the architecture of the update cells.

Virtex-4 FPGAs have several small RAM-blocks that can be grouped together to form bigger memories. We use these blocks to store all data items: spins, couplings, dilutions and external fields. The configurable logic blocks are used for random number generators and update cells.

To update one spin of a three dimensional model we need to read its six nearest neighbors, six couplings, the old spin value (for the Metropolis algorithm) and some model-dependent information such as a magnetic field or a dilution parameter for small variants of the model. All these data items must be moved to the appropriate update cells, in spite of the hardware bottleneck that only two memory locations in each block can be read/written at each clock cycle.

Let us analyze first the Ising models, considering for specifically the case L = 16. We choose to use an independent memory of size L^3 for each variable. This is actually divided into smaller memories, arranged so that reading one word from each gives us all the data needed for a single update cycle. We need $16^3 = 4096$ bits to store all the spins of one configuration. We have 16 vertical planes, and save each plane in a different memory of width 16 bits and height 16 (see Figure 3.9). In this simple case the logic resources within the FPGA allow us to update one whole horizontal plane in one clock cycle (because we mix the two bipartite sublattices of two different copies of the system, see the following discussion), and the reading rate matches requirements, as we need to read only one word from each of the sixteen memories.

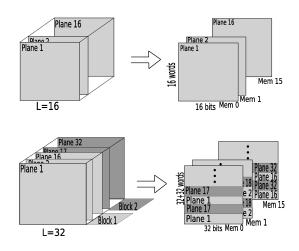


Figure 3.9 – Examples of the spin memory structure: L=16 and L=32.

The configuration is slightly more complex when the size of the lattice grows and the update of a full plane in just one clock cycle is no longer possible. In this case we split each plane into a variable number of blocks N_B , adjusted so that all the spins of each block can be updated in one clock cycle. The number of independent memories is L/N_B , as only these need to be read at the same time. The data words still have width L, while the height is $L \times N_B$ to compensate for the reduced number of memories. Considering L = 32, for example, we have a plane made of $32^2 = 1024$ spins, too large to be updated in one cycle (in the Xilinx Virtex-4 LX160). We split it in two blocks of $32 \times 16 = 512$ spins each. To read 16 lines every clock cycle we store the spins in 16 memories, each of width 32 bits and height 32×2 : the total size of the memory is still 32^3 bits.

As already remarked, we simulate two different copies of the system, that we call replicas in the same FPGA. This trick bypasses the parallelism limit of our MC algorithms (nearest neighbors cannot be updated at the same time, see [4]). We mesh the spins of the two replicas in a way that puts all the whites of one replica and the blacks of the other in distinct memories that we call respectively P and Q (see Figure 3.10). Every time we update one slice of P we handle one slice of whites for replica 1 and one slice of blacks for replica 2. Obviously the corresponding slice of memory Q contains all the black neighbors of replica 1 and all the white neighbors of replica 2.

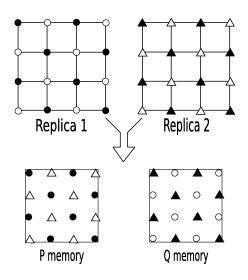


Figure 3.10 – Structure of spin configuration memories: meshing of replicas.

The amount of memory available in the FPGA limits the lattice size we can simulate and the models we can implement. In both the Virtex-4 LX160 and LX200 it is possible to simulate Edwards-Anderson models and some variants in 3D with size up to L = 88(not all smaller sizes are allowed). Because of the dramatic *critical slowing down* of the dynamics of interesting complex spin models these size limits are comfortably larger of what we can expect to be able study (even with the tremendous power made available by Janus) in a reasonable amount of (wall-clock) time: memory size is presently not a bottle-neck (as discussed in 3.2.8).

The lattice meshing scheme is maintained. With our reference FPGAs we can simulates three dimensional Potts model with at most L = 40 and a four dimensional Potts model with L = 16.

Things are even more complicated when one considers multi-state variables, as more bits are required to store the state of the system and all associated parameters. In the four state Potts model (see next section for details) the site variables need two bits and the couplings eight bits. In order to keep a memory structure similar to that outlined before we store each bit in a different memory. For example a lattice with L = 16 requires 16×2 memories for the site variables (they were sixteen in the Ising case), and 16×8 memories for the couplings.

We now come to the description of the update cells. The Hamiltonian we have written is homogeneous: the interaction has the same form for every site of the lattice, and it only depends on the values of the couplings (the fields and the dilutions, when model requires them). This means that we can write a standard update cell and use it as a black box to update all sites: it will give us the updated value of the spin (provided that we feed the correct inputs). This choice makes it easy to write a parametric program, where we instantiate the same update cell as many times as needed.

We have implemented two algorithms: Metropolis and Heat Bath. The update cell receives as input the couplings, nearest neighbors spins, field and dilution and, if appropriate, the old spin value (for the Metropolis dynamics). The cell uses these values as specified by (2.1.4) and computes a numerical value between 0 and 15 (the range varies depending on the model) used as an input to a LUT. The value read from the LUT is compared with a random number and the new spin state is chosen depending on the result of the comparison. Once again, things are slightly different for the Potts model due to the multi-state variables and couplings.

Our goal is to update in parallel as many variables as possible, which means that we want to maximize the number of cells that will be accessing the LUT at the same time. In order to avoid routing congestion at the hardware layer we replicate the LUTs: each instance is read only by two update cells. The waste in logic resources – the same information is replicated many times within the processor – is compensated by the ease of the synthesis process.

3.4 SP firmware: parallel tempering

To perform the parallel tempering algorithm (see Section 2.2.6) within a single FPGA node we can keep the firmware of the spin glass almost unchanged. The memory structure in this case is just a generalization of the standard case. Since the couplings and all other fields do not change between replicas, the only additional data that we need to store within the FPGA memories is the spin configuration of every configuration that we are going to simulate. This is achieved making the RAM Blocks grow in depth and using this extra space to store the N_T (number of temperatures in the parallel tempering) spin configurations, as if we were simulating a lattice of size $L^3 \times NT$ instead of L^2 . Although the information stored in those memories is slightly different than in the spin glass case.

When simulating only one temperature value we store the spin configurations of

two replicas in memories P and Q, meshing their black and white nodes, with both replicas being simulated at the same temperature. Using parallel tempering we want to simulate N_T copies of the same system, each evolving with a different T value, and thus a different LUT. To keep the configurations meshing trick seen before we tangle together sites belonging to replicas at different temperatures.

Although, in order to obtain the required information about the overlap we need to simulate two replicas at the same temperature. The only solution in this case is to simulate exactly the same system and temperatures in a different SP at the same time. We have already mentioned that the LUTs are replicated in order to have each of them read only by two update cells. Without parallel tempering there is just one set of LUT, from the single temperature we are simulating, that is read by all the update cells. In the case with PT the logic registers store two different sets of values, from the two temperatures being simulated in parallel. The update cells are properly arranged so that each can see either one or the other LUT. As a consequence, the spins updated by each cell will be constantly working with only one of the two available temperatures.

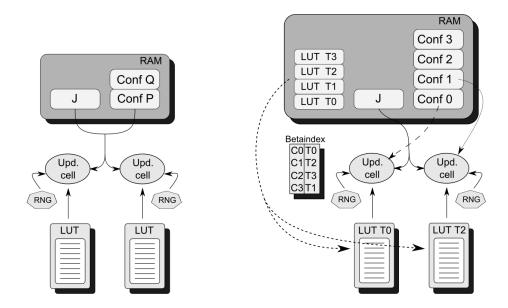


Figure 3.11 - Hardware implementation with and without parallel tempering (simplified representation). On the left we show the standard algorithm with a fixed temperature (i.e. a fixed LUT) and the relevant variables: couplings and spin configurations. The PT (on the right) needs more RAM space to store the extra LUTs and configurations, while the J memory is left unchanged. The array BETAINDEX indicates the simulation temperature of each system (source [7]).

We store all the temperature values and their corresponding pre-calculated LUTs inside the FPGA, the latter in dedicated RAM blocks (see Figure 3.11). The wo LUT sets needed by the actual simulation are written also into the LUT registers described above. An array, called BETAINDEX, keeps track of which system is being simulated

at each temperature. When two configurations have been completely updated we move to another pair of systems, once again each characterized by a (different) temperature value. The BETAINDEX values of the configurations that we are going to study point to the memory location where the corresponding LUT values are stored. Once these have been loaded into the LUT register we are ready to simulate the two new copies of the system.

The PT algorithm requires some new functions (and consequently new hardware blocks) as well. The first addition is related to the calculation (and storing) of the energy value of each configuration, necessary for the parallel tempering comparisons (see equation 2.2.20). This is done by simply running the update algorithm over each replica, but without actually updating the spin values in memories. Calculating the energy of a whole lattice takes exactly the same time as updating the same lattice. Once the energies of all configurations has been calculated, the values are used by the PT for the comparisons.

The other important feature is the calculation of the logarithms of random numbers. The PT algorithm works by comparing a random value with an exponential, namely accepting the temperature swap when:

$$r \leqslant e^{\Delta\beta\Delta E}$$
 with $r \in [0; 1[$. (3.4.1)

In this case we cannot resort to the LUT trick used for the MC updates, since it would be impossible to pre-calculate all the values of E. On the other hand, calculating the exponential value during the simulation would dramatically slow down the performances of the machine. Our best choice is thus to use the alternative equation

$$\ln r \leqslant \Delta \beta \Delta E \quad \text{with } r \in [0; 1[\tag{3.4.2})$$

exchanging the exponential function for a logarithm [8]. The logarithm operation is rather slow as well, but we don't need to wait until the simulations are over to start calculating it: since we need logarithms of random values we can generate these while the simulations are still running.

When the MC steps of all the systems have been completed, we just have to evaluate the products E and compare them with the pre-calculated $\ln r$ to decide whether to accept or reject the parallel tempering switch. If this is accepted we exchange the index values within the BETAINDEX array and move to the next pair of neighboring temperatures. Once all the temperature pairs have been evaluated we are ready to start again with the simulations, with each configuration working at a newly-assigned (hopefully reshuffled) temperature.

The PT implementation described in this section is the best choice, for small lattice sizes, because it is self-contained (within a single SP) and doesn't need too much time-wasting communications. On the other hand, due to limited FPGA memory resources,

this implementation works only with smaller lattice sizes $(L \leq 32)$ and a limited number of temperatures $(N_T \leq 128)$. For larger simulations with more temperatures it is necessary to change the approach and work with more SPs, but it is not yet implemented in Janus.

3.5 SP firmware: graph coloring

The key point of the hardware implementation of the algorithm for graph coloring is mapping of graph informations (color of each vertex and set of edges) within the FPGA memories. The update process performed for each vertex is similar to the case of the spin systems: the Metropolis algorithm is in fact the same. The challenge is extracting parallelism from from an irregular and unpredictable structure such as a random graph in order to perform parallel update of some nodes. Figure 3.12 shows a small example of random graph.

3.5.1 Memory organization

In the case of spin glasses the trick of the bipartite lattice (see Figure 2.6) allows to update whole set of the spin labelled as white in a single step and then all the black ones. As random graph is characterized by a non-fixed connectivity, is not possible to know in advance the number of edges for each vertex (that is equivalent to say that is not possible to know how many neighbours have each spin in an hypothetic irregular spin system).

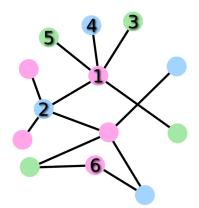


Figure 3.12 – Example of a small random graph. Detailed balance of the Monte Carlo requires that the state of the neighbour vertices should not change during the update of a given vertex. In this figure therefore is possible to update in parallel for instance vertices 2, 3, 4, 5, 6, but not 1, 2, 3, 4, 5.

The first idea is therefore to process the graph structure that we plan to study in order to find his degree of parallelism, i.e. we pre-process the input graph via a standard PC. This first step is performed to obtain a partition of the vertices in k subsets, S_k , each one containing p non adjacent vertices so that the number of vertices $|V| = k \cdot p$. This rearrangement of the graph assures that each vertex of a set S_i has no neighbours in the same set S_i and allows therefore the parallel update of all vertices of the vertices in S_i . This satisfies the detailed balance required by the Monte Carlo Metropolis algorithm. It is important to note that this operation is in principle not trivial (and not fast) when the average connectivity grows.

Any way after this preliminary operation we can suppose to have a graph in which is possible to update in parallel k vertices. The second challenge is represented by the mapping of graph structure within the memories of the FPGA. The only constraint is now given my the fact that we want update k vertices at the same time, so we would access informations related to the k updating vertices as fast as possible. It is clear that considering V, set of vertices of a given graph, and N_i , set of neighbours of the *i*-esim vertex, if we store informations by row in the form $\{V_i, N_{ij}\}$ with $0 \le i < |V|$, $0 \le j < |N_i|$ we obtain two disadvantages: i) we cannot know in advance width of the table in which we plan to store graph information (or, in the case of the FPGA, we cannot give a fixed dimension to a set of memories storing graph informations); ii) technological limitations connected with FPGA do not allow to access more than two locations per clock cycle and therefore, with this data organization we cannot take advantage from the previous parallelism extraction.

Considering to store graph data by columns is possible to take advantage from the rearrangement performed before. Each vertex has its own adjacency list stored below itself so that column *i* contains $\{V_i, N_{ij}\}$. Placing such structures one by one in a row we obtain a set of *b* blocks, $b = int\{|V|/k\} + 1$, each composed by a variable number of rows and each row with *k* elements. First row of the block b_0 contains first *k* vertices, $V_0, ..., V_{k-1}$ and following *g* rows the adjacency list ordered by columns, with $g = max\{|N_i|, 0 \le i \le k-1\}$; second block b_1 start on row g + 1 and contains $V_k, ..., V_{2k-1}$ followed by *g* rows, with $g = max\{|N_i|, k \le i \le 2k-1\}$, and so on.

In other words for each b so that $0 \leq b \leq |V|/k + 1$ we calculate

$$g(b) = \max\{|N_i| ; bk \leq i \leq k(b+1) - 1\}.$$

First row of block b is:

$$\{V_i \; ; \; kb \leqslant i \leqslant k(b+1) - 1\}$$

following lines of block b are:

$$\{N_{ij} ; kb \leqslant i \leqslant k(b+1) - 1 ; 0 \leqslant j \leqslant g(b) - 1\}$$

In case that V_i has no neighbours in a given position ij we assume to store a *dummy* word that is ignored by update engine. Blocks are stored one below the other and last row of a block contains a flag indicating the end of a block. We call this memory

structure topo-memory because in some sense it give us topological informations about the graph. It is important to note that in the topo-memory are stored only pointers and no informations about vertices colors. Pointers stored in the topo-memory point to locations of another memory structure, called *color-memory* storing colors of vertices in a given order (k by k).

A such organization allows to use the topo-memory as a read only memory (we write it only during initialization phase) and to preform read-write access to the color memory during the update phase. Another advantage of this organization comes from the possibility to access informations of a set of k vertices/neighbours in a single clock cycle (if we neglect the small latency of the internal memories) and update therefore a set of k vertices in a number of clock cycles comparable with g(b) (that is relatively small in case of graph with small mean connectivity).

A relevant disadvantage of this organization comes from the fact that we waste a non negligible part of memory storing dummy informations: in fact in a block with g(b) = 8 and k = 16 if the other 15 vertices have few or no neighbours we are wasting many memory locations. This disadvantage is strongly related with graph structure, but there is no way to optimize it because is intrinsic in the formulation of the problem of random graph coloring.

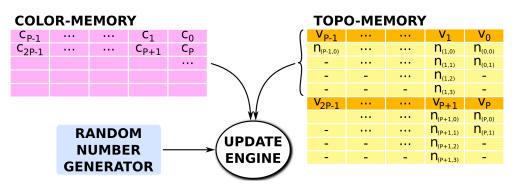


Figure 3.13 – Logic blocks of the FPGA implementation of the Monte Carlo update for a random graph.

Figure 3.13 gives a graphical representation of the hardware implementation of the graph coloring firmware. The update engine receives a set of k pointers to k locations of the color-memory and reads color informations of the current states of the updating vertices using the addresses received from topo-memory. With the following sets of pointers coming from topo-memory in set of k starts the update process: each pointer corresponds to a color of a neighbour and a counter trace the number of color conflicts generated comparing color of updating vertex with color of his neighbours. k update engines works in parallel performing therefore the parallel update of k vertices. As well in the spin system, random number generation play a key role. Conflict counting of the current set of k vertices ends when the update engine read the bit indicating

the end of the block from the topo-memory. After that, the Metropolis Monte Carlo update is performed and the new color of the k vertices is produced and stored in the color-memory. A new update process can therefore starts. It is clear that a carefully use of pipelining improve the performance.

3.5.2 Janus limitations in graph coloring

FPGAs allow a large degree of freedom in firmware design: in the case of the memory organization described above the idea to have a memory with a system of pointers and a memory with color informations is amazing if we think that we are working with hardware. State of the art of the graph coloring studies have however a memory requirement that seems to be not compatible with Janus.

The device Virtex-4 LX160 have 288 RAM blocks (we studied this prototype implementation on the first implementation of Janus). Each RAM block is 16 Kbit and can be freely configured from 16 Kwords \times 1 bit-word to 512 words \times 36 bit-words as shown in figure 3.14.

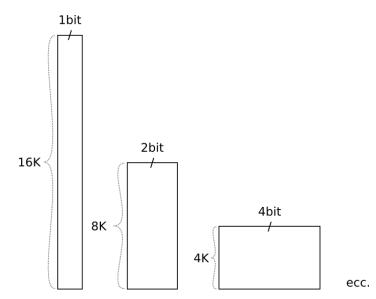


Figure 3.14 – Available memory setup in a Virtex-4 FPGA

Each word of the topo-memory (TM) is $(K \log N + 1)$ bits (the "+1" bit is the bit indicating the end of the current block). The color-memory (CM) has $\frac{N}{K} 4C_m$ rows¹.

Total TM size =
$$(K \log N + 1) \cdot \frac{N}{K} \cdot 4C_m$$
 bits

In the case of our first implementation:

TM size = $(64 \cdot 14 + 1) \cdot 256 \cdot 8 = 897 \cdot 4096$ bits

¹We consider a number of rows that is 4 times bigger than the average connectivity. This choice is may be too "conservative"

Using the Xilinx memory, we need 225 blocks with the configuration 4K words \times 4 bit-word to build the topo-memory.

Each word of the color-memory use $K \cdot \log B$ bits (with $\log B$ the number of bits used to represent a color). The CM has $\frac{N}{K}$ rows and will be replicate $\frac{K}{2}$ times in order to access k vertices data in a single clock cycle (using the dual port configuration).

Total CM size =
$$K \cdot \log B \cdot \frac{N}{K} \cdot \frac{K}{2} = \frac{K}{2} \cdot \log B \cdot N$$
 bits

In our implementation:

CM size =
$$32 \cdot 2 \cdot 2^{14} = 32 \cdot 32.768$$
 bits

Using the Xilinx memory, we need 4 blocks with the configuration 512 words \times 36 bit-word to build the CM, but we need to replicate the CM structure K/2 times, i.e. we need $4 \cdot 32$ Xilinx RAM blocks to build the complete CM parallel system.

From this very preliminary study we have two conclusions:

1. the total number of bits for the TM and the CM is:

$$(897 \cdot 4096) + (32 \cdot 32.768) = 3.674.112 + 1.048.576 = 4.722.688$$

This number is smaller than the total available bits in a Virtex-4 LX160.

2. the total number of blocks needed for TM and CM is: 225 + 128 = 353This number is bigger than the total number of blocks available in a Virtex-4 LX160.

Balancing of 1 and 2 were the optimum solution, but seems to be difficult to reach with the present realization of Janus: embedded memory are in fact the limiting factor for a more efficient implementation.

Moreover the state of the art of random graph studies requires to be able to approach graphs of more than 50000 vertices, that is far to the present realization with Janus.

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Welcome to the jungle! We got fun 'n' games, we got evrything you want, honey, we know the names. We are the people that can find whatever you may need.

Guns N' Roses

4

Architectural details of Janus

In this chapter I will present details of the Janus architecture with special attention to the structure of the Input Output Processor, IOP.

This chapter is intended as a technical document describing the hardware components and the VHDL-coded firmware: I will introduce briefly different type of FPGAs and the basic components of Janus hardware from the point of view of the "mason" building the real system, activity that I have followed during my whole PhD studies till the Janus system was finally up and running.

A relatively large section describes the VHDL entities that enable the IOP hardware; I have designed these entities in the first half of my PhD studies. I then introduce the firmware/software environment developed to test and validate the Janus hardware.

A final short section describes some engineering problems that we encountered and for which we have to provide "acceptable solutions" during the development of the system.

The VHDL code developed for the IOP and explained in this section is available on the Janus web page: http://df.unife.it/janus/ in *Tech* section.

4.1 FPGA: different flavours

The FPGA market offers devices with different features and containing different mixes of functionalities to suit a relatively wide set of different application areas. The present technology allows to embed, for instance, a simple microprocessor within an FPGA and have therefore a chip including a general purpose processor able to run an operating system and, at the same time, a non negligible quantity of configurable logic Other devices, optimized for high speed communications, embed for instance several (at present 8 to 20) high speed interfaces (e.g. PCI or Gigabit hard-cores). Other FPGAs offers a huge quantity of embedded memory.

The FPGA panorama is therefore various and the choice of a device for Janus was driven by the simple idea that the only important feature is the availability of memory and logic elements in order to store lattices as large as possible and to house the highest number of update engines. Large on chip memory size and many logic elements are obviously conflicting requirements; each FPGA family offers different trade offs.

Our preliminary prototype was developed in 2005 using a PCI development kit [1] housing an Altera Stratix S60 FPGA, containing ~ 57000 logic elements and ~ 5 MB of embedded memory¹.

	Configu	urable Logi	c Blocks ((CLBs) ⁽¹⁾	Block RAM									
Device	Array ⁽³⁾ Row x Col	Logic Cells	Slices	Max Distributed RAM (Kb)	XtremeDSP Slices ⁽²⁾	18 Kb Blocks	Max Block RAM (Kb)	DCMs	PMCDs	PowerPC Processor Blocks	Ethernet MACs	RocketIO Transceiver Blocks	Total I/O Banks	Max User I/O
XC4VLX15	64 x 24	13,824	6,144	96	32	48	864	4	0	N/A	N/A	N/A	9	320
XC4VLX25	96 x 28	24,192	10,752	168	48	72	1,296	8	4	N/A	N/A	N/A	11	448
XC4VLX40	128 x 36	41,472	18,432	288	64	96	1,728	8	4	N/A	N/A	N/A	13	640
XC4VLX60	128 x 52	59,904	26,624	416	64	160	2,880	8	4	N/A	N/A	N/A	13	640
XC4VLX80	160 x 56	80,640	35,840	560	80	200	3,600	12	8	N/A	N/A	N/A	15	768
XC4VLX100	192 x 64	110,592	49,152	768	96	240	4,320	12	8	N/A	N/A	N/A	17	960
XC4VLX160	192 x 88	152,064	67,584	1056	96	288	5,184	12	8	N/A	N/A	N/A	17	960
XC4VLX200	192 x 116	200,448	89,088	1392	96	336	6,048	12	8	N/A	N/A	N/A	17	960
XC4VSX25	64 x 40	23,040	10,240	160	128	128	2,304	4	0	N/A	N/A	N/A	9	320
XC4VSX35	96 x 40	34,560	15,360	240	192	192	3,456	8	4	N/A	N/A	N/A	11	448
XC4VSX55	128 x 48	55,296	24,576	384	512	320	5,760	8	4	N/A	N/A	N/A	13	640
XC4VFX12	64 x 24	12,312	5,472	86	32	36	648	4	0	1	2	N/A	9	320
XC4VFX20	64 x 36	19,224	8,544	134	32	68	1,224	4	0	1	2	8	9	320
XC4VFX40	96 x 52	41,904	18,624	291	48	144	2,592	8	4	2	4	12	11	448
XC4VFX60	128 x 52	56,880	25,280	395	128	232	4,176	12	8	2	4	16	13	576
XC4VFX100	160 x 68	94,896	42,176	659	160	376	6,768	12	8	2	4	20	15	768
XC4VFX140	192 x 84	142,128	63,168	987	192	552	9,936	20	8	2	4	24	17	896

Figure 4.1 – Table summarizing main features of Xilinx Virtex-4 FPGA family (source [2]).

The first two Janus prototype boards developed in 2006 had Xilinx Virtex-4 LX160 FPGA while the final realization of the system was based on Xilinx Virtex-4 LX200. Figure 4.1 summarizes the main features of Xilinx Virtex-4 FPGA, that are divided in three main families: LX, optimized for logic and memories; SX, housing a huge

¹Altera Stratix I devices embed memories organized in three different size: 574 blocks, called M512 RAM, that can be configured up to 32×18 bits, 292 blocks, called M4K RAM, configurable up to 128×36 bits and 6 blocks, called M-RAM, configurable up to $4K \times 144$ bits.

number of DSPs; FX, embedding up to two PowerPCs and high speed IO interfaces. For the aim of Janus the LX family represents the best Xilinx option since we are not particular interested in DSPs or embedded traditional processors.

The choice between Altera or Xilinx FPGAs has not been fully trivial. To first approximation both families have approximately the same amount of logic elements², but a different amount of on chip memory. Altera Stratix-II FPGAs offer ~ 8 Mb (see [3] for details) organized in three degrees of granularity allowing us to efficiently exploit only ~ 50% of it. Conversely Xilinx Virtex-4 LX200 FPGAs have ~ 6 Mb (see [4] for details) of embedded memories³ made up of relatively small blocks that we can use very efficiently for our design. Fortunately this decision based on technical features of the devices was also supported by a better price/performance ratio.

4.2 Structure of a Janus board

As described in previous chapters, Janus is a modular system; each Janus board is composed of 16 scientific processors, SP, an input/output processor, IOP, and a processing board, PB, in which SPs and IOP are plugged via high speed Samtech connectors (details of connectors and modules are shown in Figure 4.2. Each Janus board is also housed in a box providing power supply and cooling.

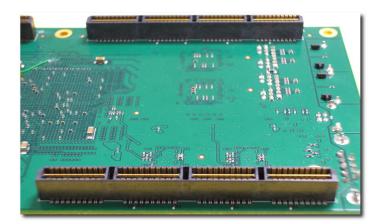


Figure 4.2 – Detail of the bottom view of the IOP with the Samtech connectors used to plug SPs and IOP to the PB.

Description of a Janus component amounts to a large extent to describe the connections of the I/O pins of the FPGA. For this reason I will present in the following paragraphs the I/O interfaces of each FPGA, that is roughly equivalent to describe

 $^{^{2}}$ We consider the largest devices of both FPGA families available when we had to make a final decision: Altera Stratix-II 180 and Xilinx Virtex-4 LX200.

³We neglect in this analysis the so-called distributed memory.

the structure of a Janus board apart for the power supply system. A paragraph is also dedicated to the description of the Janus box.

4.2.1 SP

All SPs are directly connected in a mesh (i.e. each board represents a nearest neighbours network with toroidal enclosure) on a 16 lines full duplex differential bus (16 differential input lines + 16 differential output lines). Moreover each SP has a private single ended full duplex link with the IOP composed of 10 lines.

Other single ended lines reaching the SP from the PB via connectors are:

clock:	2 input clock lines. A common clock frequency is generated by an oscil-
	lator on the PB and distributed to all modules; 2 additional FPGA I/O $$
	pins are available for feedback for each input clock; FPGA pins in close
	proximity of the clock signals are not used to reduce noise.
syncIn:	4 input lines used by FPGA firmware for synchronization (the same set
	of synchronization signals is shared by 4 SPs).
syncOut:	4 output lines used by FPGA firmware to send synchronization messages
	to the IOP (synchronization signals from SPs to IOP are point to point).
progLine:	21 lines driven by the IOP programming interface in order to configure
	each SP as needed; we use the $select map$ configuration mode that re-
	quires 8 bits data bus connected with 8 FPGA I/O pins and 13 lines
	connected to an equal number of dedicated FPGA pins driving the con-
	figuration.
spReset:	2 input lines coding various level of resets.

The rest of the pins used for input/output are dedicated to test points (16 pins), temperature sensor lying on each daughter card (3 I/O pins + 2 dedicated pins) and reset button (1 pin). These lines are however on board signals and have no connections with the PB. Four more dedicated on board pins are used for JTAG configuration that we used only during the prototype phase and for hardware debug via Chipscope [5]. More details about configuration are explained in 4.3.5 and in [6].

The total number of input/output pins⁴ for the FPGA of each SP are therefore 312, equivalent to $\sim 33\%$ of the I/O pins available. Table 4.1 summarizes pins usage in an SP module (neglecting power and ground pins).

Part of the SP area houses the 2.5 V to 1.2 V DCDC power modules. This last stage of power conversion suffered from several engineering problems, as, in the early

 $^{{}^{4}}$ Exact number calculated in view of the entries of the user configuration file UCF used by the synthesis tool.

Logical function	$\mathbf{Pin}~\#~\&~\mathbf{type}$	Direction	
Link to SP North, South, East, West	4×16 LVDS pairs	Output	
Link from SP North, South, East, West	4×16 LVDS pairs	Input	
Link to IOP	10 LVCMOS	Output	
Link from IOP	10 LVCMOS	Input	
Sync to IOP	4 LVCMOS	Output	
Sync from IOP	4 LVCMOS	Input	
Test points	16 LVCMOS	Output	
Clock	6 LVCMOS	4 Input 2 Output	
Reset	3 LVCMOS	Input	
Temperature concer	3 LVCMOS	$2 { m InOut} + 1 { m Input}$	
Temperature sensor	2 dedicated pins	-	
JTAG	4 dedicated pins	-	
Configuration channels	21 dedicated pins	-	

Table 4.1 – Pin assignment summary of the SP module.

design phases, we severely underestimated the power needed by our FPGAs. These problems (and how we solved it) are described in details in Section 4.5).

4.2.2 IOP

The hardware design of IOP is more complex then the SP because this daughter card performs numerous and various tasks and therefore houses a higher number of components. Looking at the connectors with the PB, the presence of a 10 bit bus per direction for each SP requires that a huge number of lines reach the PB.

Components on board are: 2 Gigabit PHYs (model Marvell 88E1111), 1 UART interface chip (model Maxim MAX3381E), 1 USB-bridge (model CP2102), 2 EEPROMs for FPGA configuration on boot up (model XCF32PVOG48C), 2 staging memories (model NEC μ PD44321361), 1 temperature sensor (model MAX1617A), 2 JTAG interfaces, 1 oscillator generating the clock for Gigabit PHYs, 1 LEMO connector to accept external clock and 1 reset button. Figure 4.3 shows an the IOP complete schematic design with its main components.

Pin assignment of the IOP FPGA can be summarized as follows:

- 105 pins dedicated to staging memory: 64 pins for data, 20 pins for address, 2 pins for clocks, 8 pins for parity, 1 pin for write enable and 10 pins for controls.

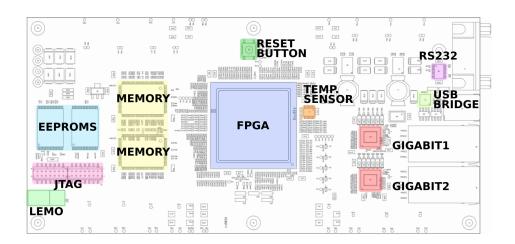


Figure 4.3 – IOP schema with highlight of the main components.

- 320 pins used for data transfer with the SPs: 10 pins per direction for each SP.
- 80 pins for synchronization signals: 4 signals reach the IOP from each SP $(4 \times 16 = 64)$ and 4 signals start from IOP and are shared with a set of 4 SPs $(4 \times 4 = 16)$.
- 8 pins for reset signals shared in the same way as the synchronization signals above.
- 44 pins for select map configuration of the SPs. We organize configuration of SPs in 2 channels: each channel has a 8 bit data bus and 6 control signals; each SPs has a chip selection bit (corresponding to 16 pins).
- 48 pins to connect FPGA to Gigabit ports. We use 2 Gigabit links: each one requires 8 pins for data to send, 8 pins for data to receive and 8 pins for controls.
- 5 pins for temperature sensor: 3 I/O pins + 2 dedicated pins.
- 1 pins for master reset of the FPGA, connected with an on board button.
- 5 pins for serial interfaces: 2 for UART interface and 3 for USB bridge interfaces.
- 8 pins for clock. We accept 2 global clocks coming from PB and for each clock we need 2 pins for feedback and 1 test point.
- 16 test points.
- 16 dedicated pins for configuration of the IOP FPGA via JTAG or via EEPROM.

The total number of the I/O pins⁵ used on the IOP is 654, equivalent to ~ 65% of the I/O available. More than 50% (452) of them must be routed to the connectors because the corresponding signals go to SPs.

⁵Exact number calculated in view of the entries of the user configuration file UCF used by the synthesis tool.

4.2.3 PB

The processing board, PB, is apparently the simplest Janus component, because it does not house complex BGAs and has a low density of components. In spite of that it performs three critical functions: distribution of power supply, generation of clocks for all Janus elements, routing of the signals linking IOP with SP and vice versa.

In order to avoid to distribute high currents we decided to feed the PB at a voltage of 48 V, a standard in the telecom industry. Each SP has an independent DCDC module (model Powergood ESC4825-15-X) converting from 48 to 2.5 V. The advantage of this solution lies in the fact that dimensions of these converters are moderate and efficiency is high (90% measured during intensive tests). This setup is currently used very close to its limit performance. We have tested a few firmware configurations for the SPs for which our converters are not able to supply the required current. Fortunately this situation was not a real problem for the physics simulation because with a careful placing of the FPGA resources all the firmwares becomes compliant with Janus hardware specification. A new conversion system able to fully support FPGA resources is therefore under test.(more details in Section 4.5).

Clock distribution is another key point played by the PB. An oscillator is placed in the middle of the board and 16 buffers assure an equalized distribution of the clock signal to each SP. We decide to combine an oscillator with a configurable frequency multiplier with zero delay and two output (model Cypress CY2302SXC-1) allowing to generate input frequency divided by 2 or multiplied by 2, 4, 8, 16 with an output range between 10 MHz and 133 MHz. This design strategy allow to upgrade the clock frequency of the system simply changing a jumper or replacing the main oscillator clock.

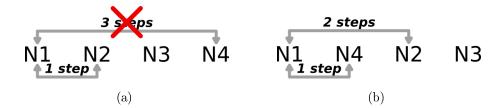


Figure 4.4 – (a) ideal 1D enclosure. (b) node replacing that allows the data to cover always at most 2 steps.

The last important role played by the PB is the routing of the signals linking IOP to SP and vice versa. As a Janus board is considered a synchronous system with relative slow frequency, is important to route signals via equalized paths so that, for instance, a broadcast message starting on the rising edge of the IOP clock reach all the SP FPGAs at the same time (within tolerance allowed by the devices). This is very relevant for the routing of the nearest neighbour links, since the most obvious placement of the SPs would imply physical paths for the links of widely different lengths as explained in 1D by Figure 4.4a. We can go around this problem with a simple change in the placement of the nodes, as shown in Figure 4.4b.

Following this idea the actual position of the SPs in a Janus board are not as in Figure 4.5a but as in Figure $4.5b^6$.

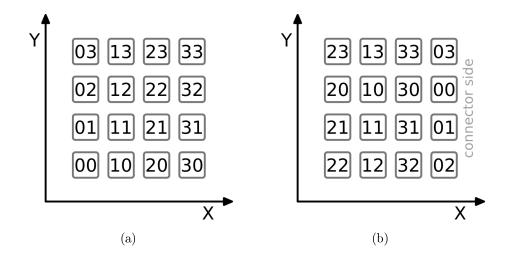


Figure 4.5 - Each SP is represented by a pair of numbers representing the Cartesian coordinates of the SP in a XY-plane. (a) Cartesian placement of a toroidal mesh.
(b) Placement of the SPs in a Janus board, partially equalizing the length of the paths between nearest neighbours.

4.2.4 Janus box

Each board is housed in a case providing power supply, plugs connecting to the host PC and a complex fan system needed remove heat and keep temperature within a reasonable working range.

Connectors are placed on the back panel of the box and are arranged as shown in Figure 4.6a. There is a power plug and a power switch and the following interfaces: two Gigabit plugs, one USB and one RS232 connector and two JTAG interfaces. In the present system setup only one Gigabit plug and optionally the serial line for debug are connected. Two JTAG interfaces are available in order to reconfigure independently the IOP FPGA or the EEPROM placed on the IOP.

Our cooling system is provided by static heat sink of different height placed on the FPGAs and a system of ten fans arranged in three lines: four at the front, four in the middle and two on back panel so air flows from the front to the back. Heat sinks are placed so that the 4 FPGAs near the front panel, where air has the lowest

⁶Note that the Y axis does not respect exactly the rules minimizing the path described in previous paragraph: the reason of that is related with the problems of routing of the signals within the PB.

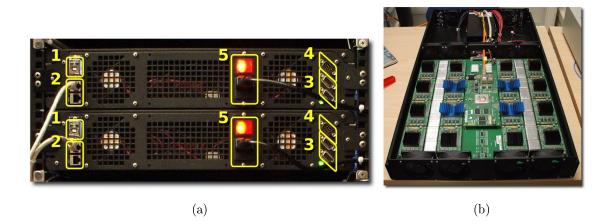


Figure 4.6 – (a) Back panel of two Janus boxes with the highlight of the available connectors: USB (1), double Gigabit channel (2), JTAG interfaces to configure the IOP FPGA or the EEPROMs (3), UART (4) and main power switch (5). (b) Janus box with cooling system.

temperature, have smaller heat sinks, 8 FPGAs placed in the middle of the board have medium-size heat sinks and 4 FPGAs close to the rear panel in places take advantage of larger heat sinks. Figure 4.6b shows the whole Janus box with a board full of IOP and SPs with fans and cooling system while Figure 4.17 shows a graph of the measured temperature within a Janus box.

4.3 The IOP in depth

In this section I will introduce a detailed description of the IOP firmware allowing the Janus user to access hardware resources described in previous section. Development of the VHDL code for IOP module was driven by some general guidelines described in 3.2.2. Figure 3.5 shows the outline of the organization of the logical blocks within the IOP. Basically the data stream coming from the I/O interfaces (UART, USB or Gigabit) is processed by *IOlink* block generating a stream of 16 bit words and one bit flagging the validity of the word. The stream generated from the IOInk is processed by the *streamRouter* that detects the word encoding the target device, flags it and forwards data words, data valid flag, and target flag to all devices of the IOP. Each device activates itself only if the word flagged by target flag encodes an ID corresponding with his own hard coded ID. Such a structure allows us to add devices on the right side of Figure 3.5 without the need to change the VHDL code of IOlink and StreamRouter.

This structure can be further developed to provide additional functionalities in the IOP without changing its overall organization. For instance the IOP might perform with very low latency some functions that require data gathered from all SPs an return its results to the SPs. The added unit performing this function would be accessed by the host PC through the IOlink and StreamRouter using a new ID. We have started to study this solution in order to allow a future implementation of the parallel tempering (see 2.2.6 and 3.4) algorithm that collects temperature information from SPs, performs the swap algorithm and sends new temperatures to the proper SPs.

Figure 4.7 shows the hierarchy of the VHDL entities in the present IOP firmware. I will present in the next paragraphs the main features of some of these entities that I coded during my PhD: in particular it is interesting to view details about clock generation, double data rate communications and all the entities included within the so called *multidev* box.

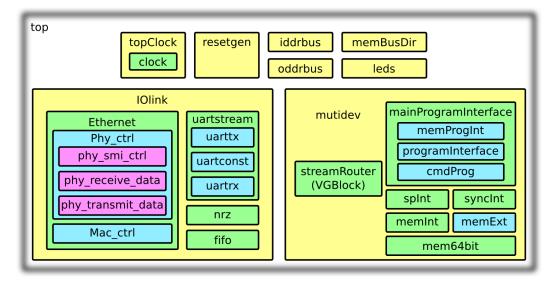


Figure 4.7 – VHDL entity hierarchy of the current implementation on the IOP.

4.3.1 Clock handling: topClock

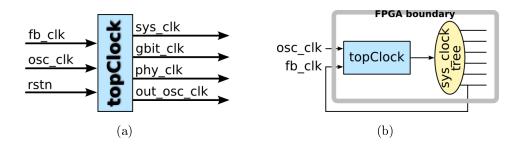


Figure 4.8 – (a) Interface of the VHDL entity handling IOP clocks. (b) Clock feedback in the Janus system.

The Xilinx Virtex-4 FPGA houses 12 Digital Clock Managers (DCMs) that deskew

clocks, multiply and divide frequencies, shift and align clock phases and reconfigure clocks dynamically. Xilinx synthesis tools offer software support to configure DCM in the form of a parametric black box that can be included in a design like a standard VHDL component.

The entity topClock, represented in Figure 4.8a, includes one of this DCM driver and some additional logic in order to generate three different clock frequencies:

sys_clk	the system clock with a frequency of 62.5 MHz (16 ns clock period);
gbit_clk	the clock for the Gigabit interface: it has a frequency of 125 MHz (8 ns clock period) and is used in a small part of the IOP design;
phy_clk	a 1 MHz clock used by the UART interface;
out_osc_clk	a copy of sys_clk connected to a test point.

It is interesting to note that the entity topClock has an input called fb_clk that is the feedback signal used by DCM to keep on the same phase the input clock osc_clk and the signal at the end of the on chip distribution tree of sys_clk. This feedback system and the fact that the input clock source is shared by each FPGA on the board are enough to assure that data transfers between nodes are synchronous (at least for the slow frequencies that we use, 62.5 MHz). Figure 4.8b shows the feedback strategy for the clock.

4.3.2 Double data rate: iddrBus and oddrBus

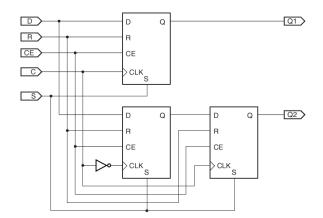


Figure 4.9 – Input DDR in SAME_EDGE mode (source [4]).

Virtex-4 devices embed dedicated registers placed near FPGA pad to implement input or output double-data-rate (DDR) registers. This feature is used by instantiating the IDDR/ODDR primitive. There is only one clock input to the IDDR/ODDR primitive. Falling edge data is clocked by a locally inverted version of the input clock. The IDDR/ODDR primitive supports different modes of operation; we use SAME_EDGE mode for both directions, input and output. Figure 4.9 shows, for instance, input DDR registers and the signals associated with the SAME_EDGE mode. Figure 4.10 shows the timing diagram of the input DDR using the SAME_EDGE mode.

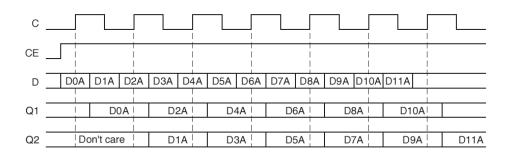


Figure 4.10 – Input DDR Timing in SAME_EDGE mode (source [4]).

The entities iddrBus and oddrBus are wrappers that instantiate respectively a set of N IDDR or ODDR Xilinx black boxes in order to obtain a double-data-rate bus of Nbits. We use them in the links between IOP and SPs to double the bandwidth.

Implementation of this technology is useful, but it becomes less reliable when the logic resource usage of the FPGA exceeds ~ 85%. In this case in fact the duty cycle of the system clock becomes appreciably asymmetric: the part of the design working with rising edge of the clock works fine because the DCM aligns phase looking at the rising edge of the input clock and the feedback clock, but 50% duty cycle constraint is not meet and logic using both edges, such as DDR I/O registers, becomes completely unstable (see also 4.5).

4.3.3 Message routing: StreamRouter

StreamRouter is a key entity in the IOP architecture (its interface is shown in Figure 4.11a): all messages coming from IOlink pass trough it and are routed to destination device following an easy protocol, called *encapsulated protocol*.

We assume that each target device has a device ID bitwise coded in a 16 bit word, called devSelMask. Each message starting from host PC (under the form of data stream) must have an header composed of three words of 16 bits coding respectively: word $\#1 \rightarrow devSelMask$ coding target device;

word $#2 \rightarrow len #0$ less significant 16 bits of the length of the following message (unit of measurement is 16 bit word);

word $\#3 \rightarrow len \#1$ most significant 16 bits of the length above.

Data following word #3 contains message for the target device. For a graphical representation of a message see Figure 3.6 in previous chapter.

StreamRouter decodes information contained in the header, removes words coding

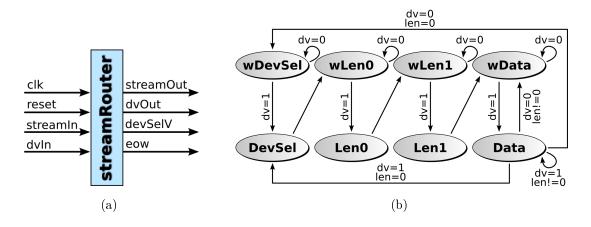


Figure 4.11 – StreamRouter block: (a) interface of the entity, (b) bubble diagram of the state machine.

the length of the message, forwards it on the data bus reaching the IOP devices with corresponding data valid flag and generates a bit flagging the data word that contains devSelMask (i.e. the target device mask).

These functionalities are implemented via a finite state machine with 8 states and 2 control signals (input data valid and reset). Behaviour of this state machine is rather simple. The Initial state is WDEVSEL: the machine waits the "first" word of the stream, that is always the devSelMask. With the first data valid the machine passes to state DEVSEL. While in state DEVSEL, device selection mask is latched and the output devSelValid is asserted. This requires one clock cycle; then the machine waits in state WLEN0 for first part of the length (less significant 16 bits). When second data valid arrives the machine goes to state LENO in which the length is saved and both outputs devSelValid and dvOut are deasserted. The same behaviour is required for states WLEN1 and LEN1 in which the machine latches the second part of the length. The machine then waits for valid data in state WDATA (outputs are deasserted) and passes to state DATA with the first data valid following the length. In state DATA the output dvOut is pulled high and the register storing the length decreases on each dvIn = 1. The machine remains in state DATA until length becomes 0. When length counter reaches 0, in correspondence with the last valid data the output signal eow (end of worm) is asserted for one clock cycle and then the machine goes back to WDEVSEL waiting for a new mask. Figure 4.11b shows the bubble diagram of the streamRouter state machine.

4.3.4 Memory controller: memExt

IOP houses two memory chips with a data word length of 32 bits and a depth of 1M words. We organize them in parallel in order to have an addressable space of 1M words of 64 bits and therefore a larger bandwidth.

As data coming from StreamRouter are organized in words of 16 bits, the entity memExt, in case of write command, has to buffer four words and write them when buffer is full. In case of read command, data coming from memory has to be split in four 16 bit words so the read process from memory cannot be continuous. Entity memExt implements therefore a state machine controlling these data transfers with the right timing. Code 4.1 shows the interface of the VHDL entity implementing memory controller on Janus. Analyzing it, we can identify 3 groups of ports beyond clock and reset.

antitu mamBut is
<u>entity</u> memExt <u>is</u> <u>generic</u> (
ID : integer
);
port (
IOP main clock (62.5 MHz) and reset
clk : <u>in</u> std_logic; reset : <u>in</u> std_logic;
reset : <u>in</u> std_loğic;
Ports to/from StreamRouter and I/O
ioDataIn : in std logic vector(15 downto 0):
ioDvIn : in std_logic;
ioDevSelV : in std_loğic;
ioGbitStop : <u>in</u> std_loğic;
<pre>ioDataIn : in std_logic_vector(15 downto 0); ioDvIn : in std_logic; ioDevSelV : in std_logic; ioGbitStop : in std_logic; ioDataOut : out std_logic; ioDataOut : out std_logic.vector(15 downto 0);</pre>
ioEot is high for the last cycle of valid data
ioEot : <u>out</u> std_logic;
ioDriver : <u>out</u> std_logic;
Ports to/from Memory
Chip Enable to activate/deactivate the NEC memories
(not used in this version)
memCe : out std_logic;
To activate the burst mode
memAdv : <u>out</u> std_logic;
Write enable active high
(but the NEC memories use an active low WE)
memWe ; <u>out</u> std_logic;
Byte enable (not used in this version)
<pre>memBe : out std_logic_vector(7 downto 0);</pre>
memAddr : <u>out</u> std_logic_vector(19 <u>downto</u> 0);
memDataIn : <u>out</u> std_logic_vector(63 <u>downto</u> 0);
<pre>memDataOut : in std_logic_vector(63 downto 0);</pre>
<pre> Ports to/from other devices (= program interface)</pre>
progDv : <u>out</u> std_logic;
progStop : <u>in</u> std_logic
);
<u>end</u> ´memExt;

Code 4.1 – VHDL entity interface of the memory controller on Janus.

First set of ports, with prefix io, are connected with the streamRouter and the IOlink: they are ioDataIn with two flags ioDvIn, indicating that the incoming 16 bit data word is a valid one, and ioDevSelV, flagging the word containing device target information; ioGbitStop is a control bit coming from IOlink triggering a break on a data transmission. ioEot flags the end of transmission, while ioDriver is the direction selection for the bidirectional ports at the top of the design (driving inout type ports).

The second set, with prefix mem, are ports directly connected to the memory chips

through the IOBs of the FPGA: memCe is a chip enable used to switch to power safe mode memories in case of inactivity, memAdv is a control bit used to activate a burst transfer mode, memWe and memBe are respectively write enable and byte enable. memAddr is the memory address for data written via memDataIn or read on port memDataOut. Note the types of the ports and their names: they are mirrored because the memories are external to the chip (i.e. memDataOut has the suffix Out because it is an output for the memory while it is an input port for the IOP).

Third set of ports, with suffix **prog** is composed of 2 control bits used by programming interface to read a configuration file, loaded in advance within the external memory, in order to use it to configure one or more SPs.

4.3.5 SP reconfiguration interface: mainProgInt

Xilinx FPGAs on the SPs are configured in Janus in the SelectMap (8-bit) slave mode [7]. Programming is performed by the IOP, that sequences all needed signals on the configuration interface. The IOP has two independent configuration interfaces, CI, called Channel A and channel B respectively, each of which programs a subset of 8 SPs. During a configuration sequence, an interface configures a subset of all SPs specified by a configuration mask. This feature allows to configure at the same time all SPs sharing the same configuration file (also called bitfile or bitstream). It is assumed in all cases that the bitfile has been loaded onto the IOP memory before the configuration sequence is started.

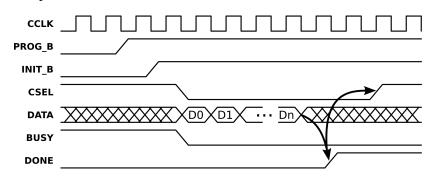


Figure 4.12 – Timing of a typical FPGA configuration sequence.

The configuration sequence is shown in Figure 4.12. In brief

- 1. prog_b is set to 0, to reset and blank the FPGA. At the same time init_b is also set to 0.
- 2. prog_b is set to 1, marking the beginning of the configuration sequence.
- init_b is also released (this is a pulled-up open-drain signal). As init_b goes to 1, the configuration mode is sampled on the mode pins. We always select 8-bit SelectMap slave mode.

- 4. At this point data loading can be started at any time. The FPGA reads the first data byte from the data bus on the first positive clock edge after csel is set to 0. In all cases, we change values on the data bus and on csel only at a falling edge of the configuration clock.
- 5. After all data bytes are read on successive clock transitions, the FPGA releases the **done** signal, signalling that all data has been received. **done** is also an opendrain signal, so it will actually go to logical 1 only if all FPGAs in the set have been configured.
- 6. After a number of clock cycles have elapsed following the transition of done the csel signal can be removed.

After a given set of FPGAs have been configured according to the sequence described above, a further set can be configured starting from point 4.

The CI receives commands from the streamRouter of the IOP. It also receives data from the IOP memory, where it is assumed that bitstream data have been loaded in advance. In normal operation, a command is issued to the CI, followed by a further command to the memory interface, requesting memory to deliver data to the CI. The CI goes through its initialization phase and then starts to deliver data bytes, as soon as they arrive from memory.

A high-level view of the CI, seen as a state-machine in shown in Figure 4.13b.

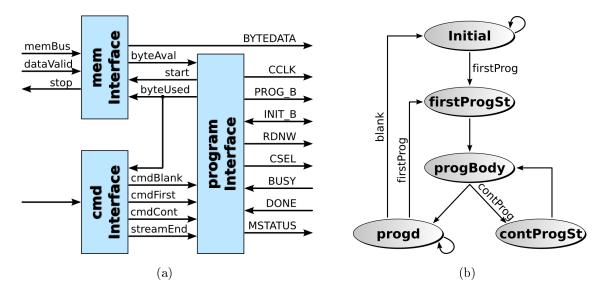


Figure 4.13 – Configuration interface: (a) top level block diagram, (b) high level state machine.

- At reset, the CI state-machine goes to an initial state that blanks the FPGA configuration and delays further configuration (by keeping both prog_b and init_b at 0.

- As the command firstProg is received, the machine goes to the firstProgSt state and properly sequences to 1 the signals described below.
- The machine then goes to the **progBody** states, where all data are sent to the FPGAs. Data bytes are sent as soon as they are received from memory.
- As soon as all data are received, the machine goes to the **progd** state, which is the state where the CI rests during normal operation of the system.
- If more program sequences are necessary on some subset of the FPGAs the statemachine goes to contProgSt where a configuration sequence that does not pulses prog_b and init_b is started. This is triggered by the command contProg.
- Alternatively, all FPGAs can be blanked going again to Initial. This is triggered by the blank command.

At the block diagram level, the CI can be described in terms of 3 main blocks, as shown in Figure 4.13a.

- Block memInterface connects the configuration machine to the memory of the IOP. It receives memory data words (64 bits) flagged by dataValid on memBus and, when appropriate delays data transfer from memory by asserting stop. The block, once started by start, deliver data bytes BYTEDATA and signals data availability on byteAval. In turn, the configuration interface signals that a data has been used (and a new one is needed) over byteUsed.
- Block cmdInterface receives a command stream from the IOP to Host interface. A command contains one of the three available opcodes, a mask that identifies the SP's involved in the configuration sequence, and a count-value, recording information on the length of the datastream. cmdInterface understands the protocol with the Host Interface and forward to the CI the decoded command (on one-hot wires), the SPs mask (not shown in figure) and a streamEnd signal that marks the time-point when the count-value, decreased every time a byteUsed is received, reaches 0.
- Block **Program Interface** is the main engine of the system. It implements the main state-machine and drives all signals for the interface, whose values are derived by the present state of the machine. The block also contains a time-base counter, used to divide the clock frequency of the main system clock by 8, in order to ensure a slow enough frequency for the configuration signals.

4.3.6 SP communication: spInt

When SPs are configured, the entity spInt takes care of communications between host PC and SPs. Behaviour of this entity is simple because it plays the role of stream router from IOP to SPs and vice versa. The interface of the VHDL entity has the

ports with prefix io (i.e. ports for data exchange with IOlink) organized as described in 4.3.4, except for ioDriver that is not required here.

Interface is different if we look at other ports: in this entity in fact for each SP (indicated as SPxy using as labels the Cartesian coordinates of the referred SP) there are the following set of ports: for the communication IOP \rightarrow SPs, dataToSP is the 16 bit data bus starting from IOP and reaching all SPs and dvToSPxy is the data valid associated with data bus. For the communication SP \rightarrow IOP, dataFromSPxy is the data bus coming from SPxy and dvFromSPxy and eotFromSPxy are control bits coding respectively the validity of the word incoming and the end of transmission relative to communication with SPxy.

A bubble diagram of the state machine controlling this data flow is represented in Figure 4.14.

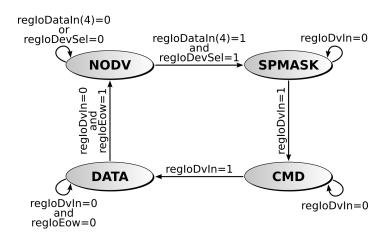
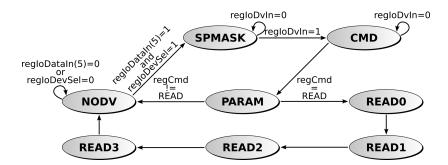


Figure 4.14 – Bubble diagram of the state machine driving IOP interface to/from SPs.

4.3.7 Synchronization device: syncInt

When SPs are running, data buses could be busy and therefore SPs could become unreachable from commands until the end of a run. To avoid the possibility to have a machine out of control and to have the chance to poll the status of each SP we implement on the IOP a synchronization interface, syncInt, that, if the SP firmware implement a compliant interface, can allow us the behaviour described above.

As described in 4.2.2, there are 4 sets of synchronization buses from IOP to SPs; each set is shared by 4 SPs: bus_y is connected to SPxy, with $x, y \in \{0, 1, 2, 3\}$ (i.e. bus_0 is connected to SP00, SP10, SP20, SP30 and so on). Buses from SPs to IOP are point to point: each SP has its own set of signals. Each bus IOP \rightarrow SP is composed of 4 sync signals (spSyncOut_h, with $h \in \{0, 1, 2, 3\}$) plus 2 reset signals (spReset_k, with $k \in \{0, 1\}$); buses SPs \rightarrow IOP are composed of 4 bits.



The state machine is represented in Figure 4.15. This finite state machine is acti-

Figure 4.15 – Bubble diagram of the state machine driving synchronization interface.

vated when a device is selected. First 16 bit word (flagged with a data valid) codes a bitwise mask selecting the SPs with which it exchanging data. The second valid word of the stream is a command that can be READ or PULSE.

If the command is READ the state machine switch to state PARAM, in which the 64 bits of synchronization coming from the 16 SPs are registered in a 64 bits register called snapShot and then reaches 4 states called READ0, READ1, READ2, READ3 in order to send snapShot bits to IOlink, split in words of 16 bits.

If the command is PULSE the state machine passes to state PARAM in which a 16 word containing a parameter is registered in P. Depending on the values of the registered parameter an impulse is sent to the corresponding SP buses following this behaviour:

 $P = 0 \times 0001 \rightarrow \text{pulse a 1 to the line labelled as spReset0}$

 $P = 0 \times 0002 \rightarrow \text{pulse a 1 to the line labelled as spReset1}$

 $P = 0x0010 \rightarrow \text{pulse a 1 to the line labelled as spSyncOutO}$

- $P = 0 \times 0020 \rightarrow \text{pulse a 1 to the line labelled as spSyncOut1}$
- $P = 0 \times 0040 \rightarrow \text{pulse a 1 to the line labelled as spSyncOut2}$
- $P = 0 \times 0080 \rightarrow \text{pulse a 1 to the line labelled as spSyncOut3}$

If the SPs implement a synchronization interface able to respond/react to this impulses, the PULSE command can generate different types of reset or different requests for status information about the running program on the SPs.

4.4 SP firmwares for Janus test

The development of a complex system such as Janus needs a test phase of the new hardware in order to verify and stress all components. FPGA, but also I/O interfaces, memories, double data rate and single data rate communication channels were tested with a set of incremental checks having the aim of discovering possible assembly faults. During my PhD I spent several months to develop this test infrastructure that I and

some other people of the Janus collaboration use during December 2007 to check all the 19 boards built by Eurotech for the Janus collaboration. The effectiveness of this test suite can be assessed by the fact that in less than one month all Janus boards were successfully put into operation.

The incremental tests developed, from the simpler to the more complex, are:

01_uart_regs

Performs a reset via the serial interface; reads two configuration registers; writes a configuration register; re-reads the written register.

Coverage: checks components driving serial link.

02_uart_int_mem

Performs a set of write and read operations on an FPGA internal memory via serial link using sequential patterns and random patterns.

Coverage: check serial link and some FPGA structures.

03_uart_ext_mem

Performs a large set of write and read operations on the external staging memory via serial link using sequential patterns and random patterns.

Coverage: checks the stability of the serial link and the integrity of the external memory.

04_gbit_regs

Reads 29 registers of the PHY driving Gigabit link and check their values. *Coverage:* checks the PHY integrity.

05_gbit_int_mem

Perform a set of write and read operations on an FPGA internal memory via Gigabit link using sequential patterns and random patterns.

Coverage: check Gigabit link and some FPGA structures.

06_gbit_ext_mem

Performs a large set of write and read operations on the external staging memory via Gigabit link using sequential patterns and random patterns.

Coverage: checks the stability of the Gigabit link and the integrity of the external memory.

07_sync

First performs configuration of all SPs via IOP configuration interface with a basic firmware, then IOP pulses a value on synchronization lines of each SP; each SP responds with the old value of the sync register x-ored with the value pulsed from IOP. When IOP sends a synchronization reset (bit 0) all synchronization lines from SP to IOP are pulled low; when IOP sends a synchronization reset (bit 1) all synchronization lines from SP to IOP are pulled high. *Coverage:* checks SPs configuration and synchronization interfaces.

08_iop_sp_comm

Resets the configuration firmware of all SPs, re-configures them and then checks single ended lines IOP \leftrightarrow SP writing data coming from host PC in an internal memory of the SPs. Re-reads data and checks their integrity.

Coverage: checks SPs configuration and all data and controls lines IOP \leftrightarrow SP.

09_loop_2_sdr

Resets the configuration firmware of all SPs and re-configures them. Initializes with different values two memories of a pair of neighbours SP_A and SP_B. Following a start command, two initialized SPs exchange between each other their memory contents. The host PC polls the status via synchronization signals. When data transfer SP_A \leftrightarrow SP_B ends, host PC dumps the memory contents and cross checks that data of SP_A are stored onto memory within SP_B and vice versa. *Coverage:* checks SPs configuration and all differential pairs SP \leftrightarrow SP in all directions.

10_loop_2_ddr

The same of $09_1oop_2_sdr$ but differential pairs SP \leftrightarrow SP works in double data rate mode.

Coverage: check stability of double data rate on differential links.

11_loop_full_board

Performs a test of data transfer among all possible ring of SPs on a Janus board. *Coverage:* checks stability in case of simultaneous load of all communication lines on a board.

12_random_wheel

First intensive test in terms of FPGA resources: we implement 37 shift registers each generating 32 random numbers per clock cycle. We store them in some buffers and we check them after a fixed number of clock cycles. This firmware requires $\sim 86\%$ of the FPGA logic and $\sim 38\%$ of the embedded RAM. Moreover this test implement a system to check temperature of the FPGA.

Coverage: checks a complex situation of medium-high load.

13_sg640

Runs the first real spin glass simulation performing 640 spin update per clock cycle on a lattice of 80^3 spins.

Coverage: checks a typical work load.

4.5 Engineering problems

Janus is a complex system using non challenging technology: the clock frequencies are low (62.5 MHz or 125 MHz), most of the data exchanges are synchronous, the communication interfaces use standard and stable protocols and the integration level is not too high.

Because of this, its development was straightforward and no really substantial problems were encountered. The only area in which non trivial engineering problems had to be faced was associated with power issues. In short, our FPGAs required more power than we were able to supply, and this fact triggered a sequence of additional problems that were partly solved, partly swept under the carpet and partly non solved. As a consequence, today we are in the odd situation in which Janus performances are limited by power problems. In this section we briefly summarize these problems and how we handled them.

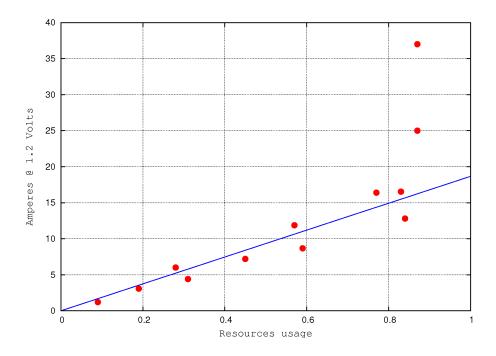


Figure 4.16 – Current absorbed by a Janus computing core (measured in Amps) as a function of logic resource usage.

As mentioned above we found a critical threshold for the current absorbed by the FPGA in the cases in which the implemented design uses more than $\sim 85\%$ of the FPGA logic resources (also called *slices* using the Xilinx lexicon). The graph in Figure 4.16 shows the current absorbed by the computing cores of each node of a Janus board (measured in Amps) as a function of logic resource usage. It is interesting to note that the measurements have been taken with different firmwares, but all of them follow basically a linear fitting curve except for the last two points corresponding to the largest resource usage of our firmwares. In the case of these firmwares it is important to note that a careful (hand made) placement of the resources within the FPGA has a non negligible impact on the current absorption.

Three additional problems arose as a consequence of the power problem, and we had to find reasonable solutions to them.

- Current absorption peak implies a power supply problem that in some cases leads to the FPGA switching off. Fortunately we discovered this problem during prototype tests so we solved it with minor modifications in the power distribution circuitries (we added one more power converter from 2.5 V to 1.2 V on the SP daughter card and more powerful power converters from 48 V to 2.5 V on the PB).
- Heat removal becomes a serious problem in the case of high density design. Each Janus box houses indeed a set of 10 fans and each FPGA houses an heat sink but because of some non optimized design choices (basically related to the placement of the fans) this set up of the box does not allow an high efficient heat removal. Figure 4.17 shows the temperature of the 16 SPs during an intensive run; it is evi-

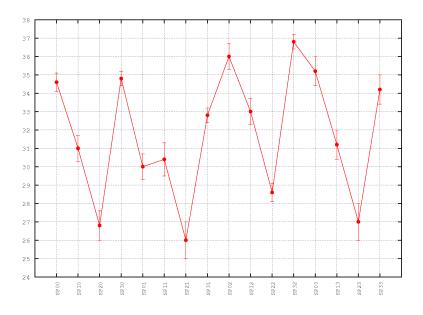


Figure 4.17 – Map of the temperature within a Janus box.

dent that temperature among modules is not homogeneous. It is remarkable that the heat problem triggers a vicious circle because temperature increase reduces the efficiency of the power supply and therefore these temperatures problems are strongly related to the discussion above. Fortunately this design imperfections do not affect the performance of the whole system and they do not disadvantage our simulations.

- We discovered that high supply currents (typical of large designs when they have a large switching rate) bias clock stability. Figure 4.18 shows in details



Chapter 4. Architectural details of Janus

Figure 4.18 – Clocks during a run using an high density design: duty cycle of the black signal is clearly asymmetric.

a scope snapshot of two clocks observed on two SPs, one running an intensive firmware and one running a low density design. The clock becomes asymmetric and the duty cycle distortion generates errors in the circuitry that uses the double data rate technology described in 4.3.2. It is remarkable that this asymmetric behaviour stops immediately when the run ends: this is the reason leading us to think that this problem is weakly connected with current absorption and power supply. As we use double data rate only to double the bandwidth between IOP and SPs we have no effective drawbacks because of this: we deactivate the double data rate communications during the run time period, in which IO data flow is relative small, ensuring therefore the correctness of the data transaction.

The engineering problems described above de facto limit the performance of our system. In principle, by re-engineering part of the system we might solve them; we expect that, by doing so, clock frequencies approximately two times higher could be supported by the system, with a corresponding performance increase. We have however decided not to put these improvements in practice, partly because of their estimated costs and – even more importantly – because we do not want to delay the use of the machine for its physics program. Obviously, if a new Janus generation is developed in the future, we will carefully apply the lessons learned.

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But I still haven't found what I'm looking for...

5 Performance and results

In this chapter I will present some relevant physics results obtained with the Janus supercomputer. In first section I summarize some physics concepts useful for the following discussions. In the second section I describe the features of the first run performed in spring 2008: special care is provided to the performance and the cost of the Janus system. Third section resumes concepts and results related with the first two important runs of Janus: the simulation of the Edwards-Anderson spin glass model and the simulation of the Potts model with four states.

Fundamental references for this chapter will be three of the most significant articles written thanks to the data produced by Janus: [1, 2, 3].

5.1 Useful concepts

5.1.1 About the equilibrium

In section 2.2.5 we look at the Monte Carlo methods as a way to find an *equilibrium* configuration for a model systems, but in some cases is interesting to discover the behaviour of systems which are *out of equilibrium* using the same tools introduced in previous chapter (see for instance 2.2.4). In studying these out-of-equilibrium systems we are commonly interested in one of two things. Either we want to know how the system relaxes from an initial state to equilibrium at a particular temperature, or we are studying a system which never reaches equilibrium because it has a driving force which continually pushes it into states which are far from equilibrium.

The statistical mechanics of out-of-equilibrium systems is a less well-developed field of study than that for equilibrium systems, and there is no one general framework to guide our calculation such as we had in the equilibrium case. As the subject stands at the moment at least, each system must be considered independently. In many cases the mathematical hurdles to formulating an accurate analytic theory of a systems behaviour are formidable.

For these reasons we decided to use the first run of the Janus system to investigate out-of-equilibrium dynamics (see Sections 5.2 and 5.4 for more details).

5.1.2 Correlation

Even if a system is not ordered, there will in general be microscopic regions in the material in which the characteristics of the material are correlated. Correlations are generally measured through the determination of a two-point correlation function

$$\Gamma(r) = \langle \rho(0) \cdot \rho(r) \rangle \tag{5.1.1}$$

where r is the spatial distance and ρ is the quantity whose correlation is being measured.

Below the phase transition, $\Gamma(r)$ becomes large for all values of r, while it rapidly decays to zero well above the critical temperature. The function $\Gamma(r)$ takes into account all the contributions to the correlation of ρ (for example, ρ can be the state of a spin), even those due to external fields. In order to measure only the part of correlation that is due to internal fluctuations of the observable (e.g. only to the interactions between spins) we define the *connected correlation function*

$$\Gamma_C(r) = \langle \rho(0) \cdot \rho(r) \rangle - |\langle \rho \rangle|^2$$
(5.1.2)

where we have discounted the overall alignment of the observable ρ . It is found that, for $T = T_C$ but close to the critical temperature, and for $r \to \infty$, this term goes as

$$\Gamma_C(r) \approx r^{(d-1)/2} e^{-r/\xi} \tag{5.1.3}$$

where d is the space dimensionality and ξ is a characteristic length of the system, known as *correlation length* associated to the observable ρ .

When approaching the transition temperature T_C from above, the interactions within the system tend to become more and more relevant, and at the same time it grows also the correlation between spatially well separated points of the system. The correlation length ξ corresponds to the average distance at which different parts of the system present a correlation on the values of a given observable ρ : that is, $\langle \rho(r_i)\rho(r_j)\rangle$ is significant for $|r_i - r_j| < \xi$ but tends to zero for larger distances. Typically ξ diverges at T_C . Equation 5.1.3 implies that the order parameter can normally fluctuate in blocks of sizes up to ξ , while fluctuations of larger size are rather improbable [4, 5]. It is also possible to consider correlations that are time-dependent. In this case the system is prepared in a given initial configuration and then the simulation is run for a time t_w (called *waiting time*). After the time t_w the configuration $\sigma(t_w)$ is stored. From now on after each MC step the following correlation function is measured:

$$C(t, t_w) = \frac{1}{N} \sum_{i} \overline{\langle \sigma_i(t+t_w)\sigma_i(t_w) \rangle}$$
(5.1.4)

where $\langle \dots \rangle$ means a thermal average (i.e. an average over different realization of the thermal noise, but the same initial configuration) and the bar means an average over different realization of the bond-disorder [6].

5.1.3 Order parameters: magnetization and overlap

Phase transitions involve an abrupt change of some macroscopic properties of the system, generally originating from a change in the microscopic structure of the system under investigation. In the simplest cases this can be signaled by the observation of a properly defined *order parameter*, i.e some property of the system which is identically zero in the disordered phase and non-zero in the ordered phase. This change of behavior when crossing the transition temperature, makes the order parameter a perfect marker to spot the change of phase in the system when the transition occurs. An order parameter is thus related to the global symmetries of the system. It is not possible to give a generic definition of an order parameter, and it typically has to be defined in a different way for each physical system of interest. A considerable ingenuity is often necessary to find a good order parameter to characterize a certain phase transition.

An example of order parameter in ferromagnetic materials (i.e. the Ising model) is the magnetization $M = 1/V \sum_{i} m_i = 1/V \sum_{i} \langle s_i \rangle$, where s_i is the value of the magnetic moment at position *i*. The value of *M* changes from non-zero below the phase transition, to zero above the critical temperature, and marks the change from the ferromagnetic to the paramagnetic states.

We have seen that a characteristic of spin glasses is the lack of a uniform ordering, substituted by a frozen disordered configuration. This phase seems to possess a nonzero local spontaneous magnetization $m_i = \langle s_i \rangle$, though the average magnetization and any staggered magnetization vanish because of the absence of regularity in the configuration. Obviously neither of these observables can be used as an order parameter. It was noted by Edwards and Anderson that a correct description of the spin glass phase should reflect the lack of global orientation spin order in the frozen phase (despite the local magnetization of the spins).

The first proper spin glass order parameter was proposed in the original EA work,

and is thus known as the Edwards-Anderson parameter:

$$q_{EA} = \frac{1}{N} \sum_{i=1}^{N} \overline{\langle s_i \rangle^2} \quad . \tag{5.1.5}$$

This term vanishes in the paramagnetic phase, but it is non-zero if the local magnetizations m_i are nonzero: this makes it a good order parameter for the transition from paramagnetic to spin glass phase. The quantity q_{EA} is actually a particular case of a more general quantity called *overlap*, a statistical mechanics tool used to measure the similarity of two different configurations or replicas of a system. Given two spin configurations a and b we define their mutual overlap as:

$$q_{ab} = \frac{1}{N} \sum_{i=1}^{N} s_i^a \cdot s_i^b \quad , \tag{5.1.6}$$

and it tells us how correlated a and b are. For example, for two replicas of an Ising spins model we find the following possible values:

$$q_{ab} = \begin{cases} 1 & \text{if } a \text{ and } b \text{ are completely correlated;} \\ -1 & \text{if } a \text{ and } b \text{ are anti-correlated;} \\ 0 & \text{if } a \text{ and } b \text{ are completely uncorrelated;} \end{cases}$$
(5.1.7)

5.2 First run details

We have followed the Monte Carlo dynamics of three cubic 80^3 spins lattices (all initialized with random values, mimicking initial high-temperature conditions) at three different temperatures below $T_c = 1.13$ [7], i.e.

$$T_1 = 0.8 = 0.70 T_c ,$$

$$T_2 = 0.7 = 0.62 T_c ,$$

$$T_3 = 0.6 = 0.53 T_c .$$

(5.2.1)

We have performed 10^{11} Monte Carlo iterations for 96 independent samples at each temperature (64 for T_2). For comparison, the previous largest simulation of the same model runs for 10^9 Monte Carlo iterations, for one temperature, 16 samples and a smaller lattice of 60^3 sites [8]. The present simulation is altogether $\simeq 2000$ times more compute intensive, allowing to come much closer to experimental time scales.

An interesting feature of the system is the size of the coherent domains. A clever trick to measure this quantity is to simulate *two* copies of the system, $\{\sigma_i^{(1)}, \sigma_i^{(2)}\}$ (evolving with the same set of couplings) for each samples. We define the overlap field at each lattice site as: $q_i = \sigma_i^{(1)} \cdot \sigma_i^{(2)}$. The key observation is that, if lattice points *i* and *j* belong to the same coherent domain, q_i and q_j tend to have the same

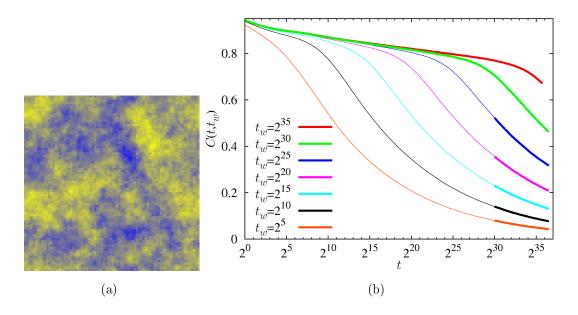


Figure 5.1 – (a) Plot of the overlap field of a sample after 2^{36} Monte Carlo steps at T = 0.8, corresponding to $\simeq 0.1$ sec. The pixel brightness at point (x, y) is the sum of the positive overlaps for all values of the corresponding z coordinate. (b) Time correlation function $C(t, t_w)$, at T = 0.6 for several values of t_w . For comparison, the longest simulation available before the present one [8] only explored a time window shown by thin lines; thick lines show the extended time window made possible by Janus (note that a log-scale in used for the x-axis).

sign. Figure 5.1(a) shows a picture of the system at T_1 after a huge number of Monte Carlo steps. The typical domain size (quantitatively estimated from the data shown in the figure) is close to 15 lattice spacings. Previous work [8] reached around 8 lattice spacing.

A very important physical quantity is the time correlation function $C(t, t_w)$ of two spin configurations, one at times t_w , the so called waiting time, the other at a later time, $t + t_w$. Specifically, (with the bar we mean the average over the samples)

$$C(t, t_w) = \overline{\frac{1}{L^3} \sum_{i=1}^{L^3} \left[\sigma_i(t+t_w)\sigma_i(t_w)\right]}.$$
 (5.2.2)

Statistical errors can be computed from the sample to sample fluctuations (hence the importance of simulating some tenths of samples).

The function $C(t, t_w)$ tells us about the memory that the spin configuration at time $t + t_w$ retains from the spin configuration at the waiting time, t_w . Of course, $C(t = 0, t_w) = 1$ (perfect memory), while $C(t, t_w)$ vanishes for large t (i.e. no memory of what happened at t_w). Once thermal equilibrium is reached, $C(t, t_w)$ should be independent of t_w . Figure 5.1(b) shows that this is far from being our case. A detailed analysis of the behavior of $C(t, t_w)$ in our simulation has provided clear indications to settle a long-standing debate on the best theoretical picture of this process [2]. The simulation was performed in March 2008 on our 16 core Janus system. 6+6 cores have been used for T = 0.8 and 0.6 and 4 cores for T = 0.7. The simulation has executed almost continuously for approximately 25 days (20 days for simulation at T = 0.7), apart from a 2 hours outage caused by snow-storm. A further Janus core was used for checks, extensively verifying the reproducibility of selected segments of the run. Simulation results (that is spin glass configuration at appropriate Monte Carlo times) were written onto disks attached to the Janus-host cluster. All in all about 4 TBytes of disk space have been used.

All physically relevant averages (such as the overlap, $C(t, t_w)$, and many others) have been computed on line on the Janus-host systems. Averages were computed as soon as data became available: these tasks are not a computational bottleneck in our case: just one or two host PCs are necessary to keep up with incoming data.

5.3 Janus performance

In this section, we assess the performance of our system in two different ways:

- the number of effective operations per second.
- the speed-up factor (mostly in wall clock time, but also on other relevant metrics) with respect to processor clusters or "arbitrary" size (i.e., assuming that, for the given simulation, the optimal number of processor is actually deployed).

Let us first consider the number of effective operations performed by second. At each clock cycle (clock frequency is 62.5 MHz), each SP processor updates 800 spins. On a traditional architecture, this could require at least the following instructions for each spin:

- 6 loads
- 1 sum (32 bits)
- 2 xor (32 bits)
- 6 sum (3 bits)
- 6 xor (3 bits)
- 1 load LUT
- 1 comp (32 bits)
- 6 updates of address pointers (1 mult, 1 sum)
- 1 store
- jump condition

We want to count only algorithm-relevant operation, so we do *not* count all load-/store and address instructions; also, we count the 6 short xor and sum operations as one each. We end up with 7 equivalent instruction for each spin update (shown in italic in the list above). This translates into a processing power of $7 \times 800 \times 62.5 \times 10^6$ operations per second, that is 350.0 Giga-ops. Our system operates 256 processors in parallel, so it performs at the level of 89.6 Tera-ops.

From this (accurately measured) value and the cost figures provided in [1] we obtain our price performance ratio of $6.56 \notin$ / Giga-ops. This figure grows to $7.98 \notin$ / Gigaops if we include the purchase costs of the items donated by our industrial partner. Using the $\neq \in$ exchange rate of 1.566 prevailing on April 4th, 2008, we obtain our final figure of 10.27 \neq Giga-ops (or 12.50 \neq Giga-ops if the costs of items donated to the project are included).

Note that our system is also extremely energy efficient at $\simeq 8.75$ Giga-ops / W (for comparison, the top entry in the Green500 list is rated at 357.23 MFlops/W).

We now compare our performances to those obtained on commercial CPUs, where spin-glass simulations have been performed so far. This analysis is mostly intended to show the impact that our machine is going to have on the spin-glass community.

As stated in Section 2.3.2, the SMSC is the elective technique to handle Monte Carlo simulation of large spin glass systems. The AMSC is only useful when dealing with smaller lattice sizes, where equilibrium properties are investigated instead of slow dynamics behavior, and a large set of sample statistics is needed.

The figures given below refer to carefully programmed routines written in C, running on an Intel[®] Core 2 DUO 64 bit CPU 2.4 GHz. In the multi-spin coding implementation we also had some advantage in using C compilers supporting 128 bit extended integer data types on 64 bit architectures (as the gcc 4.x releases do). A further warning is associated to the fact that multi-spin coding efficiency on traditional CPUs critically depend on whether the size of the simulated lattice is an integer multiple or divisor of the CPU basic data word. While we cannot claim that our implementations are optimal, our judgment is that further optimization would not improve performance by factors larger than 2 or 3, marginally affecting the qualitative picture that we are outlining.

Our SMSC routine performs at 7.0 ns/spin average update rate when simulating an L = 80 system. The AMSC routines available to us allows a 0.77 ns/spin update rate, almost independently of the lattice size on modern CPUs equipped with large cache memories. By contrast, our FPGA implementation performs at 0.020 ns/spin for the L = 80.

Table 5.1 reports the total CPU time, the wall clock time and an estimate of the total energy needed to accomplish the task of performing the 10^{11} Monte Carlo steps presented in the sections above, in various cases of different sample statistics (256 samples is the case relevant for physics, as discussed above). The three columns refer respectively to Janus, to PCs using AMSC, and to PCs using SMSC. We assumed,

	Janus	1 PC AMSC	1 PC SMSC
samples	1	1(128)	1
Wall clock time	24 d	310 y	24 y
Accumulated CPU time	24 d	310 y	24 y
Energy	$85 \mathrm{~MJ}$	$900 \mathrm{GJ}$	$70 \mathrm{GJ}$
	Janus	2 PCs AMSC	256 PC SMSC
samples	256	256	256
samples Wall clock time	256 24 d	256 310 y	256 24 y
-			

Chapter 5. Performance and results

Table 5.1 – Processing time for 10^{11} Monte Carlo steps on a 3D lattice of 80^3 points.

when estimating performances scaling with available resources, that Janus is limited to 256 nodes (SPs), while for PCs we assumed availability of infinite resources (a cluster of the optimal number of processor is assumed to be available). We consider a 100 W power consumption for a standard PC and round the figure for Janus at 40 W.

Comments on figures in the table may be summarized by the following considerations:

- in spite of its high "time per spin" efficiency, AMSC would have not permitted at all to accomplish our task;
- SMSC is the best choice on standard PCs; however estimated wall clock time in this case is long enough to discourage even the most patient researcher.
- for this application, commercial CPUs largely suffer from having a fixed architectures, as SMSC would greatly benefit of a flexible architecture as the Janus' one: we have checked that the time ratios for PCs with SMSC and Janus increase by a factor four when passing from L = 64 systems to L = 80 ones.
- as the table shows, the wall clock time in a PC cluster is largely independent of the size of the cluster itself, as scaling quickly saturates because a small part of the available parallelism is instantiated;
- our simulation on Janus corresponds to $\simeq 10^4$ CPU-years.
- our simulation is extremely energy efficient: using the numbers of Table 5.1 and standard figures in the oil industry [9], we estimate that Janus has used approximately 15 barrels of oil for our simulation campaign, to be compared with $\simeq 12000$ that would be needed if PCs were used.

5.4 Physics results overview

After its commissioning, acceptance and reliability tests in April 2008, the large Janus installation in Zaragoza was immediately used for two large scale simulations.

In the first simulation (whose results were also used for our submission to the 2008 Gordon Bell Prize), we studied the long term (in Monte Carlo time) evolution of a large (80^3) Edwards-Anderson spin glass at several temperatures clearly below the critical one. In this type of simulation, the physical focus is not on measuring physical observables at the statistical equilibrium (reaching equilibrium for such a large system is still out of question even for Janus); rather, we want to pinpoint several specific features of how the system drifts within its configuration space as it tries to move toward more energetically favourable configurations. This is a relevant subject of study, since several conjectures have been made, so it is important to compare with experimental data.

In the second large simulation, we studied the 4-state Potts model (in 3D). In this case, a much smaller lattice (16^3) was used and a serious attempt has been done to bring the system into statistical equilibrium. In this case the ultimate goal of the simulation was to characterize the structure and to measure the main parameters of the phase transition of the system.

I describe in more details both simulations and their main physical results in the following.

5.4.1 Non-equilibrium dynamics of a large EA spin glass

The results discussed in this section have been published in [1], where special attention is given to the computational aspects of the simulation and in [2, 10] where the results of the physical analysis are considered in details.

Below their glass temperature, experimental spin glasses are perennially out of equilibrium. The understanding of their sophisticated dynamical behavior is a long standing challenge both to theoretical and to experimental physics. On the simulation side, it is only possible to bring to statistical equilibrium very small systems, so if we study these small systems, we may expect that our findings are strongly affected by finite-size effects. On the other hand, several theoretical models try to highlight correlations between the structure of the equilibrium configurations of a large spin glass and the dynamics according to which that system slowly drifts toward equilibrium. If we study numerically a large spin glass system out of equilibrium we may therefore hope to extrapolate information on its equilibrium properties. In this sense, the two approaches are complementary ways to study these systems.

The standard (experimental and numerical) approach to the study of an out-ofequilibrium spin glass is the so-called direct quench. In these experiments, the spin glass is cooled as fast as possible to the working temperature below the critical one, $T < T_c$. It is let to equilibrate for a waiting time, t_w , its properties to be probed at a later time, $t + t_w$. For instance one may cool the spin glass in the presence of an external field, which is switched off at time t_w . The so called thermoremanent magnetization decays with time, but the larger t_w is, the slower the decay.

In fact, there is strong evidence that, if the cooling is fast enough, the thermoremanent magnetization depends upon t and t_w only through the combination t/t_w , for a very large time window.

The time evolution is believed to be caused by the growth of coherent spatial domains. Great importance is ascribed to the size of these domains. Domain size is characterized by the coherence length $\xi(t_w)$, which can be measured experimentally and has to be correctly estimated from numerical simulation. Therefore, in a large simulation, one has to measure observables that lead to the definition and evaluation of $\xi(t_w)$. This is the main raw-data of the simulation, that can then be used to compare with theoretical expectation. In the rest of this section I focus on this work.

We simulated the dynamics of the Edwards Anderson model, defined in (2.1.5). The system was simulated on a lattice of linear size L = 80 and on a smaller lattice with L = 40, at several temperatures below and above the critical temperatures (which is not known with high precision but was estimated in previous works to be close to $T_c \simeq 1.1$). Several samples of the systems were simulated and a huge number of Monte Carlo steps (up to 10^{11}) was performed. The parameters of our simulation are summarized in table 5.2.

L	Т	MC steps	N_s
80	0.6	10^{11}	96
80	0.7	10^{11}	63
80	0.8	10^{11}	96
80*	0.9	$2.8 imes 10^{10}$	32
80	1.1	$4.2 imes 10^9$	32
80*	1.15	2.8×10^{10}	32
80*	0.7	10^{10}	768
40	0.8	2.2×10^9	2218

Table 5.2 – Parameters of our simulations. The overall wall-clock time needed was less than six weeks. We highlight with * the simulations performed after completion of [3]. Recall that we take the critical temperature from [11], $T_c = 1.109(10)$. The full analysis of spin configurations was performed off-line.

We wrote to disk the spin configurations at all times of the form $[2^{i/4}] + [2^{j/4}]$, with integer *i* and *j* (the square brackets stand for the integer part). Hence, our *t* and t_w are of the form $[2^{i/4}]$. These spin configurations are the starting point for any analysis. It is common practice, followed in this work, to define real replicas. These are two statistically independent systems, $\{\sigma_x^{(1)}\}\$ and $\{\sigma_x^{(2)}\}\$, evolving in time with the very same set of couplings. Their overlap field at time t_w is:

$$q_x(t_w) = \sigma_x^{(1)}(t_w)\sigma_x^{(2)}(t_w)$$
(5.4.1)

The overlap is the basic quantity from which most observables are derived. For instance, the spin glass order parameter is

$$q(t_w) = \frac{1}{N} \sum_{x} q_x(t_w)$$
 (5.4.2)

The mean value of $q(t_w)$ over many samples vanishes in the non-equilibrium regime, when the system size is much larger than the coherence length $\xi(t_w)$. So the spin glass susceptibility,

$$\chi_{SG}(t_w) = N\overline{q^2(t_w)} \tag{5.4.3}$$

that steadily grows with the size of the coherent domains, is a first approximate way to measure the correlation length. More accurate estimation of the correlation length are based on so called single time correlation functions

$$C_4(r, t_w) = \frac{1}{N} \overline{\sum_x q_x(t_w) q_{x+r}(t_w)}$$
(5.4.4)

The long distance decay of $C_4(r, t_w)$ can be written according to the following functional form:

$$C_4((r, t_w) = \frac{1}{r^a} f(r/\xi(t_w))$$
(5.4.5)

In other terms, if we are able to estimate correctly ξ we are in a position to reliably understand (from numerical or experimental data) the functional behaviour of f, thereby clarifying the approach to equilibrium of the system.

In order to measure ξ in a model-independent way, we compute the integrals

$$I_k(t_w) = \int_0^\infty dr r^k C_4(r, t_w)$$
 (5.4.6)

If we take (5.4.5) into account, we see that

$$\xi_{k,k+1} = \frac{I_{k+1}(t_w)}{I_k(t_w)} \propto \xi(t_w)$$
(5.4.7)

In principle this equation is valid for any k. In practice, for small k there is a systematic error because (5.4.5) is only valid for large values of r. However, if we increase k (thereby suppressing the contribution of small r to the integral) we meet larger statistical error. In practice, the best trade-off is for k = 1 (or, equivalently, $\xi_{1,2}$). This is shown in Figure 5.2a. Using this technique, we are able to measure

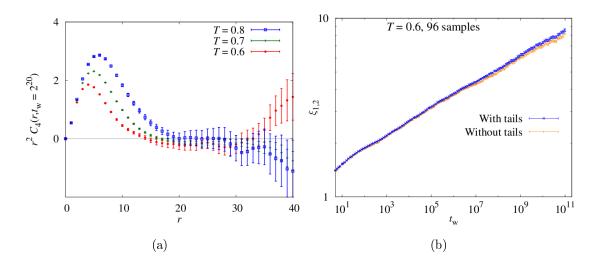


Figure 5.2 – (a) The spatial autocorrelation of the overlap field for $t_w = 220$ and three subcritical temperatures, as computed in our L = 80 lattice. (b) Result of computing $\xi_{1,2}$ in two different ways for our 96 samples at T = 0.6. In the orange curve we stop the integration at the cutoff point where relative error of C_4 is greater than one third. In the blue curve we estimate the contribution of the tail from that point on extrapolating with another fit. The difference is small, but with the second method the power law behavior of $\xi_{1,2}(t_w)$ lasts longer.

 $\xi_{1,2}(t_w)$ (see Figure 5.2b). To check the reliability of our result, we can compare with the estimate for ξ that we obtain from 5.4.3. The comparison is given in Figure 5.2b, showing that the two results are consistent but that the measurement based on ratios of $I_k(t_w)$ has much smaller errors.

Based on this preliminary analysis, a large number of physics results have been derived for our spin glass system. Since the details of the analysis are very technical in nature, we address the interested reader to the original papers.

5.4.2 The 4-state Potts model and its phase structure

The results discussed in this section have been published in [3, 12]

We have simulated three dimensional cubic lattices with linear sizes L = 4, 6, 8 and 16 (I explain below why simulation of large and small lattices are needed). Because spin-glass simulations have very long relaxation times, we used the parallel tempering (PT) algorithm (described earlier in this thesis) to speed up the dynamical process that brings the system to thermal equilibrium and eventually explores it. Physical quantities are only measured after the system has been brought to equilibrium.

The Monte Carlo dynamics uses single-spin updates and temperature swaps. The single-spin updates are carried out with a sequential heat bath (HB) algorithm. We define a Monte Carlo sweep (MCS) as N sequential trial updates of the HB algorithm

(i.e. every spin in the lattice undergoes a trial update once). The PT algorithm (applied to a given realization of the quenched disorder, that we will call a sample) is based on simulating a number of copies of the system with different values of the temperature but the same interactions (the same set of J_{ij}). Exchanging the temperature of two copies with adjacent temperatures with a probability that respects the detailed balance condition is the crucial mechanism of PT. The result is that each copy of the system drifts in the whole allowed temperature range (that has been decided a priori). When a copy is at a high temperature it equilibrates fast and so each time it descends to low temperature it is likely to be in a different valuey in the energy landscape.

The simulation of the smaller lattices, with L = 4 and 6, was performed on standard PCs, while Janus was used to simulate the larger lattices of linear size L = 8 and L = 16. For this specific simulation, one Janus SP processor (one FPGA) is about 10^3 times faster that an Intel Core2DuoTM processor [13]. Janus has allowed us to thermalize a large number of samples for bigger sizes than would have been feasible on a standard computer. The computational effort behind our analysis amounts to approximately 6 years CPU time on a 2.4 GHz Intel(R) Core2Duo(TM) processors for L = 8 and thousands of CPU-years for L = 16.

Data input and output is a critical issue for Janus performance, so we had to carefully choose how often to read configuration data; in general, we end up taking fewer measurements than in simulations on a traditional PC. Having fewer (but less correlated) measurements does not affect the quality of our results. We read and analyze values of physical observables every 2×10^5 MCS. On the larger lattices, we perform a PT step every 10 MCS while on the smaller lattices this value is 5. In a standard computer the PT algorithm takes a negligible amount of time, compared to a whole MCS. However, in Janus the clock cycles needed by one PT step are more than those needed for a MCS. For this reason we chose to increase the number of MCS between two PT steps. However, this number should not be too large, as we do not want to negatively affect the PT efficiency. A preliminary analysis has been carried over to test how the PT parameter would affect the simulation results, and we have selected a value that seems to be well optimized (see Table 5.3).

When performing such a large simulation, for which most of the expected results live in a *unknown space* it is important to perform reasonable sanity checks on the results of the simulation.

A standard test that a set of configurations are extracted from a thermalized sample takes a given physical quantity and averages (first over the thermal noise and then over the quenched disorder) over logarithmically increasing time windows. Equilibrium is reached when successive values converge. We emphasize that it is crucial for time to be plotted on logarithmic scale.

A typical quantity that is analyzed in this way is the correlation length of the

L	J	N_{samples}	MCS	$[eta_{\min},eta_{\max}]$	N_{β}	$N_{\rm HB}$	$N_{ m m}$
4	:	1000	3.2×10^5	[2.0, 6.0]	9	5	10^{3}
6	;	1000	8×10^5	[2.5, 5.0]	7	5	10^{3}
8	,	1000	2×10^8	[2.7, 4.2]	16	10	2×10^5
16	6	1000	8×10^9	[1.7, 4.1]	32	10	2×10^5

Table 5.3 – For each lattice size we show the number of disorder samples that we have analyzed, the number of MCS per sample, the range of simulated inverse temperatures $\beta = 1/T$, the number of (uniformly distributed) β values used for PT, the number of MCS performed between two PT steps (N_{HB}), and the number of MCS between measurements (N_m).

spin-glass order parameter ξ . ξ is obtained by a rather complex set of mathematical operations, that we describe here (albeit quickly) because this is a central quantity in the simulation. One starts with a definition of the Potts model with p = 4 states in the so-called simplex representation: the possible spin states are associated to one of the p unit vectors S_a pointing to the corners of a hyper-tetrahedron in a p-1 (3, in our case) dimensional space. The following equality holds:

$$S_a \cdot S_b = \frac{p\delta_{ab} - 1}{p - 1}$$
(5.4.8)

so the Hamiltionian

$$H = -\sum_{\langle ij \rangle} I_{ij} \delta_{s_i, s_j} \tag{5.4.9}$$

can be written in term of $S_a \cdot S_b$.

Using this representation for the spins, we define the overlap between two replicas $(S_i^{(1)} \text{ and } S_i^{(2)})$ of the same system in Fourier-space as

$$q^{\mu\nu}(k) = 1/N \sum_{i} S_{i}^{(1)\mu} S_{i}^{(2)\nu} e^{ikR_{i}}$$
(5.4.10)

where S_i^{μ} is the *i*-component of the spin along direction μ . From this quantity, the spin-glass susceptibility is defined as

$$\chi_{q(k)} = N \sum_{\mu,\nu} \overline{\langle |q^{\mu\nu}(k)|^2 \rangle}$$
(5.4.11)

and, finally, the correlation length is derived as a function of χ evaluated at two specific points in Fourier space $(k_m = (2\pi/L, 0, 0))$:

$$\xi = \frac{1}{2sin(k_m/2)} \left(\frac{\chi_q(0)}{\chi_q(k_m)} - 1\right)^{1/2}$$
(5.4.12)

As a check of thermalization, ξ is plotted in Figure 5.3 at the lowest simulated temperature (the hardest case for thermalization). We see that the values of the correlation

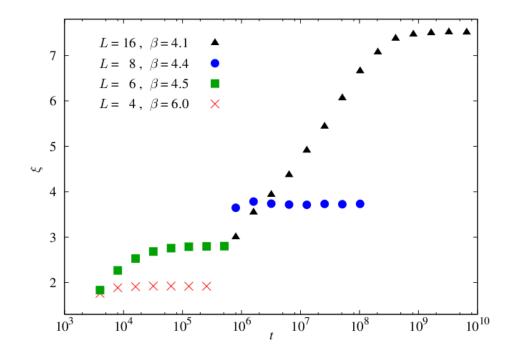


Figure 5.3 – A thermalization test. We show the behavior of the time dependent spin glass correlation length as a function of Monte Carlo time. We have averaged the correlation length using a logarithmic binning procedure. We show data for the lowest temperature simulated for each size.

length reach a clear plateau for all sizes, strongly suggesting that our samples have reached thermal equilibrium. This analysis also provides useful information about the number of sweeps that have to be discarded at the beginning of the Monte Carlo. Finally, the absence of wiggles in the plot, after thermalization has been reached is a fair indication that the quality of our random numbers in good enough for our simulation.

Once we are reasonably sure that a set of thermalized configurations is available, we can proceed to measure the critical temperature. The most accurate way to do so is to apply a so-called finite size scaling analysis. The theoretical background of this method is too complex to be treated here in details. The relevant point of the analysis is that when the system is at its critical temperature, its correlation length, if measured in units of the lattice size is independent of the lattice size. So, when $\beta \simeq \beta_{crit}$, the following relation

$$\frac{\xi(sL,\beta_{crit})}{sL} = \frac{\xi(L,\beta_{crit})}{L}$$
(5.4.13)

must hold for any value of s. One must remember that this equation is only valid in the limit in which both L and sL go to infinity.

Using this approach one measures ξ for several values of β on two lattices of sizes L and sL and plots the ratios defined by (5.4.13). The value of β where the two curves cross is the critical temperature. A better way to perform this analysis is to repeat

this procedure for several values of the lattice size, since (5.4.13) must be valid for any value of s. In principle, we expect that *all* curves meet at the same value of β . In practice, this does not happen, since – as remarked early – (5.4.13) is only valid for large lattices (for which we are not able to compute thermalized samples. The best we can do is to consider the different crossing values for β as a measure of the systematic error of the procedure.

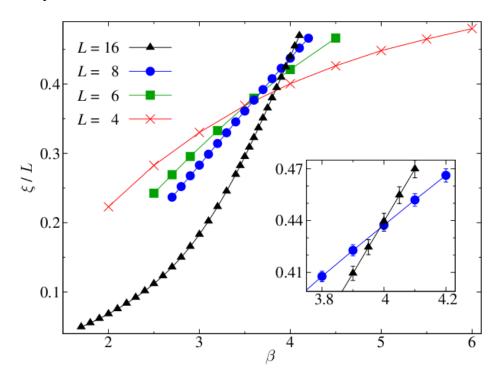


Figure 5.4 – The spin glass correlation length divided by L as a function of β for L = 4, 6, 8 and 16. In the inset we magnify the crossing between the L = 8 and L = 16 curves.

Our results are in Figure 5.4; we plot the correlation length divided by system size for different lattice sizes as a function of the temperature. According to (5.4.13), data should cross if there is a transition. There are clear crossings in the data, though these occur at different temperatures for different sizes. Even though the data represents a considerable computing effort, it is still not enough to be able to extrapolate reliably the intersection temperatures to infinite size. Hence our results are consistent with a second order transition at a finite temperature (whose value is close to $\beta \simeq 4$), but we are not able to measure accurately the asymptotic value of $\beta_{\rm crit}$. Full details on this work are in [3].

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Conclusions

Reconfigurable computing implies a completely different approach to computing with respect to the well known long-standing Von Neumann model. In place of a fixed hardware structure that performs a given programmed algorithm, we configure uncommitted hardware resources so they behave in such a way that they process a data stream according to the requirements of a given algorithm.

There are obvious large potential advantages in this approach; the main advantage is the fact that a given budget of functional units (or logical operators) can be fully configured to perform useful work for the considered algorithm. In recent years, reconfigurable computing has become a viable approach, at the hardware level, as Field Programmable Logic Arrays (FPGA) have established themselves as a technically sound, well understood and well supported technological paradigm for real-life reconfigurable computing. Progress in FPGAs has been dramatic in the last decade and their pace of development is expected to remain stable in the foreseeable future.

In spite of these encouraging situation, reconfigurable computing is still a niche area in high-performance computing (and even more so in conventional computing, with the only possible exception of embedded systems). This is so for reasons that are probably as obvious as the potential advantages of this technology. In brief, a reconfigurable devices defines a potentially enormous space of possible configurations, in which only a very tiny subspace is the optimal (or maybe, even just satisfactory) work-point for the implementation of a given algorithm. Locating this work-point is a non-trivial endeavour for human intervention, that requires very specific skills and long and tedious coding with appropriate Hardware Description Languages. Automatic code-transformation tools, that start from a program written in a traditional programming language and develop an appropriate configurable computing widespread in the computing community, but at present their are hopelessly inadequate. A second key problem in the area is the lack of well-defined, standardized interfaces between the reconfigurable partition of a system and its host computer.

In this thesis, I have described my work in the framework of a large project (the Janus project) that again shows the truly huge advantages offered by reconfigurable

computing for specific applications. The developments described here have increased computing power available to a relevant and well-established application area in statistical physics by almost three orders of magnitude, redefining the state-of-the-art for Monte Carlo simulations in condensed matter. Quite frankly, Janus has reached these very important goals because the people that have developed the project (including the writer of this document) have accepted to go through all the needed, complex and tedious development steps. From this point of view, this work is frankly gives very thin contributions to the broad area of making reconfigurable computing easier to approach. However, Janus had to go around all these hurdles, so Janus experience, including the definition of a reasonably flexible interface between Janus and its host, may be an useful contribution.

The Janus system was developed by a collaboration of research centers and universities in Spain (Zaragoza, Madrid and Extremadura) and Italy (Ferrara and Rome) with the industrial partnership of Eurotech. Its main goal was the development of an efficient system for Monte Carlo simulations of frustrated systems in condensed matter, such as spin glasses. The specific computational features of the relevant simulation algorithms is particularly well adapted to the architecture of currently available reconfigurable devices.

During my PhD studies I have been deeply involved in the several activities associated to the project:

- I have contributed to the overall architectural design of the system;
- I have defined the architectural structure of the interface between Janus and its host computer;
- I have developed in all details the firmware defining the Input Output processor for the system;
- I have defined and developed in all details the set of test programs that have been used to verify the correct behaviour of the Janus hardware;
- I have coordinated the installation and commissioning of the large Janus system installed in Spain and of the smaller machine in Ferrara;
- I have worked on the development of a Janus-based application for the solution of the coloring problem in large random graphs. This is an interesting attempt to adapt reconfigurable hardware to a problem in which addressing is not as regular as in spin-glasses. Encouraging results have been reached, even if it is clear that present FPGAs have strong limitations to support these algorithms.

Physics simulations performed using Janus marks a new standard in the simulation of complex discrete systems by allowing studies that would take centuries on traditional computers. The system that we have developed has a very large sustained performance, an outstanding price-performance ratio and an equally good power-performance ratio. Janus is not a small scale laboratory prototype, but a large high-performance system, able to run large simulations extending the state-of-the-art in the field and is an example of high-performance reconfigurable computing.



Notes on the IOP communication strategy

This appendix describes three ideas that were considered for the structure of the data stream protocol used in the communications between Janus and the host PC. These ideas were discussed within the collaboration and a decision was made taking into account the robustness of the protocol, the overhead in terms of bandwidth and the complexity of the corresponding hardware block. This document, largely taken from an early internal working document, discusses all these issues.

A.1 Overview of the IOP structure

As described in Sections 3.2.3 and 4.3, the IOP made up of the lOlink that plays the role of the input/output interface and a set of others "objects/devices" with different functionalities that have to be driven via Gigabit from a host PC. Figure 3.5 shows a block diagram of the IOP structure.

The lOlink receives as input a 16-bit data stream (dataIn) and a data valid (dvIn) flagging the 16-bit words that the host PC sends to one of the Janus devices. All of these use a module, called StreamRouter (SR) that is a VHDL entity with the logical role to "scan" the data stream in order to recognize a "special" semi-word identifying the target device and assert a signal (called devSelVal) used to select target device.

Another common assumption of the three communications protocol is also that each target device is labelled with an ID. The "special" semi-word of the previous paragraph is therefore a 16-bit mask (bitwise mode) that the StreamRouter recognizes and flags as a bitwise device selector.

StreamRouter analyzed obviously the other valid data words (containing informations for the devices) and forwards them with the suitable data valid (dvOut) to the target devices (on the right of Figure 3.5).

A.2 First idea: the stuff byte

In this protocol it is assumed that a special semi-word (having for instance the value 0xAAAA) is inserted into the stream every time we want to send an information coding the destination of the message. The StreamRouter can therefore easily identify the "stuff semi-word" and consider the following 16-bit data word as the target devices where addresses the following message. Is similar to an escape.

If at the time t the valid data coming from the |Olink| is $D_t = 0xAAAA$ then the StreamRouter waits the next data D_{t+1} (time t+1):

- if $D_{t+1} = D_t$ (green frame in the Figure A.1) then the data D_{t+1} is not labelled as a control data and goes to the target device with the value 0xAAAA. The value of the dvOut is consequently set to high value and the devSelVal is deasserted (see Figure A.2 for a timing diagram of this case);
- if $D_{t+1} \neq D_t$ (violet frame in the Figure A.1) then the data coming from the lOlink is labelled as control-data with the value D_{t+1} . The value of the devSelV is set to high and the dvOut is pulled low (see Figure A.3 for a timing diagram); in this case the data D_{t+1} works as a device selection mask (devSelMask).

In both cases the first data D_t works as "stuff semi-word", behaviour that gives the name to this simple protocol. A typical stream of data is shown in figure A.1.

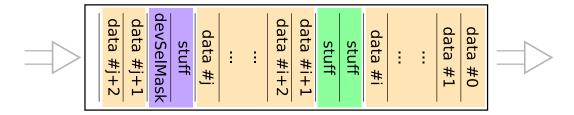


Figure A.1 – Example of a stream with stuff byte.

A.2.1 Which stuff-value?

Using this data stream organization the address of a target device cannot have the same value of the stuff word, or in other words, a mask coding a target device cannot assume the same value of the stuff: in this case indeed the data D_{t+1} is recognized

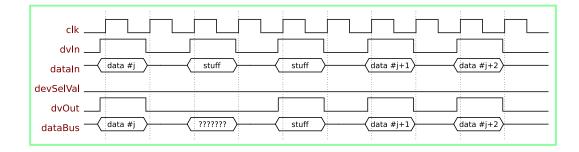


Figure A.2 – Timing diagram in the case of the green frame of Figure A.1.

clk				
dvln				
dataIn	— (data #i)	stuff	 (data #i+1)	data #i+2
devSelVal				
dvOut				
dataBus			 data #i+1	(data #i+2)

Figure A.3 – Timing diagram in the case of the violet frame of Figure A.1.

by the StreamRouter as a data of the previous message instead as a mask for a new message.

This means that one of the $2^{16} = 65536$ possible values coded in the data D_{t+1} should be used as stuff semi-word and cannot be used as mask to select target devices. Ideally the mask that does not make sense is the null mask (0x0000) because the null mask is the mask that selects no device. A natural value for the stuff seems to be therefore the null mask (0x0000).

On the other hand the stuff should have a value that must be not to frequent in the data stream in order to not degrade the bandwidth. For the spin-glass simulation the value 0x0000 could be statistically many present in the data stream. Regular structures as 0xAAAA or 0x5555 are sure a good stuff (are statistically few present in the stream) but are "possible" masks.

A solution for this problem could be to implement a software data encoder/decoder that analyzes the stream on the host PC and converts the stuff from 0x0000 to 0xAAAA or 0x5555 and vice versa before the transmission.

A.2.2 Performance problem

To implement this data stream protocol the host PC sending data to a Janus device have to check each outgoing data word to insert the stuff semi-word where/when needed. This task should be performed run-time with an undefined and unpredictable frequency. This may degrade the performance of the software driving Janus because it could not possible to exploit contiguous memory buffers to pack data in advance.

A.3 Second idea: the tagged stream

The basic idea of this stream protocol is to introduce a tag before each message from host PC to a Janus device. This tag distinguishes between command message and data message. A 16-bit header gives therefore us two different information:

- a tag codifying if the following data words are data or commands;
- a length that the StreamRouter uses as a counter of valid words before to start to receive a new command.

The data stream is structured as shown in Figure A.4: in case of command message, the 16-bit word following the header (green in Figure A.4) will be the mask codifying the target devices.

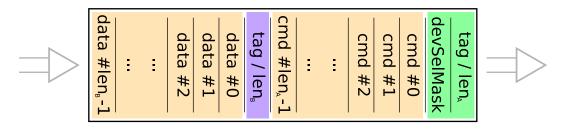


Figure A.4 – Sample of a tagged stream.

A.3.1 Remarks

In this case the Janus software can perform an efficient static memory buffer allocation because the protocol does not require changes of the message contents. The possible performance problems introduced with the stuff byte protocol are here therefore solved.

A second comment is related to the logic structure of this stream organization: we allows indeed the StreamRouter to distinguish between command messages and data messages, that is not required and in some cases could introduce dependences between Janus target devices and StreamRouter. Moreover, this protocol could appear as an overkill for a system such as Janus, because the command messages in Janus are few and very simple, so they do not require a special class of data exchange.

A.4 Third idea: the encapsulated stream

This third communication protocol proposed to handle the Janus data transaction can be considered a generalization of the previous one. Each data set to send from the host PC to a Janus device is called *message* with no difference between command messages and data messages. Each message, moreover has an header (violet frame in Figure 3.6) that gives to the StreamRouter information about the destination (mask containing the ID of the target device) and the length of the message. Details of the protocol are described in 3.2.3 and 4.3.

A.4.1 Pros and cons

We focused our attention on this last protocol and we accept to use it for Janus because:

- it is relatively simple to realize both the hardware side and the software side. Sending header information (basically the mask of the target device and the length of the message) does not require to scan the entire data segment of the message. As described above this allow us to exploit the allocation of buffers into the memory and to send them with no bandwidth disadvantages.
- a message length of 2 16-bit words means that we need a 32 bit counter into the FPGA, but allow us a stream of ~ 10^9 words of 16 bit (~ 64 GB). It is remarkable that a data transmission that load a cubic configuration with lattice size L = 100 for each FPGA, with their couplings, requires $16 \times 4 \times 100^3 + \varepsilon \sim 8$ MB.

Using this protocol remains in all cases the problem that the Janus system become out of control until all the data of a requested operation are not delivered. A stop/reset command of the system is indeed not allowed during a command using this protocol.

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